

**Development of Active Photocatalysts for CO₂ Conversion
by High-Pressure Torsion**

(高圧ねじりによる CO₂ 変換用活性光触媒の開発)

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March 2023

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Abstract

Excessive CO₂ emission from fossil fuel usage has resulted in global warming and environmental crises. To solve this problem, photocatalytic conversion of CO₂ to CO or useful components is a new strategy that has received significant attention. The main challenge in this regard is exploring photocatalysts with high efficiency for CO₂ photoreduction, as discussed in chapter 1. Severe plastic deformation (SPD) through the high-pressure torsion (HPT) process was used in this study to develop novel active photocatalysts for CO₂ conversion. These active photocatalysts were designed based on four main strategies (i) oxygen vacancy and strain engineering, (ii) stabilization of high-pressure phases, (iii) synthesis of defective high-entropy oxides (HEOs), and (iv) synthesis of low-bandgap high-entropy oxynitrides (HEONs).

In chapter 2, to overcome the drawbacks of BiVO₄ for CO₂ conversion, oxygen vacancies and lattice strain are simultaneously introduced in this oxide using the HPT process. The processed material not only shows a low recombination rate and enhanced conduction band level but also exhibits bandgap narrowing. The oxygen-deficient and highly-strained BiVO₄ shows a high photocatalytic CO₂ conversion rate with an activity comparable to the P25 TiO₂ benchmark photocatalyst. The enhancement of photocatalytic activity is discussed based on the modification of band structure, enhanced light absorbance, the lifetime of excited electrons, and the role of oxygen vacancies as activation sites for CO₂ photoreduction. This chapter introduces a feasible pathway to develop active photocatalysts for CO₂ conversion by lattice strain and defect engineering.

In chapter 3, the nanocrystalline high-pressure TiO₂-II polymorph, synthesized by HPT, shows higher photocurrent and better photocatalytic activity for CO₂ to CO conversion compared to the anatase phase. The photocatalytic activity of the material is further enhanced by the thermal annihilation of oxygen vacancies generated in the bulk of TiO₂-II during the high-pressure treatment. The high potential of the TiO₂-II phase for CO₂ conversion is discussed based on the light absorbance, band structure, charge carrier mobility and CO₂ adsorption.

In chapter 4 and inspired by the inherent defective and strained structure of HEOs, photocatalytic CO₂ conversion is examined on a dual-phase TiZrNbHfTaO₁₁ synthesized by a two-step mechanical alloying by HPT and high-temperature oxidation. The HEO, which had various structural defects, showed simultaneous photocatalytic activity for CO₂ to CO and H₂O to H₂

conversion without the addition of a co-catalyst. The photocatalytic activity of this HEO for CO₂ conversion was better than conventional photocatalysts such as anatase TiO₂ and BiVO₄ and similar to P25 TiO₂. The high activity of HEO was discussed in terms of lattice defects, lattice strain, light absorbance, band structure, photocurrent generation and charge carrier mobility to activation centers. The current study confirms the high potential of HEOs as a new family of photocatalysts for CO₂ conversion.

In chapter 5, a HEON was introduced as an active photocatalyst for photoreduction. The material had a chemical composition of TiZrNbHfTaO₆N₃ and was produced by HPT processing followed by oxidation and nitriding. It showed higher photocatalytic CO₂ to CO conversion compared to corresponding high-entropy oxide, benchmark photocatalyst P25 TiO₂, and almost all catalysts introduced in the literature. The high activity of this oxynitride, which also showed good chemical stability, was attributed to the large absorbance of light and easy separation of electrons and holes, the low recombination of charge carriers, and the high CO₂ adsorption on the surface. These findings introduce high-entropy oxynitrides as promising photocatalysts for CO₂ photoreduction.

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Chapter 1. Introduction

1.1. Background

Nowadays, environmental crises especially global warming caused by CO₂ emissions from burning fossil fuels and humankind activities are considered one of the most significant challenges in the world. Reduction of CO₂ to reactive CO gas or useful components and fuels such as CH₄ and CH₃OH using photocatalysts is one of the clean and new strategies which is developing rapidly [1-3]. The function of photocatalysts is the acceleration of reduction and oxidation reactions under light irradiation. Reduction and oxidation reactions occur on the surface of a photocatalyst using the electrons and holes, respectively to promote the CO₂ conversion [1-3]. In photocatalytic CO₂ conversion, photoexcited electrons contribute to the conversion of CO₂ to reactive and useful components such as carbon monoxide (CO), formic acid (HCOOH), formaldehyde (HCHO), methanol (CH₃OH) and methane (CH₄) [1-3]. The mechanism of photocatalytic CO₂ conversion is shown in Figs. 1a. To perform the reactions, a photocatalyst should have some features including high light absorbance, appropriate band structure, low recombination rate of electrons and holes, easy migration of charge carriers, and high surface affinity to adsorb CO₂ with abundant active sites [3,4].

Semiconductors such as TiO₂ [5-7], g-C₃N₄ [8,9], and BiVO₄ [10-12] are typical photocatalysts that have been engineered by various strategies to enhance the catalytic efficiency for CO₂ conversion. Doping with impurities [13-15], defect engineering [16,17], strain engineering [18,19], mesoporous structure production [20], and heterojunction introduction [21,22] are some of the most promising strategies which have been used so far to improve the optical properties and catalytic activity of various photocatalysts. Among these strategies, doping with impurities is the most investigated and feasible method, but impurities can increase the recombination rate of electrons and holes [13-15]. Therefore, finding new strategies to improve the photocatalytic activity and suppress the recombination rate of electrons and holes is a key issue.

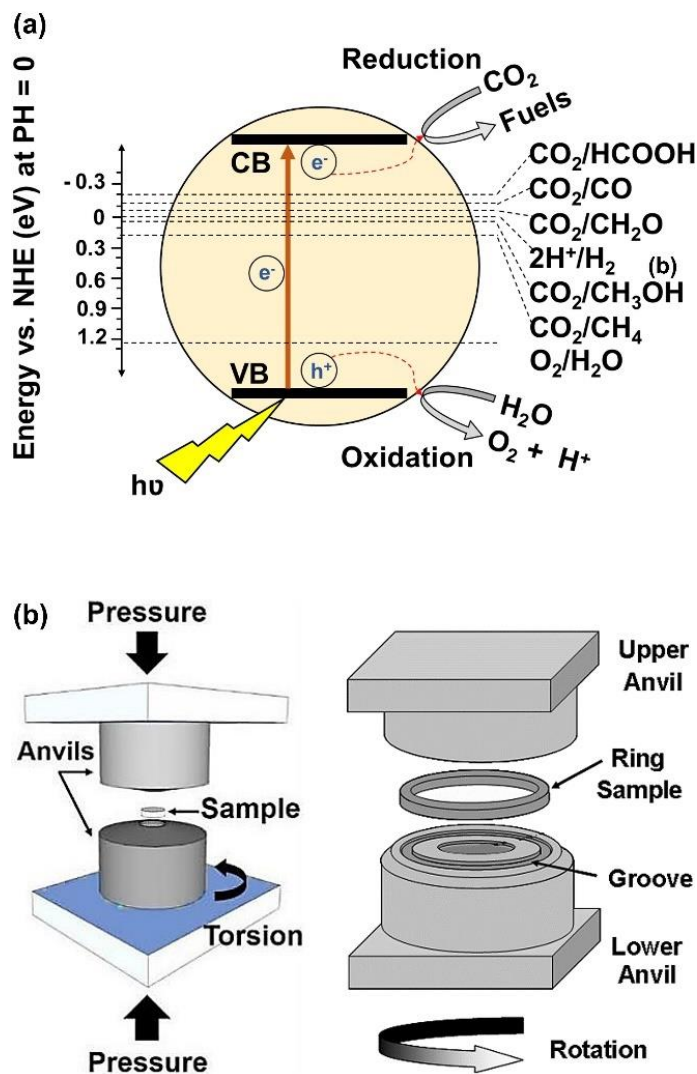


Fig. 1.1. Schematic illustration of (a) photocatalytic CO₂ conversion and (b) high-pressure torsion [41].

In this regard, severe plastic deformation (SPD) through the high-pressure torsion (HPT) method, which is schematically shown in Fig. 1.1b and mainly used for nanostructuring of metallic materials showed the potential to improve the electronic structure of semiconductors, narrow their bandgap and enhance the photocatalytic activity for water splitting [23-30] and dye degradation [31-34]. Fig. 1.2 shows the impact of HPT on photocatalytic H₂ evolution and dye degradation. In Fig. 1.2a it is observed that after HPT processing, the H₂ generation amount for ZrO₂ increased significantly [26]. Fig. 1.2b demonstrates the HPT method improved the photocatalytic activity of

Al_2O_3 effectively while the initial powder of this material could not degrade the dye [32]. It can be observed that the photocatalytic activity of Al_2O_3 after HPT is considerable compared to TiO_2 as a conventional semiconductor for this application [32]. Despite the application of HPT for these two photocatalytic processes, there are not any reports in the literature on the potential of this method to improve or synthesize the materials for photocatalytic CO_2 conversion.

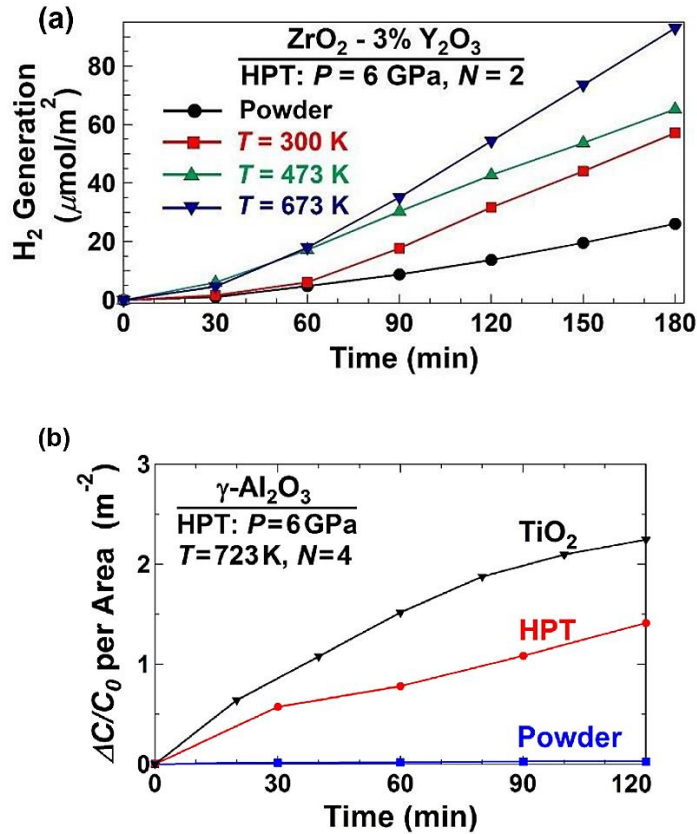


Fig. 1.2. Impact of HPT on photocatalytic H_2 evolution and dye degradation. (a) Photocatalytic H_2 production amount versus irradiation time for ZrO_2 and (b) photocatalytic dye degradation amount versus irradiation time for Al_2O_3 before and after HPT processing compared with TiO_2 [26,32].

1.2. High-pressure torsion (HPT)

The SPD process is usually used to form ultrafine-grained and nanostructured materials with enhanced mechanical and functional properties [35-37]. SPD has various methods including the HPT method, introduced by Bridgman in 1935 for the first time [35]. In this method, torsional strain under high pressure is applied to induce large plastic strain in various ranges of materials [37]. In the HPT process, a disc-shape sample that has typically a 10 mm diameter or a ring sample

is inserted between two anvils under high pressure and strained by rotation of anvils against each other [38,39]. A schematic illustration of HPT is shown in Fig. 1.1b. The applied shear strain (γ) to the sample can be calculated by the following equation [40].

$$\gamma = \frac{2\pi r N}{h} \quad (1.1)$$

where r , N and h are the distance from the center of the disc or ring, the number of turns and the sample height, respectively.

In addition to grain refinement, the HPT is widely used for hardening of pure metals [41,42], mechanical alloying of miscible and immiscible systems [43,44] and plastic deformation of hard-to-deform materials [45-47]. Phase transformation, consolidation of powders, the introduction of defects such as oxygen vacancies and dislocations, narrowing of the bandgap and improvement of functionalities were frequently reported after HPT processing [48]. The HPT method also can be utilized for the plastic deformation of ceramics which are hard and brittle materials at ambient temperature [35]. The presence of covalent or ionic bonding in ceramics results in their lower grain size after HPT processing compared to metals [35]. Moreover, such bonding features result in the formation of large densities of defects such as vacancies and dislocations which can improve the properties and functionality of ceramics [35].

Despite the high potential of ceramics for various applications, there are some limited publications on the effect of HPT on the structure, properties and functionality of ceramics. These HPT-processed ceramics, including oxides, nitrides, oxynitride and borides, have been investigated for photocatalysis [23-34], phase transformation [49-55], electrocatalysis [56], photocurrent generation [57,58], dielectric properties [59,60], bandgap narrowing [61], optical properties [62-65], mechanical properties [66,67], Li-ion batteries [68] and microstructural features [69-71]. Furthermore, the HPT method was also utilized to synthesize new ceramics for various applications [27,29]. A list of ceramics treated by HPT for various properties and applications is given in Table 1.1.

Table. 1.1. Summary of major publications about ceramics treated by HPT and their major applications or features.

Materials	Investigated properties and application	Reference
Photocatalysis		
TiO ₂ -II	Photocatalytic activity for hydrogen production	[23]
ZnO	Photocatalytic activity for dye degradation	[31]
γ -Al ₂ O ₃	Photocatalytic activity for dye degradation	[32]
MgO	Photocatalytic activity for dye degradation	[33]
ZrO ₂	Photocatalytic activity for hydrogen production	[26]
SiO ₂	Photocatalytic activity for dye degradation	[34]
CsTaO ₃ , LiTaO ₃	Photocatalytic activity for hydrogen production	[24]
GaN-ZnO	Photocatalytic activity for hydrogen production	[25]
TiHfZrNbTaO ₁₁	Photocatalytic activity for hydrogen production	[27]
TiZrNbTaWO ₁₂	Photocatalytic activity for oxygen production	[30]
TiZrHfNbTaO ₆ N ₃	Photocatalytic activity for hydrogen production	[29]
TiO ₂ -ZnO	Photocatalytic activity for hydrogen production	[28]
Reviews		
	Review on HPT	[40]
Oxides	Review on HPT of oxides	[48]
Phase transformation		
SiO ₂ , VO ₂	Phase transformation	[49]
TiO ₂	Grain coarsening and phase transformation	[50]
ZrO ₂	Phase transformation by ball milling and HPT	[57]
ZrO ₂	Allotropic phase transformations	[52]
TiO ₂	Plastic strain and phase transformation	[53]
BN	FEM modeling of plastic flow and strain-induced phase transformation	[54]
BN	Coupled elastoplasticity and plastic strain-induced phase transformation	[55]
Electrocatalysis		
TiO ₂ -II	Electrocatalysis for hydrogen generation	[56]
Photocurrent		
Bi ₂ O ₃	Enhanced photocurrent generation	[57]
TiO ₂ -II	Visible light photocurrent generation	[58]
Dielectric properties		
BaTiO ₃	Optical and dielectric properties	[59]

CuO	Dielectric properties	[60]
Bandgap investigation		
ZnO	Bandgap narrowing	[61]
Optical properties		
Y ₂ O ₃	Optical properties	[62]
CuO, Y ₃ Fe ₅ O ₁₂ , FeBO ₃	Optical properties and electronic structure	[63]
Cu ₂ O, CuO	Middle infrared absorption and X-ray absorption	[64]
CuO, Y ₃ Fe ₅ O ₁₂ , FeBO ₃	Optical properties	[65]
Mechanical properties		
α -Al ₂ O ₃	Microstructure and mechanical properties	[66]
Fe _{53.3} Ni _{26.5} B _{20.2} , Co _{28.2} Fe _{38.9} Cr _{15.4} Si _{0.3} B _{17.2}	Microstructure and mechanical properties	[67]
Lithium-ion batteries		
Fe ₃ O ₄	Lithium-ion batteries	[68]
Microstructural features		
ZnO	Plastic flow and microstructural instabilities	[69]
YBa ₂ Cu ₃ O _y	Microstructural investigation	[70]
Fe _{71.2} Cr _{22.7} Mn _{1.3} N _{4.8}	Microstructural features	[71]

1.3. Fundamentals and mechanism of photocatalytic CO₂ conversion

In photocatalysis, which is also called artificial photosynthesis, electrons in the valence band absorb the light photons, separate from the holes and transfer to the conduction band of the photocatalyst to form the electron-hole charge carriers. The charge carriers then migrate to the surface of the photocatalyst and finally take part in various chemical reactions [3]. In this process, both reduction and oxidation reactions occur on the surface of the photocatalyst, provided that the thermodynamic and kinetic conditions are satisfied [3]. From the thermodynamic point of view, a photocatalytic process will perform when the reduction and oxidation reaction potentials are between the valence band and the conduction band of a photocatalyst [3]. For the reduction reactions, potentials lower than the valence band top are desirable, while for the oxidation reactions, potentials higher than the conduction band bottom are desirable. From the kinetic point of view, the electron-hole separation time should be fast, and their recombination time should be slow. Here, the fundamentals of three main photocatalytic reactions are mentioned.

In photocatalytic CO₂ conversion, electrons in the conduction band along with protons (H⁺) produced from water oxidation in the valence band, take part in various reduction reactions leading

to the formation of CO, CH₄ and some other hydrocarbons as shown in Fig. 1.1a. The first step in the photocatalytic CO₂ conversion process is the adsorption of CO₂ molecule on the surface of the photocatalyst which occurs in three modes including oxygen coordination, carbon coordination and mixed coordination [72,73], as shown in Fig. 1.3. These CO₂ adsorption modes determine the reaction pathway for photocatalytic CO₂ conversion. For instance, the oxygen coordination mode, with the bidentate bonding of oxygen atoms with the surface of the photocatalyst, leads to the fabrication of formate anion as an intermediate product and formic acid as the final product [73]. If the CO₂ molecules have the carbon coordination mode, with the monodentate binding of carbon and photocatalyst surface, the reaction results in carboxyl radical production [73]. After the adsorption step of CO₂ to the surface of the photocatalyst, the conversion reactions occur via different pathways.

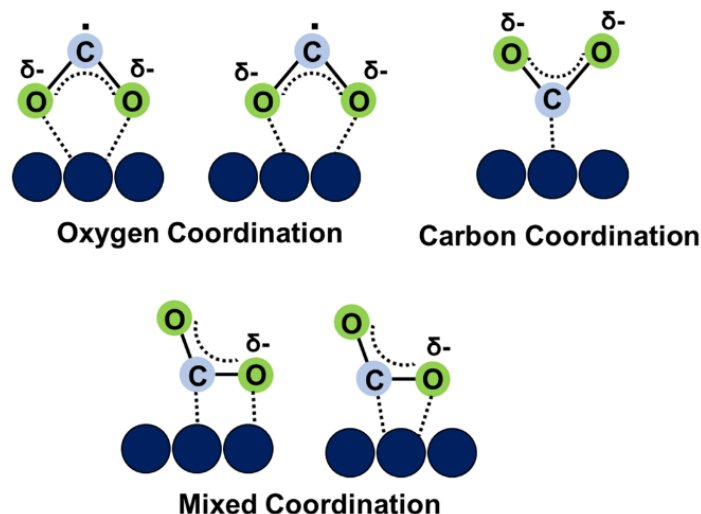


Fig. 1.3. CO₂ adsorption modes on the surface of photocatalyst.

It was reported that there are three main conventional pathways for CO₂ photoreduction: (i) carbene pathway (ii) formaldehyde pathway and (iii) glyoxal pathway [74,16]. The reactions for these three pathways are presented in Table 1.2. In all these pathways, the CO₂^{•-} radical is the main intermediate product which is formed by the reaction of adsorbed CO₂ and an electron. In the carbene pathway, CO₂^{•-} converts to CO[•] and leads to the production of CO, CH₂, CH₄ and CH₃OH. In this pathway, carbon coordination is the main binding mode. In the formaldehyde pathway, CO₂^{•-} converts to [•]COOH to produce HCOOH, CH₃ OH and CH₄ [74,75]. In this pathway,

$\text{CO}_2^{\bullet-}$ attaches to the catalyst by monodentate or bidentate binding. In monodentate binding, either one oxygen in $\text{CO}_2^{\bullet-}$ or one carbon atom in $\text{CO}_2^{\bullet-}$ are bonded to the surface of the photocatalyst. In the bidentate mode, both oxygen atoms in $\text{CO}_2^{\bullet-}$ are bonded to the photocatalyst surface [74]. In the glyoxal pathway, C_2 components are formed by some complicated reactions, although there are limitations to producing C_2 due to the low selectivity of the photocatalyst. In this pathway, $\text{CO}_2^{\bullet-}$ converts to the formyl radical HOC^{\bullet} to form $\text{C}_2\text{H}_2\text{O}_2$, $\text{C}_2\text{H}_4\text{O}_2$, $\text{C}_2\text{H}_4\text{O}$, CO and CH_4 [74,16]. In all these pathways, CO_2 adsorption and activation are key parameters to adjust the selectivity of photocatalysts and this should be empowered by some strategies such as oxygen vacancy generation and strain engineering [73].

Table 1.2. Reactions of CO_2 photoreduction pathways [74].

Carbene Pathway	Formaldehyde Pathway	Glyoxal Pathway
(1) $\text{CO}_2 + e^- \rightarrow \text{CO}_2^{\bullet-}$	(1) $\text{CO}_2 + e^- \rightarrow \text{CO}_2^{\bullet-}$	(1) $\text{CO}_2 + e^- \rightarrow \text{CO}_2^{\bullet-}$
(2) $\text{CO}_2^{\bullet-} + e^- + \text{H}^+ \rightarrow \text{CO} + \text{OH}^-$	(2) $\text{CO}_2^{\bullet-} + \text{H}^+ \rightarrow \text{COOH}^{\bullet}$	(2) $\text{CO}_2^{\bullet-} + e^- + \text{H}^+ \rightarrow \text{CHOO}^-$
(3) $\text{CO} + e^- \rightarrow \text{CO}^{\bullet-}$	(3) $\text{COOH}^{\bullet} + e^- + \text{H}^+ \rightarrow \text{HCOOH}$	(3) $\text{CHOO}^- + \text{H}^+ \rightarrow \text{HCOOH}$
(4) $\text{CO}^{\bullet-} + e^- + \text{H}^+ \rightarrow \text{C} + \text{OH}^-$	(4) $\text{HCOOH} + e^- + \text{H}^+ \rightarrow \text{H}_3\text{OOC}^{\bullet}$	(4) $\text{HCOOH} + e^- \rightarrow \text{HOC}^{\bullet}$
(5) $\text{C} + e^- + \text{H}^+ \rightarrow \text{CH}^{\bullet}$	(5) $\text{HCOOH}_2^{\bullet} + e^- + \text{H}^+ \rightarrow \text{HCOH} + \text{H}_2\text{O}$	(5) $\text{HOC}^{\bullet} + \text{OH}^- \rightarrow \text{C}_2\text{H}_2\text{O}_2$
(6) $\text{CH}^{\bullet} + e^- + \text{H}^+ \rightarrow \text{CH}_2$	(6) $\text{HCOH} + e^- \rightarrow \text{H}_2\text{C}^{\bullet}\text{O}^-$	(6) $\text{C}_2\text{H}_2\text{O}_2 + e^- + \text{H}^+ \rightarrow \text{H}_3\text{O}_2\text{C}_2^{\bullet}$
(7) $\text{CH}_2 + e^- + \text{H}^+ \rightarrow \text{CH}_3^{\bullet}$	(7) $\text{H}_2\text{C}^{\bullet}\text{O}^- + \text{H}^+ \rightarrow \text{H}_2\text{OHC}^{\bullet}$	(7) $\text{H}_3\text{O}_2\text{C}_2^{\bullet} + e^- + \text{H}^+ \rightarrow \text{C}_2\text{H}_4\text{O}_2$
(8) $\text{CH}_3^{\bullet} + e^- + \text{H}^+ \rightarrow \text{CH}_4$	(8) $\text{H}_2\text{OHC}^{\bullet} + e^- + \text{H}^+ \rightarrow \text{CH}_3\text{OH}$	(8) $\text{C}_2\text{H}_4\text{O}_2 + e^- + \text{H}^+ \rightarrow \text{H}_3\text{OC}_2^{\bullet} + \text{H}_2\text{O}$
(9) $\text{CH}_3^{\bullet} + \text{OH}^- \rightarrow \text{CH}_3\text{OH}$	(9) $\text{CH}_3\text{OH} + e^- + \text{H}^+ \rightarrow \text{CH}_3^{\bullet} + \text{H}_2\text{O}$	(9) $\text{H}_3\text{OC}_2^{\bullet} + e^- + \text{H}^+ \rightarrow \text{C}_2\text{H}_4\text{O}$
	(10) $\text{CH}_3^{\bullet} + e^- + \text{H}^+ \rightarrow \text{CH}_4$	(10) $\text{C}_2\text{H}_4\text{O} + \text{h}^+ \rightarrow \text{H}_3\text{OC}_2^{\bullet} + \text{H}^+$
		(11) $\text{H}_3\text{OC}_2^{\bullet} \rightarrow \text{CH}_3^{\bullet} + \text{CO}$
		(12) $\text{CH}_3^{\bullet} + e^- + \text{H}^+ \rightarrow \text{CH}_4$

1.4. Objectives and outline of thesis

Numerous studies in recent years reported that the HPT method can be used to improve the photocatalytic activity of ceramics for dye degradation and water splitting [23-34]. Bandgap narrowing for the easier transition of photoexcited electrons from the valence band to the conduction band, increasing the light absorbance in the visible region of light, electronic band structure alignment, decreasing the recombination rate of electrons and holes and accelerating the electron-hole separation and migration are some phenomena responsible for the high activity of photocatalysts processed or synthesized by HPT [23-34]. Since there are significant efforts

worldwide to reduce the CO₂ emissions by its conversion, in the present thesis, the HPT method is utilized for the first time to develop active photocatalysts for CO₂ conversion by using four main strategies.

- 1) Simultaneous strain and oxygen vacancy engineering
- 2) Introduction of high-pressure phases
- 3) Formation of defective high-entropy phases
- 4) Production of low-bandgap high-entropy oxynitride phases

This thesis contains six chapters:

In chapter 1, the research background of this thesis is summarized. A brief description of the HPT method and mechanism of photocatalytic CO₂ conversion is presented.

In chapter 2, BiVO₂ as a conventional photocatalyst is treated by HPT using simultaneous strain and oxygen vacancy engineering. Impact of HPT on oxygen vacancy concentration and lattice strain is evaluated by various characterization methods. Electronic structure, charge carrier dynamics and photocatalytic CO₂ conversion are investigated to realize the influence of employed strain and oxygen vacancy engineering on catalytic efficiency.

In chapter 3, the high-pressure TiO₂-II phase is synthesized by HPT to investigate its potential for photocatalytic CO₂ conversion for the first time. Impact of this new phase on light absorbance, narrowing the bandgap, charge carrier dynamics and photocatalytic activity of TiO₂ for CO₂ conversion is evaluated.

In chapter 4, a new high-entropy oxide is produced using the HPT method and subsequent high-temperature oxidation and evaluated for photocatalytic CO₂ conversion for the first time. This oxide is characterized by various technique in terms of crystal structure, microstructure, light absorbance, electronic structure, charge carrier dynamic and photocatalytic activity for CO₂ conversion.

In chapter 5, to decrease the bandgap and improve the photocatalytic activity of fabricated high-entropy oxide in previous chapter, a high-entropy oxynitride is synthesized and evaluated by various characterization methods for photocatalytic properties and activity for CO₂ conversion.

In chapter 6, the conclusions and outlook of this thesis are summarized.

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Chapter 2. Enhanced CO₂ conversion on highly-strained and oxygen-deficient BiVO₄ photocatalyst

2.1. Introduction

Among various photocatalysts developed for photocatalytic CO₂ conversion, BiVO₄ (as a mixed-metal oxide or salt) has recently received significant attention due to its high stability and low direct bandgap (2.4-2.5 eV) [1,2]. The main drawbacks of BiVO₄ which limit its photocatalytic CO₂ conversion efficiency are its low conduction band energy level and high rate of electron and hole recombination [2,3]. To improve the photocatalytic CO₂ conversion on BiVO₄, most reported studies introduced heterojunctions as a strategy: BiVO₄/C₄N₃ [3], BiVO₄{010}/Au/Cu₂O [4], BiVO₄/C/Cu₂O [5], ZnIn₂S₄/BiVO₄ [6], BiVO₄/Bi₄Ti₃O₁₂ [7], and Cd_xZn_{1-x}S@Au/BiVO₄ [8]. Doping with other elements such as Bi, Cu and Ag [9-13] and formation of vanadium-deficient layers [14] were also reported as other strategies to enhance the photocatalytic CO₂ reduction on BiVO₄. Despite the well-known importance of engineering the defects and lattice strain on photocatalytic CO₂ conversion, there have been limited studies on this issue for pure BiVO₄ without the addition of impurities.

There are some reported methods in the literature to introduce the defects and strain in photocatalysts. Thermal treatment under hydrogen or inert atmospheres at high temperature, chemical reduction, vacuum activation, ultraviolet irradiation, phase transformation via fast heating, plasma etching, lithium-induced conversion and low-valence metal doping are some methods used to introduce lattice defects such as oxygen vacancies [15-17]. Most of these methods need chemical reactions, high temperature or impurity addition. In the case of BiVO₄ for photocatalytic CO₂ conversion, doping with other elements or producing oxygen- and vanadium-deficient layers by chemical reactions are the methods utilized for oxygen vacancy generation [9-14]. There are few studies on strain engineering of BiVO₄-based photocatalysts, although strain engineering can be conducted by the formation of layered and core-shell structures or by chemical methods such as solvothermal technique [18,19].

In this chapter, BiVO₄ was treated by a high-pressure torsion (HPT) method to simultaneously introduce both defects and strain and improve the electronic structure, carrier dynamics and photocatalysis CO₂ conversion. The HPT method, in which a material is torsionally

strained under high pressure [20], was used in this study because the method can introduce defects and strain in oxides without a need for impurity atoms [21,22]. The highly-strained and oxygen-deficient BiVO₄ exhibited narrow optical bandgap, enhanced conduction band bottom, low electron-hole recombination rate and high photocatalytic CO₂ conversion, which was comparable to the conversion rate on P25 TiO₂ as a benchmark photocatalyst. This study introduces an effective strategy based on defect and strain engineering to enhance the photocatalytic CO₂ conversion on BiVO₄, although the strategy is easily applicable to a wider range of photocatalysts for CO₂ conversion.

2.2. Experimental procedure

2.2.1. Sample synthesis

Initial BiVO₄ with 99.9% purity was purchased from Sigma Aldrich and further processed by HPT, as schematically shown in Fig. 2.1a [20]. For HPT processing, about 410 mg of initial powder was first compressed under 380 MPa to form pellets with 10 mm diameter and 0.8 mm thickness. The pellets were then compressed between two HPT anvils under a pressure of 6 GPa at room temperature (300 K). When the pressure was stabilized, the sample was plastically strained by rotating the lower anvil with respect to the upper anvil for either 0.25, 1 or 4 turns with a rotation speed of 1 rpm. It should be noted that larger plastic deformation is introduced in the sample with increasing the number of HPT turns [20]. The appearance of the initial powder and of the three samples processed by HPT after compacting to disc shapes are shown in Fig. 2.1b. The color of samples became darker with increasing the number of HPT turns and changed from yellow to orange and rose, respectively. These color changes indicate that photons with lower energies can be absorbed after HPT processing, as shown in the Itten color wheel of Fig. 2.1c [23].

2.2.2. Characterization

The initial powder and the samples processed by HPT for 0.25, 1 and 4 turns were characterized by various techniques, as described below.

- Crystal structure was examined by X-ray diffraction (XRD) utilizing the Cu K α radiation ($\lambda = 0.1542$ nm wavelength) and a zero-diffraction amorphous silicon holder with a scanning speed of 1 °/min. A Rigaku SmartLab 9kW AMK machine equipped with automatic alignment systems for beam, detector and sample (height and angle) was used for the XRD analysis. The

precision of diffraction angle, confirmed by using a standard crystalline silicon sample, was better than 0.01° . XRD profiles were analyzed using PDXL software to determine the lattice parameters and lattice volume. Moreover, the crystallite size was determined by the Halder-Wagner method [24].

•

$$\left(\frac{\beta}{\tan^2\theta}\right)^2 = \frac{K\lambda}{D} \left(\frac{\beta \cos(\theta)}{\sin^2\theta}\right) + 16\varepsilon^2 \quad (2.1)$$

where, β , θ , K , λ , D and ε are full width at half maximum after subtracting the instrumental broadening, Bragg angle, shape factor, wavelength, crystallite size and lattice distortion (caused by planar and line defects), respectively.

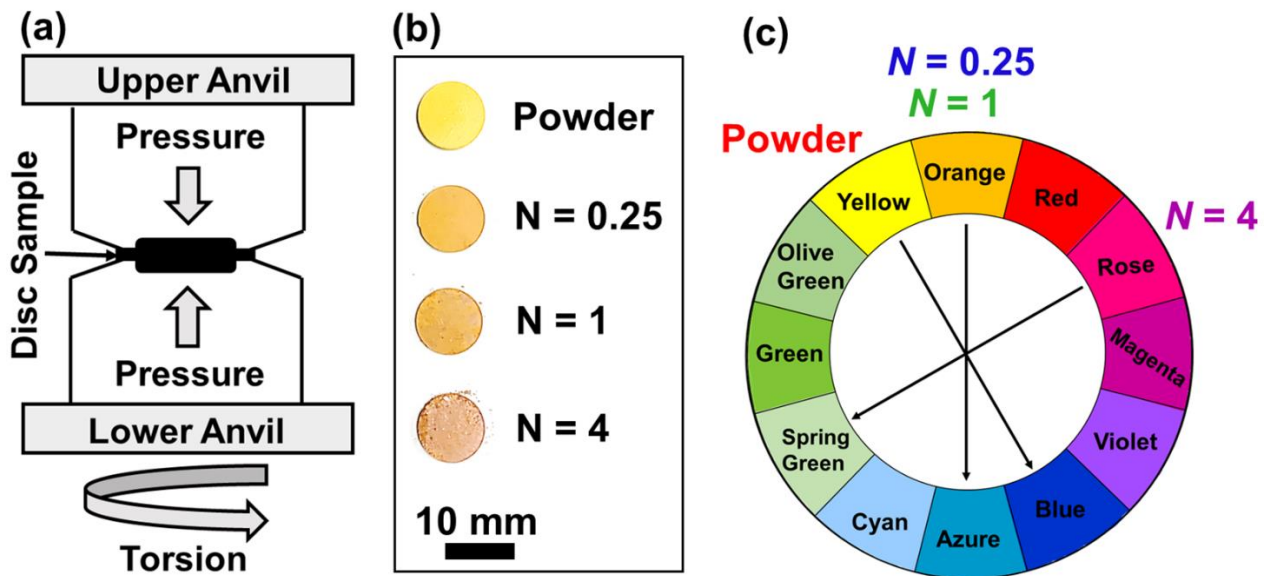


Fig. 2.1. (a) Schematics of HPT method, (b) appearance of BiVO₄ samples before and after HPT processing for 0.25, 1 and 4 turns, and (c) color of samples in comparison with Itten color wheel. Samples in (b) were compacted to disc form to show color changes clearly.

• Crystal structure was also examined by Raman spectroscopy using a laser source with $\lambda = 532$ nm wavelength. A Renishaw inVia Raman WiRE 4 machine equipped with an optic alignment

tool was used. The precision of Raman shift, confirmed using a standard single-crystal silicon sample, was better than 0.1 cm^{-1} .

- Particle size was measured by the dynamic light scattering (DLS) method utilizing a Zetasizer Nano-S machine equipped with a 4 mW He-Ne laser (633 nm) with 173° diffraction angle. The specific surface area was measured from the average particle size as: surface area (m^2g^{-1}) = $6000 / [\text{average particle size (nm)} * (\text{density (gcm}^{-3}\text{)})]$.
- Microstructure was investigated by transmission electron microscopy (TEM) using the bright-field (BF) images, dark-field (DF) images, selected area electron diffraction (SAED) analysis, high-resolution images and fast Fourier transform (FFT) diffractograms. For these observations, small amounts of samples were crushed in ethanol and dispersed on a carbon grid and examined in an aberration-corrected TEM under a voltage of 200 keV.
- X-ray photoelectron spectroscopy (XPS) using the Al $K\alpha$ radiation with a wavelength of $\lambda = 0.834 \text{ nm}$ was used to determine the oxidation state of elements and the presence of oxygen vacancies. Peak deconvolution of O 1s XPS spectrum was also performed to calculate the oxygen vacancies concentration [25].
- Electron paramagnetic resonance (EPR) using a 9.4688 GHz microwave source was used to investigate the formation of oxygen vacancies.
- Light absorbance of samples was examined by UV-vis diffuse reflectance spectroscopy and then the optical bandgap was calculated by the Kubelka-Munk analysis.
- The valence band top was determined using the ultraviolet photoelectron spectroscopy (UPS) with the He-I UV irradiation and a DC bias of -4 V. The conduction band bottom was determined by adding the bandgap to the valence band top.
- Recombination rate of electrons and holes was investigated by steady-state photoluminescence (PL) emission spectroscopy with a 325 nm laser light source.
- Time-resolved photoluminescence decay (PL decay) with a 285 nm laser source was performed to investigate the average lifetime of excited electrons.

2.2.3. Photocurrent test

Photocurrent measurement was conducted under the full arc of a xenon lamp (300 W, Asashi Spectra, Japan, HAL-320W) in a 1 M Na_2SO_4 electrolyte to study the electron separation potential. A thin film of BiVO_4 samples on FTO (fluorine-doped tin oxide, $8 \Omega/\text{sq}$, Aldrich) glass

was used as a working electrode, and Pt wire and Ag/AgCl were utilized as the counter and reference electrodes, respectively. To measure the photocurrent generation, an electrochemical analyzer (ALS C. Ltd., Japan, BAS1230C) was used in potentiostatic amperometry mode under a voltage of 0.7 V versus Ag/AgCl (0.5 min light ON and 1 min light OFF). To prepare the thin film samples, 5 mg of BiVO₄ were crushed in 0.2 mL ethanol and dispersed on 2.25 mm thick FTO glass with 15×25 mm² surface area and baked in a muffle furnace at 473 K for 1 h, as described in detail in an earlier study [26].

2.2.4. Photocatalysis test

For the photocatalytic CO₂ conversion test, 120 mg of each sample was dissolved in 500 mL of deionized water and 1 M NaHCO₃ under CO₂ bubbling with a flow rate of 3 mLmin⁻¹. Photocatalytic test was performed in an 858 mL cylindrical quartz photoreactor with continuous flow, as described earlier [27]. A high-pressure mercury light source with 0.5 Wcm⁻² light intensity (Sen Lights continuous flow Corporation, HL400BH-8, 400 W) with no filtering was used for irradiation. The process was first performed without irradiation for 2 h to be sure about the absence of reaction products. The light source was then placed in the inner space of the reactor and the reactor temperature was adjusted to 288 K by a water chiller. The reaction mixture was stirred continuously by a magnetic stirrer and the CO₂ gas was injected into the reactor by an inlet hole on the top of the reactor. The reaction products continuously flowed out of the reactor through an outlet hole on the top of the reactor, as shown in Fig. 2.1a. The outlet gas from the reactor was divided into two parts and moved either to vent or a gas chromatograph (Shimadzu GC-8A, Ar Carrier). The formation of CO and CH₄ was determined by a flame ionization detector equipped with a methanizer (Shimadzu MTN-1). To measure H₂ and O₂ production rate, a thermal conductivity detector was also used. Fig. 2.2a shows the system used for photocatalytic CO₂ conversion along with the relevant spectral composition of the mercury light source in Fig. 2.2b. It should be noted that a blank test without the catalyst addition but with light irradiation was also conducted to be sure about the absence of reaction products.

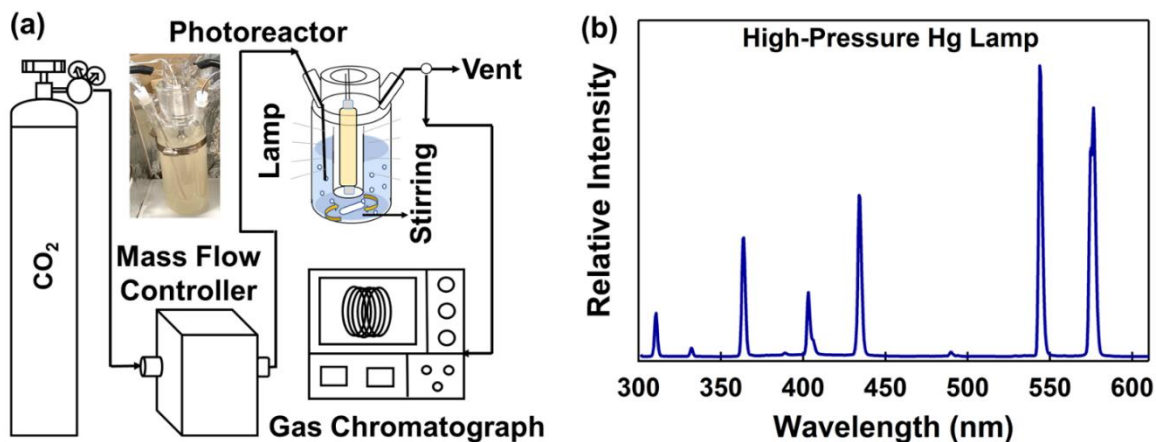


Fig. 2.2. (a) Schematics of photocatalytic experiment for CO₂ conversion, (b) spectral composition of light source employed for photocatalytic experiment.

2.3. Results

2.3.1. Lattice strain and microstructure

Fig. 2.3a shows the crystal structure analysis of four samples using the XRD method. It is evident that the initial powder and the HPT-processed samples have a monoclinic phase (C12/c1 space group with $a = 0.519$ nm, $b = 1.170$ nm, $c = 0.509$ nm; $\alpha = 90^\circ$, $\beta = 90.38^\circ$ and $\gamma = 90^\circ$). XRD profiles in a higher magnification are shown in **Fig. 2.3b** for the (004) atomic plane with a diffraction angle of $\sim 30.5^\circ$. There is a peak shift to higher angles after HPT processing and the shift becomes more significant with increasing the number of turns (i.e., with increasing plastic deformation). The peak shift is an indication of lattice expansion, which can be due to lattice strain and vacancy-type defect formation [28,29]. Furthermore, a peak broadening occurs, and the broadening becomes more significant with increasing the number of turns. The peak broadening is usually due to the formation of dislocation-type defects or planar imperfections such as grain boundaries [30]. **Fig. 2.3c** shows the crystallite size of four samples measured by the Halder-Wagner method. It is evident that the crystallite size increases from 65.7 nm for the initial powder to 24.6, 12.2 and 8.8 nm for the samples proceeded by HPT for 0.25, 1 and 4 turns, respectively. **Fig. 2.3a** also summarizes the volumetric lattice strain which was calculated as the increase in the lattice volume of the HPT-processed sample compared to the lattice volume of the initial powder, $\Delta V/V_0$. The estimated volumetric lattice strains for the samples proceeded by HPT for 0.25, 1 and 4 turns are 0.52%, 0.68% and 0.86%, respectively.

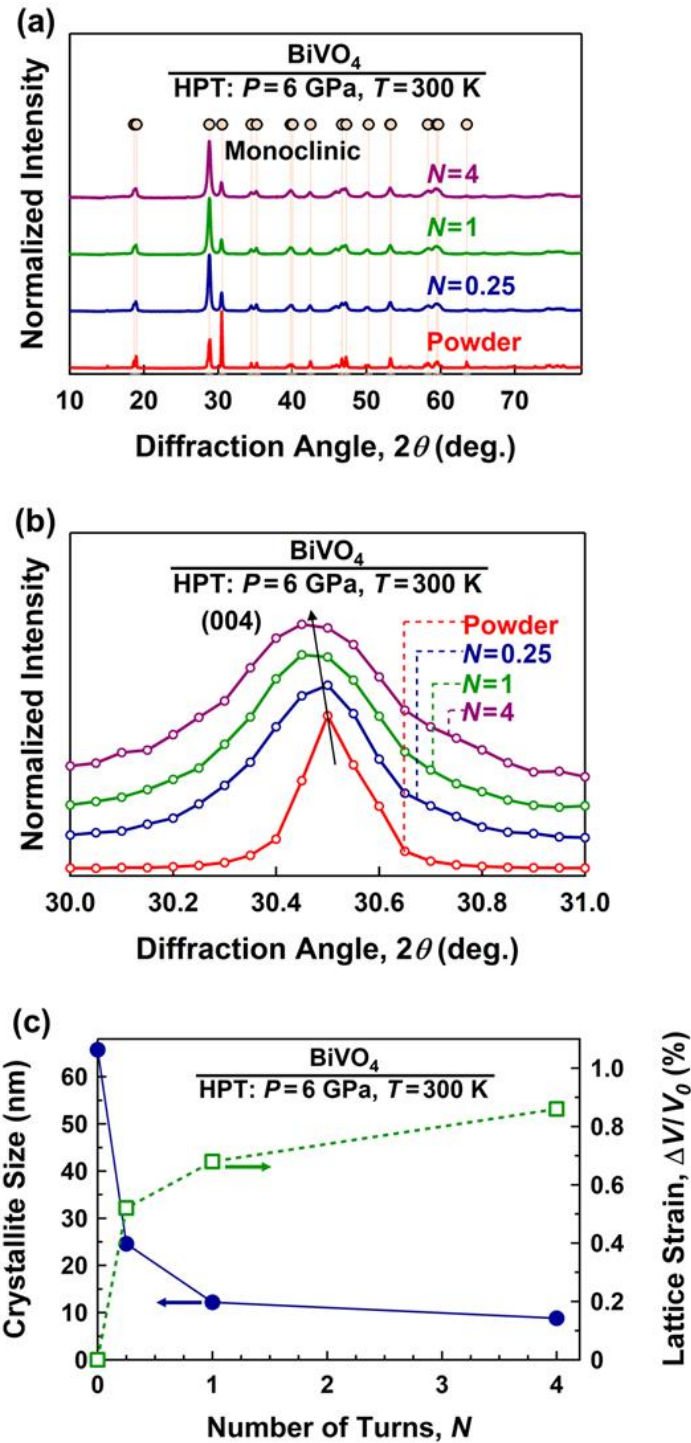


Fig. 2.3. Formation of strained nanocrystals in monoclinic phase BiVO_4 by HPT processing. (a) XRD analysis, (b) XRD analysis in high magnification for (004) atomic plane and (c) crystallite size and volumetric lattice strain versus number of HPT turns for powder and samples processed for 0.25, 1 and 4 turns.

Fig. 2.4a illustrates the crystal structure analysis using the Raman spectroscopy method for the four samples. All peaks in the Raman spectra correspond to the monoclinic phase in good agreement with the XRD analyses. The peaks located at 124 and 210 cm^{-1} are relevant to the external mode of BiVO_4 ; the peaks located at 330 and 365 cm^{-1} are relevant to asymmetric ($\delta_{\text{as}}(\text{VO}_4^{3-})$) and symmetric ($\delta_{\text{s}}(\text{VO}_4^{3-})$) deformation modes of the VO_4^{3-} , respectively; and the peak located at 827 cm^{-1} is due to stretching vibration mode of the V-O bonds in the VO_4^{3-} (ν_{s}) [31]. Since the peak at 827 cm^{-1} shows the length of V-O bonds, it is usually used to analyze the lattice strain in BiVO_4 [32].

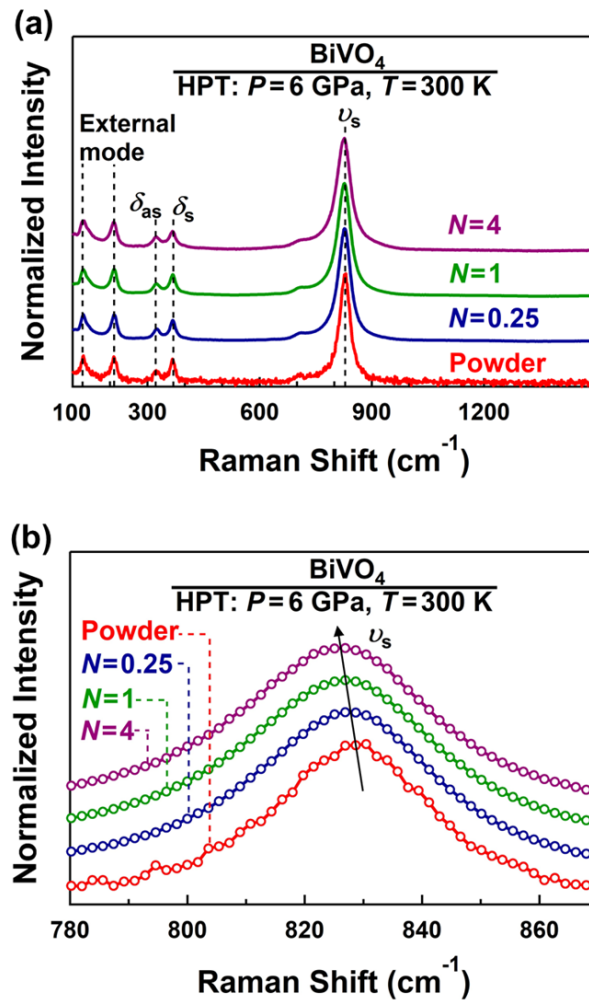


Fig. 2.4. Introduction of lattice strain in BiVO_4 by HPT processing. (a) Raman spectroscopy, and (b) Raman spectroscopy in high magnification for stretching vibration mode of V-O bonds for initial powder and samples processed for 0.25, 1 and 4 turns.

Fig. 2.4b shows the Raman spectra of stretching vibration mode of the V-O bonds in higher magnification. It is evident that there is a peak shift to the lower wavenumbers after HPT processing and the shift is enhanced by increasing the number of turns. This shift indicates that a lattice expansion (tensile strain) occurs in the samples [33].

The distribution of particle size of the four samples measured by DLS is shown in **Fig. 2.5**. The average particles size for the initial powder is 370 nm, but it increases to 420, 540 and 540 nm after HPT processing for 0.25, 1 and 4 turns, respectively. It is concluded that despite decreasing the crystallite size by HPT processing, the specific surface area decreases after HPT processing: $3.8 \text{ m}^2\text{g}^{-1}$ for the initial powder and 3.4, 2.6 and $2.6 \text{ m}^2\text{g}^{-1}$ for the samples processed by HPT for 0.25, 1 and 4 turns, respectively. Such a decrease in specific surface area, which is a consequence of high applied pressure and large plastic deformation, was reported in a wide range of HPT-processed materials [20,21].

Fig. 2.6 shows the microstructure analysis by TEM for four samples in the BF, DF and SAED modes. It should be noted the dark and bright colors in the BF and DF images are due to the diffraction contrasts. The white regions in the DF images correspond to grains having the diffracted beam indicated by arrows in the SAED profiles. The SAED profile has a dotted pattern, and only one grain with the submicrometer size is observed in the initial powder. After HPT processing for 0.25 turns, a combination of dotted and ring patterns is observed in the SAED profiles, indicating grain refinement to the nanometer sizes partially occurs, as is also evident in the DF images with some nanosized white regions. The ring pattern of SAED analyses after 1 and 4 turns and the presence of many white nanosized regions in DF images confirm that the initial submicrometer grain sizes are refined to the nanometer sizes. The average grain size after 4 HPT turns is 15 nm, which is in reasonable agreement with the crystallite size value measured by the XRD analysis in **Fig. 2.3c**. Such a decrease in the crystal size is a consequence of large plastic deformation, as reported in a wide range of HPT-processed oxides [21,22,26].

The nanostructural evolution is shown more clearly in **Fig. 2.7** using the TEM high-resolution images for (a) the initial powder, and for the samples processed by HPT for (b) 0.25, (c, d) 1 and (e, f) 4 turns. The high-resolution images confirm the presence of monoclinic lattice structure in all samples in good agreement with the XRD analysis and Raman spectroscopy. Moreover, while the initial powder contains large crystals, numerous nanograins are formed after HPT processing. In addition to the formation of nanograins with high-angle grain boundaries,

dislocations are also detected within some crystals, as shown in Fig. 2.7d. The presence of large fractions of grain boundaries and dislocations and resultant lattice distortion is consistent with the XRD peak broadening. The evolution of nanostructure to a defective and distorted state is supposed to contribute to an easier electron and hole separation in this material [16].

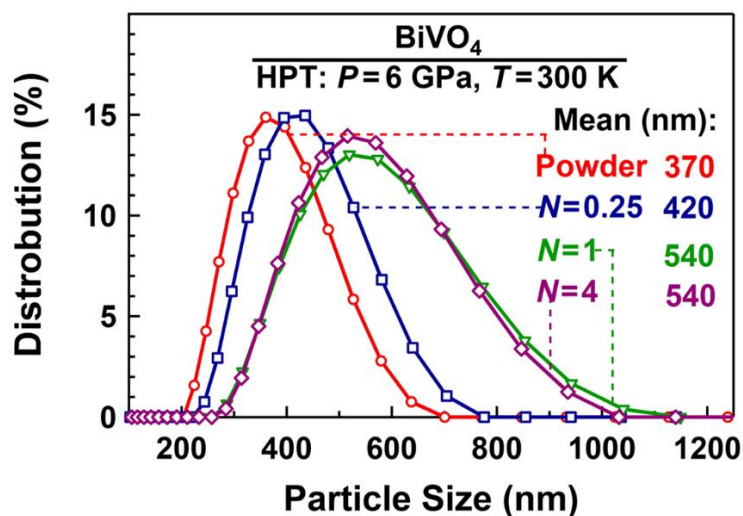


Fig. 2.5. Increasing average particle size of BiVO₄ by HPT processing. Particle size distribution, determined using DLS method, for initial powder and samples processed for 0.25, 1 and 4 turns.

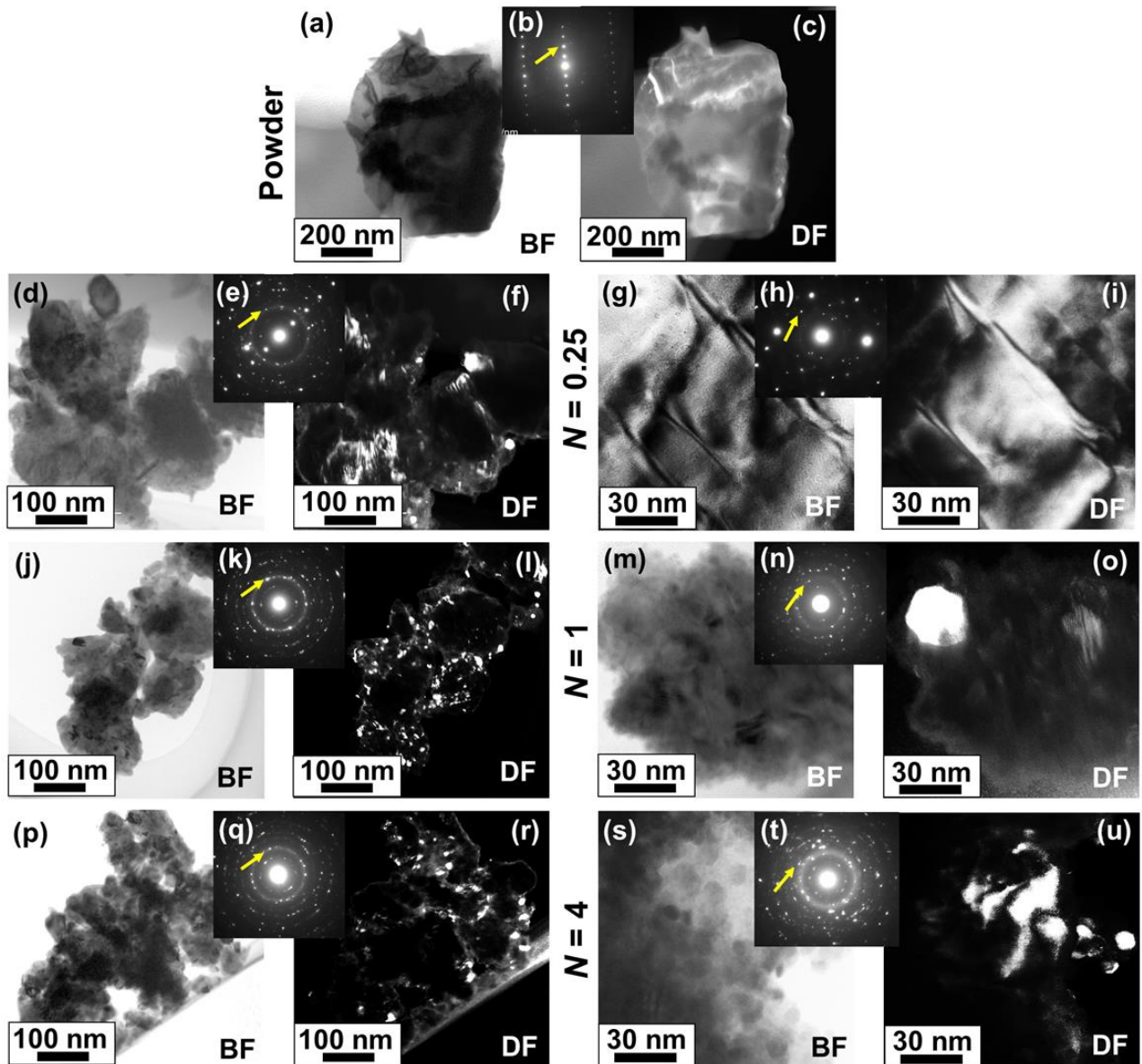


Fig. 2.6. Formation of nanograins in BiVO_4 by HPT processing. TEM BF images, SAED analyses and DF images taken with diffracted beams indicated by arrows in SAED analyses for (a-c) initial powder and samples processed for (d-i) 0.25, (j-o) 1 and (p-u) 4 turns.

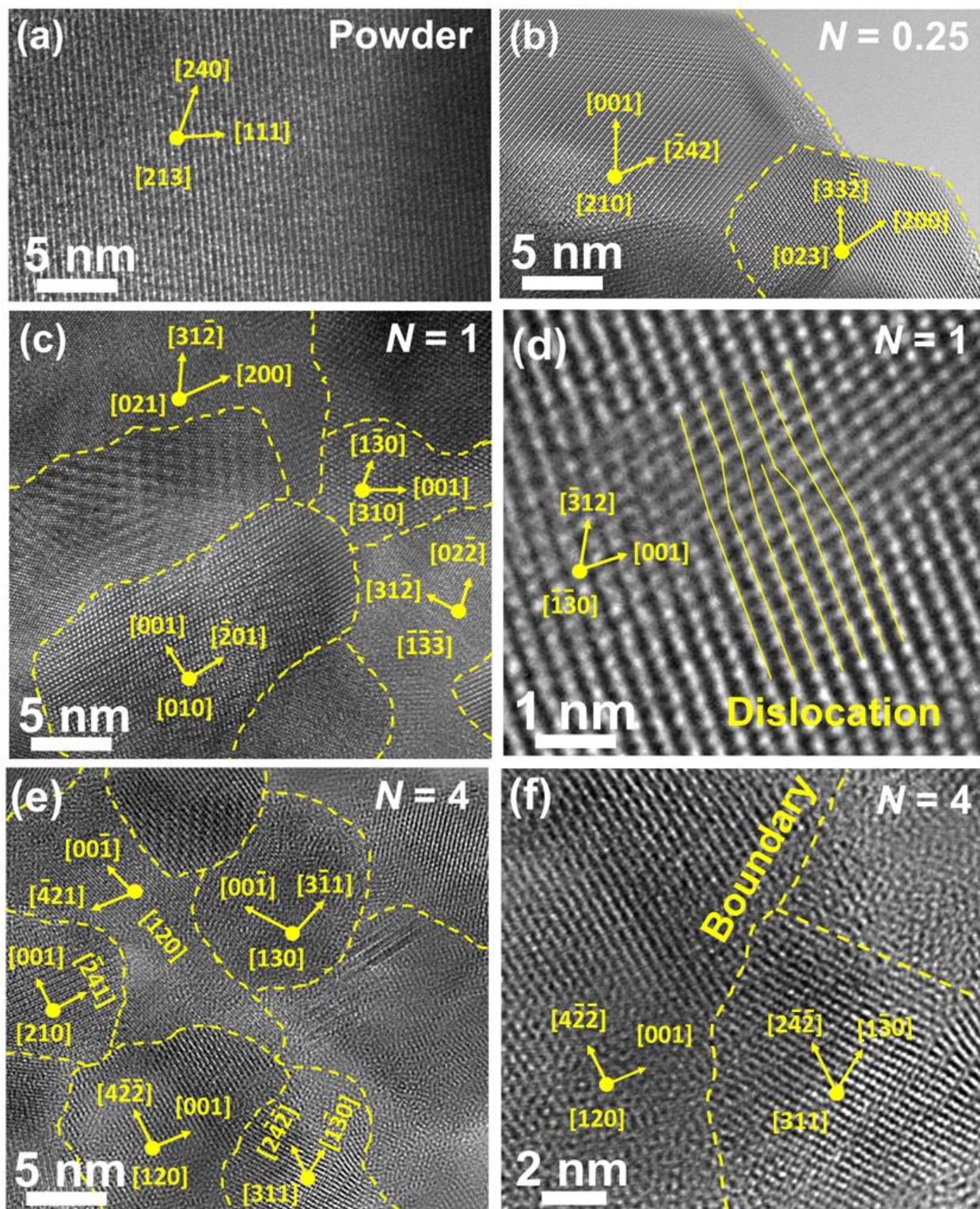


Fig. 2.7. Coexistence of grain boundaries and dislocations in BiVO_4 after HPT processing. TEM high-resolution images for (a) initial powder and for samples processed for (b) 0.25, (c, d) 1 and (e, f) 4 turns.

2.3.2. Oxygen vacancy formation

The changes of the color of samples from yellow for the initial powder to orange for the samples processed with 0.25 and 1 turns and to rose for the sample processed with 4 turns (Fig. 2.1c) indicates that the color centers such as oxygen vacancies should have formed after HPT processing [26]. The lattice expansion, confirmed by XRD peak shift in Fig. 3b and by Raman spectroscopy in Fig. 2.4b, also indirectly suggests that vacancy-type defects can be formed by HPT processing [28,29,32]. In addition to these evidences, the presence of oxygen vacancies can be analyzed by XPS and EPR analyses.

The XPS spectra of initial powder and samples processed by HPT are shown in Fig. 2.8 for (a) Bi 4f, (b) V 2p and (c) O 1s. Examination of XPS data confirms that the main oxidation states of three elements are Bi^{3+} , V^{5+} and O^{2-} . The Bi 4f and V 2p spectra do not show a clear change after HPT processing. However, the O 1s spectra in Fig. 2.8c after HPT processing show small shoulders at high energies, indicating the formation of oxygen vacancies [25]. To quantify, the concentration of oxygen vacancies using XPS, peak deconvolution for the O 1s spectra was conducted by considering two components in the spectra: (i) oxygen at lattice sites O_L , and (ii) oxygen vacancies O_V . The peak positions for O_L and O_V in Fig. 2.8c were considered at 529.9 and 530.7 eV, respectively [34,35]. It is evident that the intensity of oxygen vacancy peak systematically increases with increasing the number of HPT turns. The intensity ratio for O_V to O_V+O_L , as an indication for surface oxygen vacancy concentration, is plotted in Fig. 2.8d for the four samples. The concentration of surface oxygen vacancy reaches ~15% after HPT processing for 4 turns, although this value may be overestimated due to the influence of adsorbed oxygen on O 1s spectra [35].

The EPR spectra for the four samples are shown in Fig. 2.8e, in which the electron spin resonance intensity is plotted against the dimensionless magnetic moment or g factor ($g = h\nu/\mu_B B_0$, h : Planck constant, ν : frequency, μ_B : magnetic field; B_0 : Bohr magneton constant [36,37]). For all samples, peak pairs with opposite curvatures are observed, while the intensity of these peaks increases with increasing the number of HPT turns. For the sample processed for 4 turns, the turning points of these peaks are close to a g factor of 2. Since it was reported that a g factor of 1.978 for BiVO_4 corresponds to oxygen vacancies [39], it is concluded that some amounts of oxygen vacancies are present after HPT processing and particularly after 4 turns. The formation of vacancies by HPT processing, which was reported in both metallic [40,41] and non-metallic

materials [22,23], is due to the plastic deformation effects on vacancy formation and due to the high-pressure effects on suppressing the vacancy annihilation.

2.3.3. Electronic structure

UV-vis spectroscopy and UPS analysis were employed to examine the light absorbance, bandgap and electronic band structure of BiVO₄. Fig. 2.9a shows the light absorbance of the four samples. All samples mainly absorb the light in the UV region, but the light absorbance in the visible region improves after HPT processing and particularly after 4 turns, indicating a better electron-hole separation after HPT processing. Light absorbance edge tends to transfer from the blue light region to the green light region after HPT processing for 4 turns, which is in agreement with the prediction of the Itten color wheel in Fig. 2.1c. The Kubelka-Munk analysis is shown in Fig. 2.9b for estimation of the direct bandgap. A bandgap narrowing occurs by HPT processing and the bandgap decreases from 2.4 eV for the initial powder to 2.3 eV for the samples processed with 0.25 and 1 turns and to 2.1 eV for the sample processed with 4 turns. Since no phase transformations occur in BiVO₄ by HPT processing, such a bandgap narrowing should be due to the introduction of lattice defects and strain [42-45,18].

Fig. 2.9c shows the UPS spectra of four samples which are used to determine the top of the valence band from the vacuum level. The top of the valence band from the vacuum level was calculated as cutoff energy from the Fermi level (E_0 indicated by arrows in Fig 2.9c) - valence band top energy from the Fermi level (E_B indicated by arrows in Fig. 2.9c) - UV He-I energy (21.2 eV) [46]. The top of the valence band for the initial powder and the samples processed by HPT for 0.25, 1 and 4 turns is -6.4, -6.0, -6.0 and -6.0 eV versus the vacuum level, respectively. By the addition of bandgap values to the top of the valence band, the bottom of the conduction band can be determined as -4.7, -3.7, -3.7 and -3.9 eV for the initial powder and the samples processed by HPT for 0.25, 1 and 4, respectively.

Fig. 2.9d shows the band structure of four samples, while the numbers in the figure are the positions of the top of the valence band and the bottom of the conduction band versus the normal hydrogen electrode (NHE): energy vs. NHE = -4.4 - energy vs. vacuum [46]. For the initial powder, the bandgap, the top of the valence band and the bottom of the conduction band are in agreement with the reported data in the literature [47].

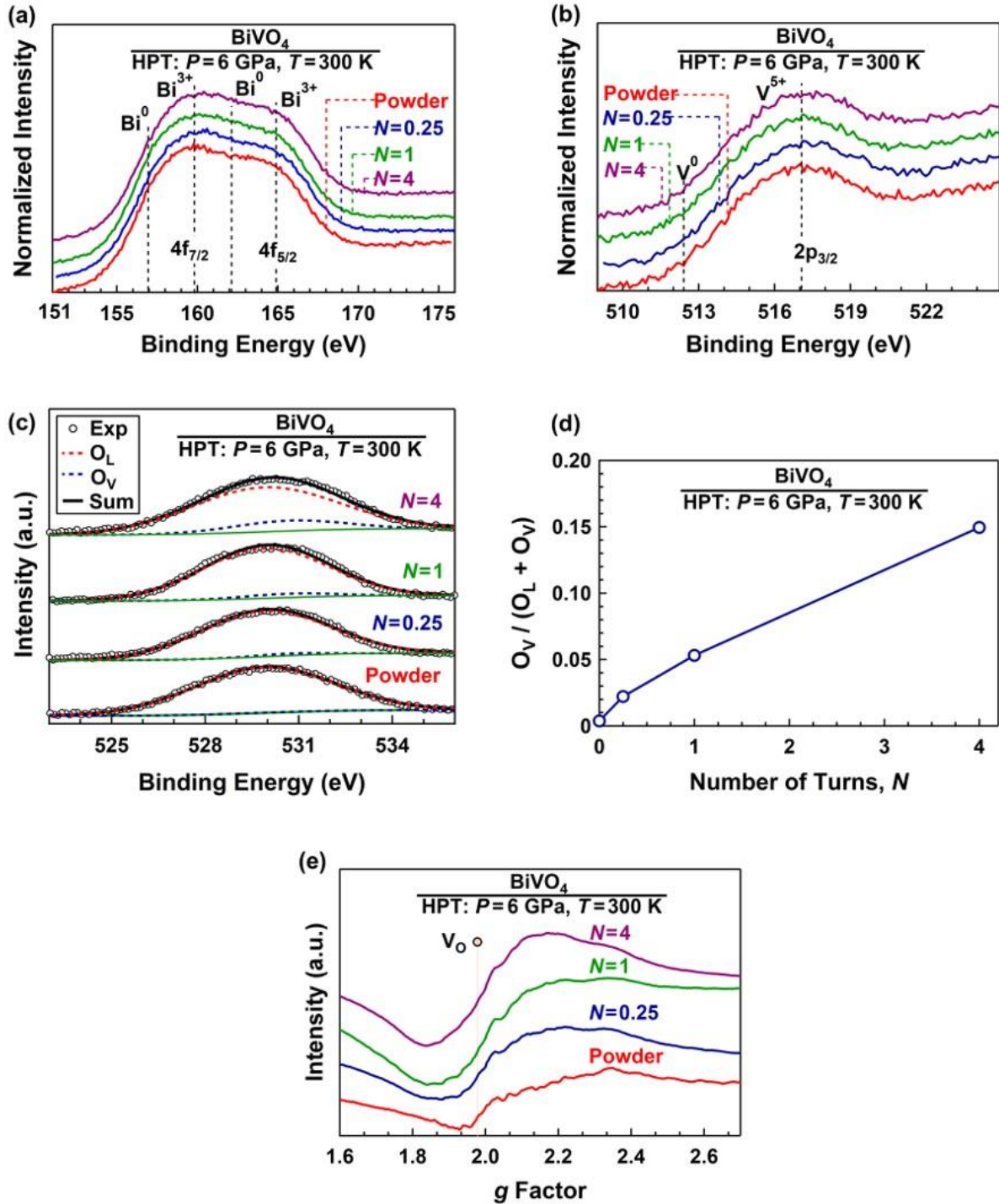


Fig. 2.8. Formation of oxygen vacancies in BiVO₄ after HPT processing. XPS spectroscopy of (a) Bi 4f, (b) V 2p and (c) O 1s and its peak deconvolution; (d) ratio of surface oxygen vacancy peak (O_V) to sum of lattice oxygen and vacancy peaks (O_L + O_V) plotted versus number of HPT turns; and (e) EPR spectroscopy for initial powder and samples processed for 0.25, 1 and 4 turns.

After HPT processing, in addition to bandgap narrowing the bottom of the conduction band increases, suggesting that the problem of low conduction band bottom of BiVO₄ [2,3] can be successfully solved by the introduction of lattice defects and strain. As shown in Fig. 2.9d, the electronic structure of all four samples can support the CO₂ conversion reactions by considering the relevant chemical potentials for different reactions [48,49]. However, the HPT-processed samples have a higher overpotential for CO₂ conversion, which is thermodynamically desirable for photocatalysis.

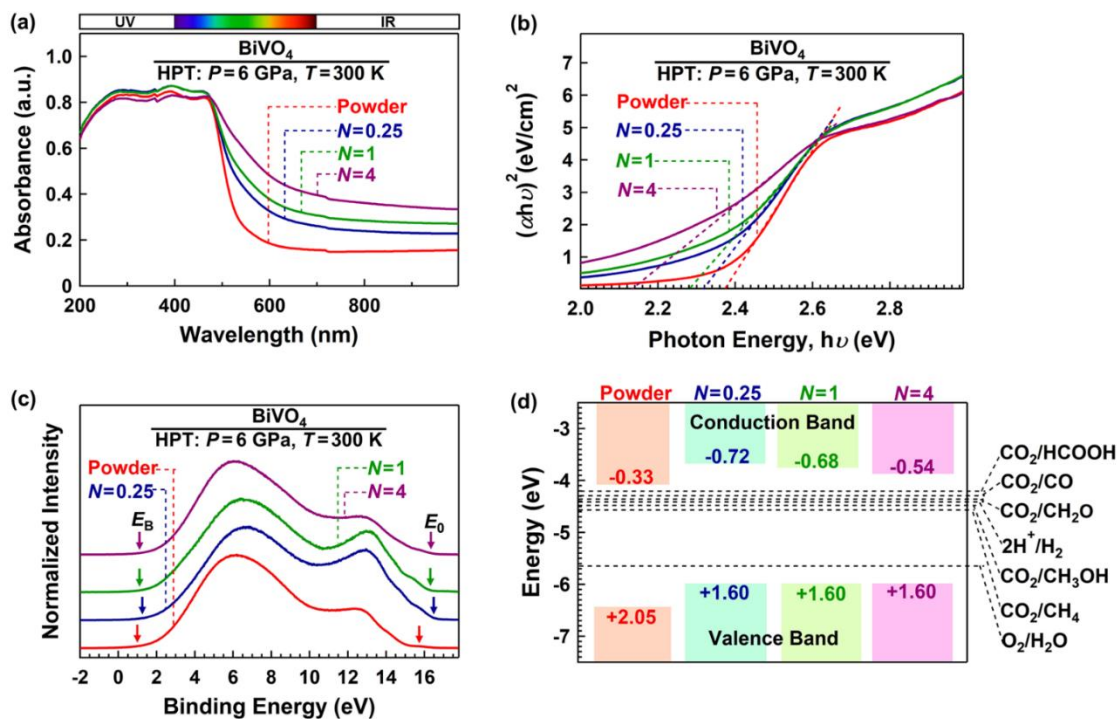


Fig. 2.9. Bandgap narrowing and improvement of electronic band structure of BiVO₄ for photocatalytic CO₂ conversion by HPT processing (a) UV-vis spectroscopy, (b) Kubelka-Munk analysis for direct bandgap calculation (α : light absorption, h : Planck's constant, ν : photon frequency), (c) UPS spectroscopy using bias of -4 V to measure the top of the valence band, and (d) electronic band structure including chemical potentials for CO₂ reduction reactions for initial powder and samples processed for 0.25, 1 and 4 turns. Arrows on left and right in (c) indicate the valence band top energy (E_B) and cutoff energy (E_0) shifted 4 eV lower from Fermi level, respectively. Numbers in (d) shows energy levels versus normal hydrogen electrode (NHE).

While earlier studies used the heterojunctions to solve the problem of the electronic structure of BiVO₄ [3-8], this study confirms that engineering the lattice defects and strain is effective to improve the electronic structure of pure BiVO₄ without any need for a second phase or impurity atoms.

2.3.4. Charge carrier dynamics

Fig. 2.10 shows (a) PL spectroscopy, (b) PL decay analysis and (c) photocurrent measurements which can give information about charge carrier dynamics including the recombination of electrons and holes, mobility of charge carriers and defect trapping [49]. Fig. 2.10a shows a significant decrease in PL intensity after HPT processing, indicating that the recombination rate of electrons and holes decreases effectively after HPT processing. These results suggest that the HPT-induced defects and strain act positively as electron-hole separation sites rather than recombination sites [4,11,26]. The PL peak for all samples appeared at a peak around 650 nm (1.9 eV) which is less than the bandgap of samples. Although electron-hole recombination in BiVO₄ can happen radiatively or nonradiatively, the emission detected around 650 nm in Fig. 2.10a was attributed to the radiative recombination of electrons and holes on defects [39,48]. PL decay spectra shown in Fig. 2.10b follow a relation shown below [49].

$$I(t) = A_1 \exp\left(-\frac{t}{\tau_1}\right) + A_2 \exp\left(-\frac{t}{\tau_2}\right) \quad (2.2)$$

where, $I(t)$, A_1 , A_2 , τ_1 and τ_2 are PL decay intensity at time t , the amplitude of the first exponential function, amplitude of the second exponential function, fast decay time and slow decay time, respectively. The average lifetime can be calculated using the following equation [49].

$$\tau_{ave} = \frac{A_1 \tau_1^2 + A_2 \tau_2^2}{A_1 \tau_1 + A_2 \tau_2} \quad (2.3)$$

Values of A_1 , A_2 , τ_1 , τ_2 τ_{ave} for the initial powder and the samples processed by HPT are presented in Table 2.2. The average electron lifetime decreases from 12.90 ns for the initial powder to 11.48, 10.22 and 9.24 ns for the samples processed for 0.25, 1 and 4 turns, respectively. The decrease in the lifetime for HPT samples can be attributed to the formation of oxygen vacancies. Oxygen vacancies on the surface of photocatalyst act as shallow traps or active sites for the fast migration of electrons and participation in a reaction. Therefore, their formation can lead to a decrease in the excited electron lifetime and an increase in the photocatalytic activity [26,35]. It

should be noted that the oxygen vacancies in the bulk can also lead to electron migration and a decrease in the lifetime, but they do not necessarily improve the photocatalytic activity [16,50,51].

Photocurrent measurements in Fig. 2.10c demonstrate that the current density is the highest for the initial powder and decreases with HPT processing and increasing the number of turns. Since photocurrent generation is attributed to the migration of electrons from the surface under an external voltage, Fig. 2.10c suggests that electron separation from the surface of BiVO₄ becomes weaker after HPT processing. Such a decrease in photocurrent may not be due to enhanced electron-hole recombination because the PL intensity (as an indication of recombination) decreases after HPT processing, as shown in Fig. 2.10a. The decreases in photocurrent density after HPT processing should be due to the action of surface vacancies as electron trapping sites. It was shown earlier that polarity and localization of photoexcited electrons of BiVO₄ can lead to electron trapping in this material [52]. Such trapping can be enhanced on oxygen vacancies, making these defects active sites with sufficient electrons for chemical reactions (particularly if they form near the bottom of the conduction band and act as shallow donors) [53]. Theoretical studies suggested that a change of the oxygen vacancy states from the shallow donors to localized intragap states can reduce the photoconductivity and photocurrent of BiVO₄ [53], but this is not the case in this study because the PL spectra in Fig. 2.10a show similar peak energies for all four samples.

Table 2.2. Decreasing electron lifetime in BiVO₄ by HPT processing. PL decay parameters and electron lifetimes for initial powder and for samples processed for 0.25, 1 and 4 turns.

Sample	τ_1 (ns)	τ_2 (ns)	A_1	A_2	τ_{ave} (ns)
Powder	2.17	14.92	56.56	43.44	12.90
$N = 0.25$	1.92	13.08	53.19	46.81	11.48
$N = 1$	1.78	12.40	64.26	35.74	10.22
$N = 4$	1.62	10.35	48.09	51.91	9.24

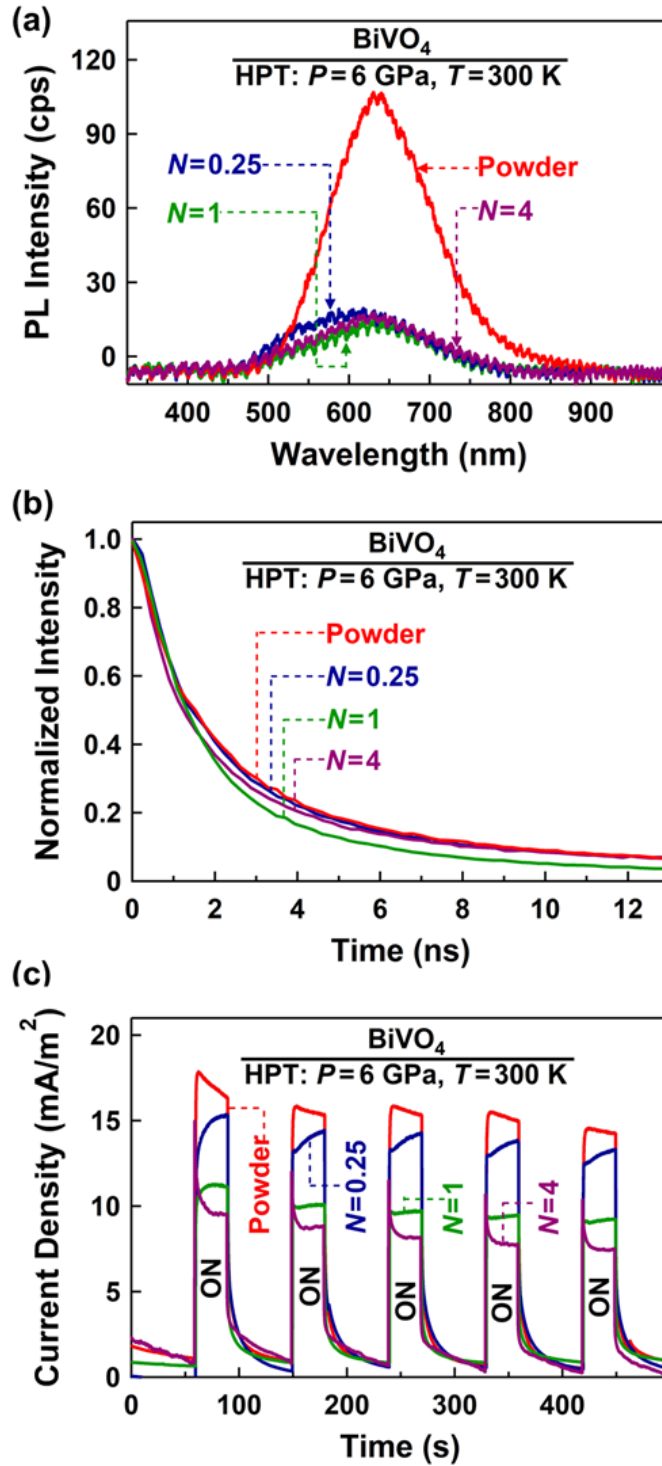


Fig. 2.10. Decreasing recombination of electrons and holes in BiVO₄ after HPT processing. (a) Steady-state PL emission spectra, (b) time-resolved PL decay spectra, and (c) photocurrent density versus time for initial powder and samples processed for 0.25, 1 and 4 turns.

2.3.5. Photocatalytic activity

Photocatalytic CO₂ to CO conversion activity of the initial powder and the samples processed by HPT for 0.25, 1 and 4 turns are shown in Fig. 2.11a versus the catalyst mass unit and in Fig. 2.11b versus the catalyst surface area unit. Since the photocatalytic reactions occur on the surface of photocatalysts, normalizing the CO production rates per surface area is more reasonable to show the differences. It should be noted that the photocatalytic tests were conducted two times in Figs. 2.11a and 2.11b for the initial powder and for the sample processed with $N = 4$ to confirm the reproducibility of data. All samples could convert CO₂ to CO, but the photocatalytic performance of HPT-processed samples is better than the initial powder. The sample processed for 4 turns shows the highest activity. The average of CO production rate for the powder is $2.27 \pm 0.24 \mu\text{molh}^{-1}\text{g}^{-1}$ ($0.60 \pm 0.06 \mu\text{molh}^{-1}\text{m}^{-2}$) which increases to 3.21 ± 0.23 , 3.35 ± 0.27 and $3.97 \pm 0.38 \mu\text{molh}^{-1}\text{g}^{-1}$ (0.96 ± 0.07 , 1.27 ± 0.10 and $1.45 \pm 0.08 \mu\text{molh}^{-1}\text{m}^{-2}$) for the samples proceeded for 0.25, 1 and 4 turns, respectively. Here, it should be noted that no hydrocarbons such as CH₄ and no oxidation products such as O₂ were detected within the detection limits of the experiments. The absence of the oxidation products should be a reason for the gradual decrease in the activity during the time. It should be mentioned that the result of the blank test before the photocatalytic activity test (addition of catalysts without light irradiation) shows that the production rate of CO is zero (the data point at $t = 0$ h). Moreover, a blank test without the addition of catalyst with light irradiation shows no CO production within 8 h, as shown in Fig. 2.11a and 2.11b. The improvement of photocatalytic activity of BiVO₄ after HPT processing should be due to the presence of oxygen vacancies as active sites for adsorption and conversion of CO₂ molecules as well as due to the effect of lattice strain on the improvement of band structure for CO₂ photoreduction [2,54].

Fig. 2.11c demonstrates the structure stability of powder and samples processed by HPT after photocatalytic test using the XRD analysis. All samples are stable and retain their monoclinic structure after use for photocatalytic CO₂ conversion. No chemical reaction, corrosion or degradation products are detected by XRD analysis. Since some photocatalysts are negatively influenced by photocorrosion phenomenon, finding the photocorrosion-resistant catalysts is a critical issue [55,56]. In some BiVO₄-based catalysts, it was reported that the photocorrosion and formation of bismuth oxides can occur by dissolving vanadium into the solution through the

accumulation of holes on the surface [55,56]. The absence of such photocorrosion products confirms the stability of current HPT-processed photocatalysts.

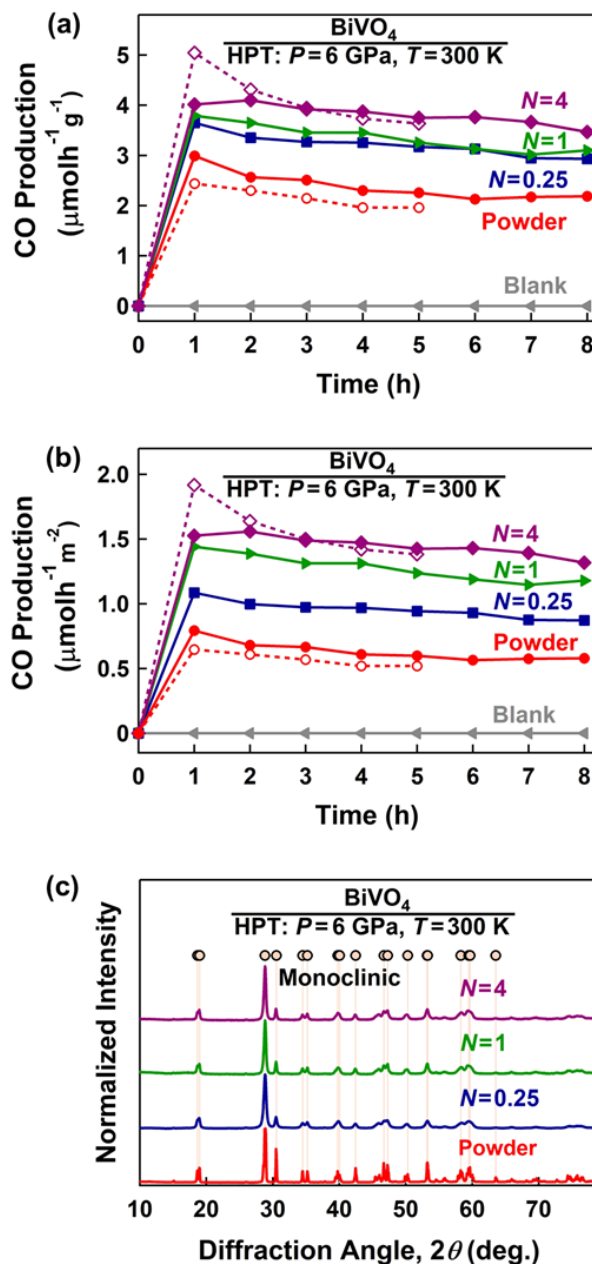


Fig. 2.11. Improvement of photocatalytic CO₂ conversion on highly-stable BiVO₄ after HPT processing. CO production rate under UV irradiation versus time normalized by (a) mass unit and (b) surface area of catalysts; and (c) XRD profiles after photocatalytic test for initial powder and samples processed for 0.25, 1 and 4 turns. Photocatalytic tests were repeated two times for initial powder and sample processed for 4 turns.

2.4. Discussion

This chapter shows that the simultaneous introduction of strain and oxygen vacancies into BiVO₄ can lead to improving the light absorbance, narrowing the bandgap, modifying the electronic band structure, decreasing the recombination rate of electrons and holes and finally enhancing the photocatalytic activity for CO₂ conversion. Unlike earlier strategies that mainly used heterojunctions [3-8] or doping [9-13] to deal with the low level of the conduction band in BiVO₄ [2,3], this study shows that the level can be enhanced even in pure BiVO₄ by the accumulation of oxygen vacancies and lattice strain. Due to the presence of unpaired electrons in surface oxygen vacancies, they can act as adsorption sites to uptake CO₂ as a Lewis acid [43,57]. These vacancies also contribute to the activation of adsorbed CO₂ and H₂O for conversion to intermediate and final products by weakening the carbon-oxygen and oxygen-hydrogen bondings [43]. The oxygen vacancies on the surface do not usually act as recombination sites for electrons and holes, but they trap the excited electrons and provide active sites for photocatalytic reactions [57-59]. Here, it should be noted that the concentration of oxygen vacancies in BiVO₄ as activation sites should be an optimum level because large fractions of oxygen vacancies can reduce V⁵⁺ to V⁴⁺, resulting in electron scattering due to a large ionic diameter of V⁴⁺ [29,60,61]. A high concentration of V⁴⁺ can lead to decreasing the length of effective diffusion of holes and accordingly reduce the photocatalytic activity [62]. Despite these expected effects of oxygen vacancies, the overall effect of oxygen vacancies in this study is the improvement of photocatalytic activity. On the other hand, lattice strain can lead to electronic band structure modification such as changing the level of the conduction band bottom or decreasing the bandgap [18,19,45]. Therefore, in the current HPT-processed BiVO₄ photocatalysts, increasing the light absorbance, decreasing the bandgap, and increasing the bottom of the conduction band energy are the results of the contribution of lattice strain, while decreasing the electron-hole recombination and improvement of surface catalytic activity are the main contribution of oxygen vacancies.

Here, to have a clear view regarding the improvement of photocatalytic CO₂ conversion of highly-strained and defective BiVO₄, some comparisons can be made. Fig. 2.12 compares the CO production rate of this material with P25 TiO₂, which is a typical benchmark photocatalyst. Note that both materials were examined exactly under similar conditions by the current authors. The average of photocatalytic CO production of P25 TiO₂ is $4.63 \pm 0.33 \mu\text{molh}^{-1}\text{g}^{-1}$ which is close to the productivity of HPT-processed BiVO₄ after 4 turns ($3.97 \pm 0.38 \mu\text{molh}^{-1}\text{g}^{-1}$). Here, it should

be noted that the surface area of P25 TiO₂ is 38.70 m²g⁻¹ which is much higher than the surface area of HPT-processed BiVO₄ (2.6 m²g⁻¹). Despite the negative effect of the HPT method on the reduction of the surface area [21,28], Fig. 2.12 confirms the efficiency of the defect and strain engineering in improving the photocatalytic activity to high levels comparable with one of the most active photocatalysts. Another comparison is given in Table 2.3, in which photocatalytic activity of HPT-processed BiVO₄ was compared with some data for BiVO₄-based photocatalysts [4-13]. It should be noted that data in Table 2.3 should be compared with care because of the differences in experimental conditions in various studies (e.g., light source, photoreactor type, the mass of catalyst, etc.). Despite these differences, the range of reported photocatalytic CO production rate for these materials is between 0.05-13.5 μmolh⁻¹g⁻¹, while the HPT-processed BiVO₄ shows one of the highest activities.

Finally, it is worth mentioning that the HPT method used in this study for the generation of defects and strain is applicable to any kind of conventional oxides [20,21] and multi-component ceramics and salts [63]. Although the method was used earlier by the authors to improve the photocatalytic CO₂ conversion by synthesis of high-pressure phases [27] or high-entropy oxides [28], its current application to introduce strain and defects is applicable to a wider range of materials to develop advanced photocatalysts for CO₂ conversion. Since the specific surface area is usually small after HPT processing, other alternative mechanical or chemical routes may be employed in the future for simultaneous generation of strain and defects in photocatalysts, while keeping their surface area large.

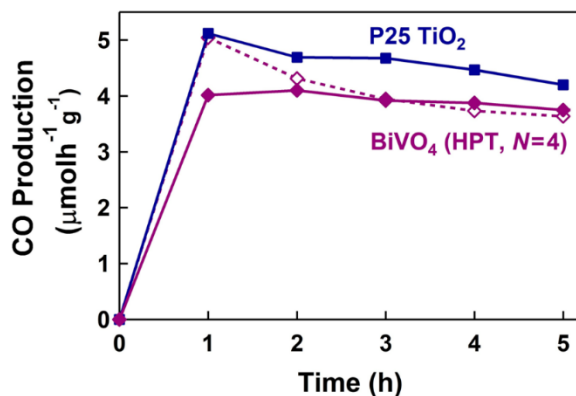


Fig. 2.12. High photocatalytic CO₂ conversion activity of highly-strained and oxygen-deficient BiVO₄ in comparison with P25 TiO₂ as benchmark photocatalyst. CO production rate versus irradiation time for BiVO₄ processed by HPT for 4 turns and P25 TiO₂.

Table 2.3. Photocatalytic CO evolution rate of BiVO₄ processed by HPT for 4 turns compared with other BiVO₄-based photocatalysts.

Photocatalyst	Catalyst Mass (mg)	Light Source	CO Production ($\mu\text{molh}^{-1}\text{g}^{-1}$)	Reference
BiVO ₄ {110}-Au- Cu ₂ O	100	300 W Xe (> 420nm)	1.12	[4]
BiVO ₄ {010}-Au- Cu ₂ O	100	300 W Xe (> 420nm)	2.02	[4]
BiVO ₄ {010}-Cu ₂ O	100	300 W Xe (> 420nm)	0.45	[4]
BiVO ₄ {010}-Au-(Cu ₂ O-Au)	100	300 W Xe (> 420nm)	2.08	[4]
BiVO ₄ -Cu ₂ O	-----	300 W Xe (> 420nm)	~0.5	[5]
BiVO ₄ -C- Cu ₂ O	-----	300 W Xe (> 420nm)	3.01	[5]
ZnIn ₂ S ₄ -BiVO ₄	100	300 W Xe	4.75	[6]
BiVO ₄ -5 % Bi ₄ Ti ₃ O ₁₂	10	300 W Xe	~6.75	[7]
BiVO ₄ -10 % Bi ₄ Ti ₃ O ₁₂	10	300 W Xe	~13.5	[7]
BiVO ₄ -20 % Bi ₄ Ti ₃ O ₁₂	10	300 W Xe	~6.75	[7]
Bi-BiVO ₄	10	300 W Xe (> 420nm)	~0.75	[9]
Cu-Bi-BiVO ₄	10	300 W Xe (> 420nm)	~11	[9]
BiVO ₄	10	300 W Xe (> 420nm)	0.42	[11]
Bi-BiVO ₄	10	300 W Xe (> 420nm)	1.29	[11]
Ag-Bi-BiVO ₄ (Ag/Bi = 0.3)	10	300 W Xe (> 420nm)	3.4	[11]
Ag-Bi-BiVO ₄ (Ag/Bi = 0.6)	10	300 W Xe (> 420nm)	5.19	[11]
Ag-Bi-BiVO ₄ (Ag/Bi = 0.9)	10	300 W Xe (> 420nm)	4.54	[11]
1.0%Cu-BiVO ₄	50	400 W High-Pressure Hg	2	[12]
0.5%Cu-BiVO ₄	50	400 W High-Pressure Hg	4.1	[12]
0.3%Cu-BiVO ₄	50	400 W High-Pressure Hg	3.9	[12]
Au-BiVO ₄	50	300 W Xe (> 420nm)	~0.05	[8]
CdS-BiVO ₄	50	300 W Xe (> 420nm)	~0.75	[8]
Cd _{0.5} Zn _{0.5} S-BiVO ₄	50	300 W Xe (> 420nm)	~1.2	[8]
CdS-Au-BiVO ₄	50	300 W Xe (> 420nm)	~1.1	[8]
Cd _{0.2} Zn _{0.8} S-Au-BiVO ₄	50	300 W Xe (> 420nm)	~1.6	[8]
Cd _{0.5} Zn _{0.5} S-Au-BiVO ₄	50	300 W Xe (> 420nm)	~2.15	[8]
Cd _{0.8} Zn _{0.2} S-Au-BiVO ₄	50	300 W Xe (> 420nm)	~1.3	[8]
BiVO ₄ -2%Co	30	25 W 254 nm Ultraviolet	0.62	[13]
BiVO ₄ -5%Co	30	25 W 254 nm Ultraviolet	0.68	[13]
BiVO ₄ -10%Co	30	25 W 254 nm Ultraviolet	0.83	[13]
BiVO ₄ -15%Co	30	25 W 254 nm Ultraviolet	1.04	[13]
BiVO ₄ -20%Co	30	25 W 254 nm Ultraviolet	2.08	[13]
BiVO ₄ (HPT: N = 4)	120	400 W High-Pressure Hg	3.97 ± 0.38	This study

2.5. Conclusion

A photocatalyst with low bandgap, improved band structure and low recombination rate of electrons and holes was produced by simultaneous introduction of oxygen vacancies and lattice strain in BiVO₄ via a high-pressure torsion process. The material showed high photocatalytic activity for CO₂ to CO conversion with an activity comparable to P25 TiO₂ as a benchmark photocatalyst. This study suggests that simultaneous engineering of lattice strain and defects is an effective strategy to produce active photocatalysts for CO₂ conversion.

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Chapter 3. High-pressure TiO₂-II polymorph as an active photocatalyst for CO₂ to CO conversion

3.1. Introduction

TiO₂ has three main polymorphs under ambient conditions (rutile and anatase with the tetragonal crystal structure and brookite with the orthorhombic crystal structure) [1], and among them, the anatase phase shows the highest photocatalytic activity due to its high electron effective mass and large specific surface area [2]. However, the large bandgap of anatase (3.0-3.2 eV) and its limited activity (due to short lifetime of charge carriers) have led to numerous studies to narrow its optical bandgap and improve its photocatalytic activity. The main focus of these studies is the addition of a third element or dopants to TiO₂ such as nitrogen [3], sulfur [4], iodine [5], fluorine [6], nickel [7], chromium [7], vanadium [7], cobalt [7], copper [8], silver [9] and gold [10]. Although addition of a third element to TiO₂ is an effective technique to reduce the bandgap, the method may not enhance the photocatalytic activity significantly due to the impurity-induced recombination effects [11]. Therefore, there are new trends to improve the photocatalytic CO₂ conversion on pure TiO₂ by nanostructure and polymorphic control and without addition of impurity elements [12-15].

In this study, to improve the photocatalytic CO₂ conversion on TiO₂ by nanostructure control and without addition of any impurity elements, the high-pressure columbite (TiO₂-II) phase with the orthorhombic crystal structure is introduced in TiO₂ (see the pressure-temperature phase diagram of TiO₂ in Fig. 3.1a [1]). A high-pressure torsion (HPT) method [16], as schematically shown in Fig. 3.1b, is used to stabilize the TiO₂-II phase because earlier studies confirmed the high capability of this method for controlling the polymorphic phase transformation in various oxides [17,18] including those with photocatalytic properties such as TiO₂ [19], ZnO [20] and ZrO₂ [21]. This first application of high-pressure TiO₂-II phase for photocatalytic CO₂ conversion confirms its higher activity compared with the pure anatase phase.

3.2. Experimental section

3.2.1. Sample preparation

About 250 mg of TiO₂ anatase powder (Sigma Aldrich) with the purity level of 99.8% (impurities in mass ppm: 0.3 Be, 23.3 Na, 14.7 Mg, 116.6 Al, 1773.1 K, 26.0 Ca, 12.1 V, 0.6 Mn, 19.7 Fe, 2.7 Zn, 0.4 Sr, 144.8 Zr, 20.0 Pd, 18.7 Sn, 5.8 Ba, 13.0 Pt, 10.2 Pb, 22.9 Bi) and 50-250 nm particle sizes was compacted in the form of 10 mm disc under a pressure of 380 MPa. The compacted disc was then pressurized between two HPT anvils under a pressure of 6 GPa at ambient temperature ($T = 300$ K) and concurrently strained by rotating the two anvils with respect to each other for $N = 3$ turns with a rotation speed of 1 rpm (see the principles of HPT in [16,17]). Since the HPT method can generate oxygen vacancies in oxides [18,19], the HPT-processed samples were annealed in the air atmosphere at 773 K for 1 h to minimize the effect of oxygen vacancies on photocatalytic CO₂ conversion. The initial powder, the sample after HPT processing and the sample after annealing were crushed, as shown in Fig. 3.1c, and examined in detail using different techniques, as described below.

3.2.2. Characterization

The specific surface area of the crushed samples was examined by nitrogen gas adsorption using the Brunauer-Emmett-Teller (BET) technique. The estimated surface areas were 10.2, 6.2 and 6.8 m²/g and no porosities were detected within the detection limits of the BET method.

X-ray diffraction (XRD) using the Cu K α radiation with a wavelength of $\lambda = 0.1542$ nm and Raman spectroscopy using a 532-nm laser source were utilized to investigate the crystal structures and phase transformations. The XRD profiles were evaluated by the Rietveld method using the FullProf Suite software and by refining the profiles using the Gaussian and Lorentzian functions (“Thompson-Cox-Hastings pseudo-Voigt” function) [22]. The average lattice strain was determined from the width of Gaussian function $H_G(\Delta 2\theta)$, and the average crystallite size (D) was determined from the width of Lorentzian function $H_L(\Delta 2\theta)$, using the following relationships (θ : Bragg angle) [23,24].

$$D = \frac{2\lambda}{\pi H_L(\Delta 2\theta) \cos\theta} \quad (3.1)$$

$$\varepsilon = \frac{H_G(\Delta 2\theta)}{8\sqrt{\ln 2/\pi} \tan\theta} \quad (3.2)$$

Transmission electron microscopy (TEM) in the bright-field (BF) mode, dark-field (DF) mode, selected area electron diffraction (SAED) and high-resolution imaging combined with fast Fourier transform (FFT) was performed to clarify the micro/nanostructural features of phases. For TEM, the samples were crushed in ethanol and dispersed on carbon grids and examined by an aberration-corrected microscope in an acceleration voltage of 200 kV.

Electron paramagnetic resonance (EPR) was conducted at ambient temperature using a 9.4688 GHz microwave source to study the electron spins and evaluate the presence of oxygen vacancies and formation of Ti^{3+} states.

X-ray photoelectron spectroscopy (XPS) was conducted using a Mg K_{α} source to investigate the electronic states of different elements, examine the formation of oxygen vacancies and estimate the bottom of valence band.

Ultraviolet-visible (UV-vis) diffuse reflectance spectroscopy was performed to investigate the light absorbance of samples and calculate the bandgap using the Kubelka-Munk analysis. A combination of UV-vis spectroscopy and XPS was used to determine the band structure of samples.

The steady-state photoluminescence (PL) emission spectroscopy with a 325-nm laser source and time-resolved photoluminescence decay (PL decay) with a 285-nm laser source were used to investigate the recombination of electrons and holes and decay time of excited electrons.

Attenuated total reflectance - Fourier transform infrared spectroscopy (ATR-FTIR) was conducted at room temperature using a Jasco FT/IR-610 device.

Temperature-programmed desorption (TPD) was conducted using a commercial Micotrac Bel instrument equipped with a quadrupole mass spectrometric detector and thermal conductivity detector. For TPD, 30 mg of samples were first placed in pure (99.99%) CO_2 with a gas flow rate of 100 mL/min under 0.1 MPa at 323 K or 773 K for 2 h. After evacuation with a vacuum pump at 323 K for 15 min, the samples were then heated from 323 K to 773 K with a heating rate of 5 K/min and the TPD signals were recorded.

3.2.3. Photocurrent test

Photocurrent experiment was performed using a Xe lamp in 1 M Na_2SO_4 electrolyte to investigate the photoactivity. In these tests, thin film of TiO_2 sample on FTO (fluorine-doped tin oxide) glass, Pt wire and Ag/AgCl were used as working, counter and reference electrodes, respectively. An electrochemical analyzer was utilized to investigate the photocurrent generation

by potentiostatic amperometry during time (30 s light ON and 60 s light OFF). For making the thin films, 5 mg of samples were crushed in 0.2 mL ethanol and uniformly dispersed on 0.1-mm thick FTO glass with 15×25 mm² surface area and baked in furnace at 473 K for 1 h.

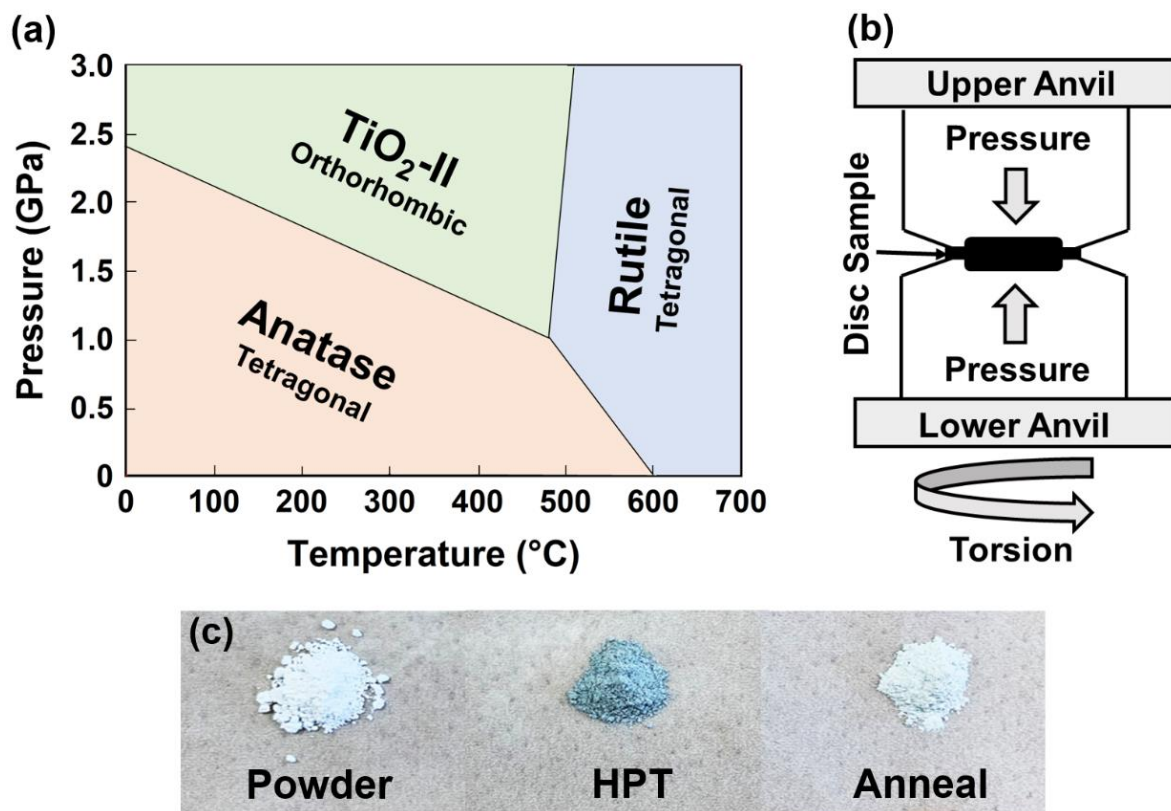


Fig. 3.1. (a) Pressure-temperature phase diagram of TiO₂, (b) schematic of high-pressure torsion (HPT) process and (c) appearance of three samples before HPT, after HPT and after annealing.

3.2.4. Photocatalytic test

For photocatalytic CO₂ conversion, 100 mg of samples were dispersed in 30 mL solution of 1 M NaHCO₃ in a CO₂ flow reactor system, and after bubbling with CO₂ (3mL/min) in dark for 1 h, irradiated with a 300 W Hg-Xe light source (Newport, Arc Lamp Housing 66905, 300 W), as schematically shown in Fig. 3.2a. The photoreactor, as shown in Fig. 3.2a, was made from quartz, and had a cuboid shape with a total volume of 105 mL (30×50×70 mm³ internal dimensions with a quartz thickness of 10 mm). The reactor had three holes on top for CO₂ gas inflow and outflow and for sampling the atmosphere of reactor for chemical analysis. The light intensity on the

photocatalysts was 0.63 W/cm^2 and the spectral composition of light source, examined with Hamamatsu Photonics PMA-11, is shown in Fig. 3.2b. The gas phase was analyzed using a gas chromatograph (Shimadzu GC-8A). The CO and CH₄ production were evaluated using a flame ionization detector equipped with a methanizer (Shimadzu MTN-1) and H₂ and O₂ production were evaluated using a thermal conductivity detector. Here, it should be noted although some systematic studies suggested an optimum TiO₂ concentration of 1 g/L for photocatalytic reactors [25], a higher concentration of 3.3 g/L was used in this study due to the low specific surface area of HPT-processed samples.

3.3. Results and discussion

3.3.1. TiO₂-II phase formations

Examination of phase transformations using the XRD profiles and Raman spectra are shown in Figs. 3.3a and 3b, respectively. As shown in Fig. 3.3a, the initial powder contains the anatase phase, while new peaks corresponding to the TiO₂-II phase appear after HPT processing. The TiO₂-II peaks remain in the XRD profile after annealing and some other peaks corresponding to the rutile phase appear. Quantitative analysis of phase fractions using the Rietveld method, as given in Table 3.2, confirms the high stability of TiO₂-II phase even after annealing at 773 K. The Raman spectra also confirms that the initial powder mainly contains an anatase phase, while a TiO₂-II phase is formed after HPT processing (see the peaks at wavenumber of 171, 283, 316, 340, 357, 428, 533 and 572 cm⁻¹). A rutile phase is also formed after HPT processing and the intensity of its peaks increases after annealing. These results are reasonably consistent with the reported effect of pressure on TiO₂-II phase formation and the temperature effect on the formation of rutile phase, as shown in Fig. 3.1a [1]. However, the stability of TiO₂-II phase at ambient pressure and the formation of rutile phase at relatively low temperature compared to the phase diagram should be due to the effect of straining by HPT, as discussed in earlier publications [26].

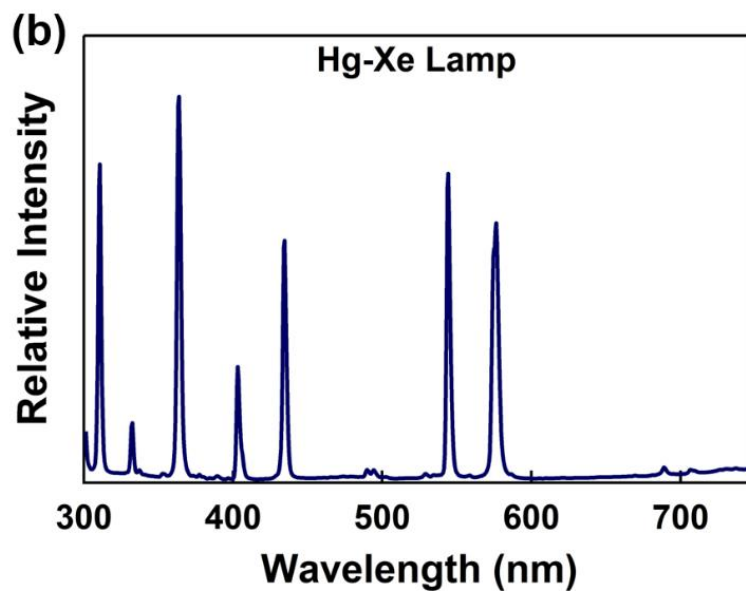
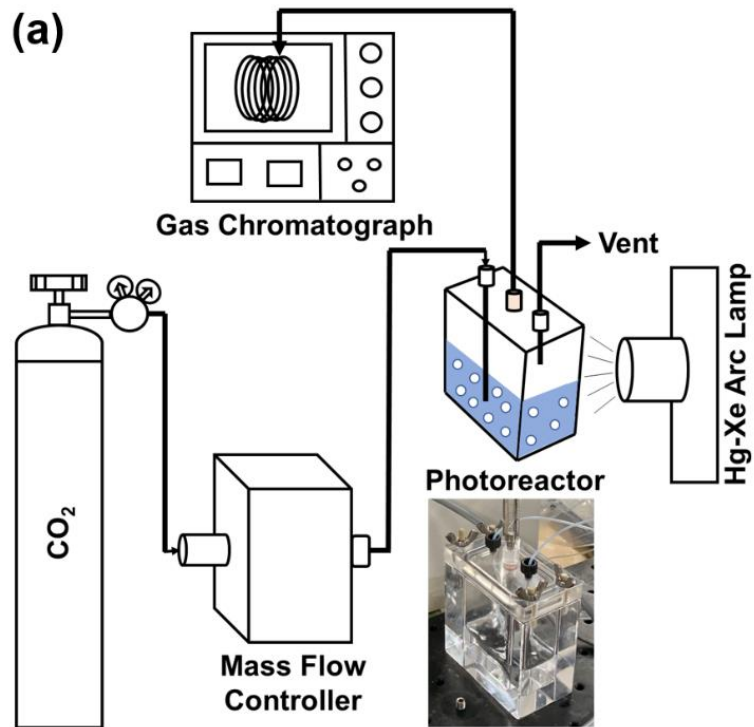


Fig. 3.2. (a) Schematic illustration of photocatalytic CO₂ conversion system and photo of its photoreactor, and (b) spectral composition of He-Xe lamp used for photocatalytic test.

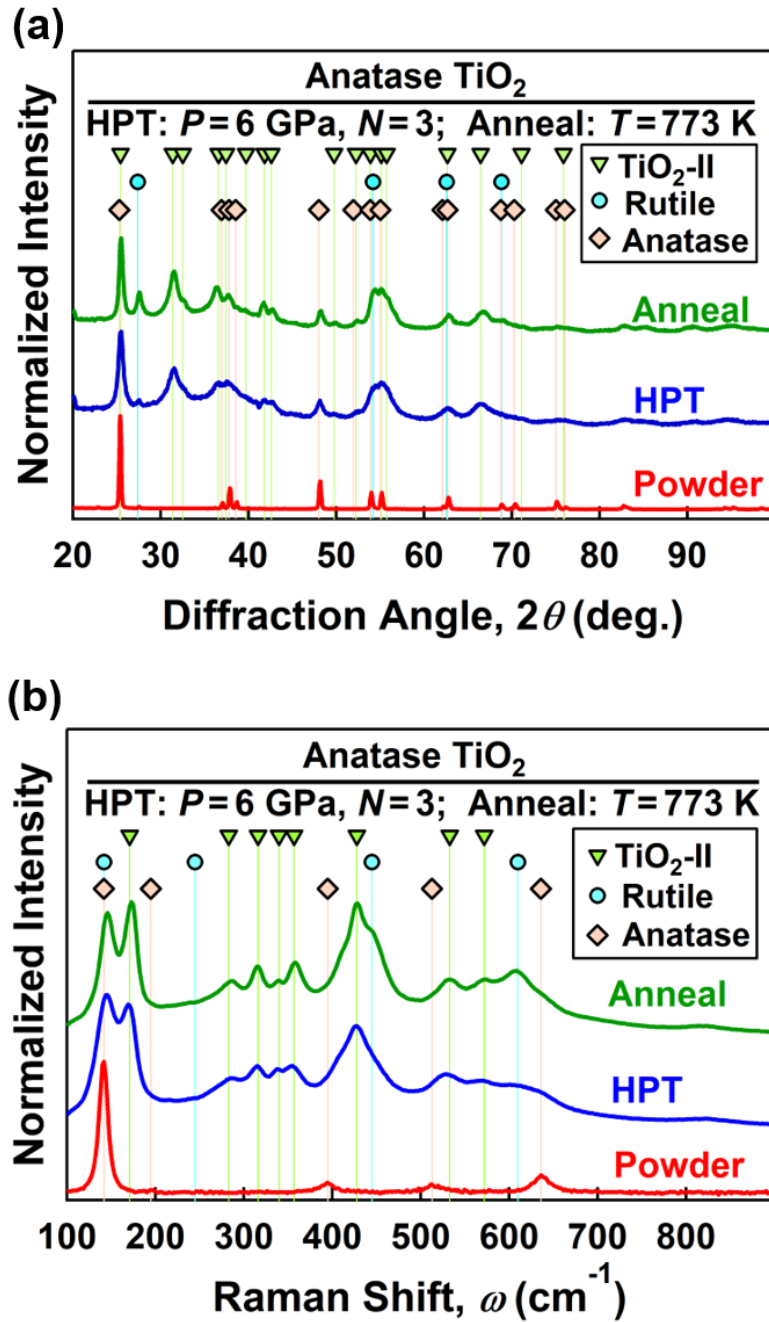


Fig. 3.3. Formation of TiO₂-II phase after HPT processing and its stability during annealing. (a) XRD profiles and (b) Raman spectra for initial powder and for samples processed by HPT and annealing.

Table 3.2. Fraction of anatase, TiO₂-II and rutile phases, average lattice strain and average crystallite size determined from XRD profiles, and mean grain size measured by TEM for initial powder and for samples processed by HPT and annealing.

Sample	Phase Fractions (wt%)			Lattice Strain (%)	Crystallite Size (nm)	Grain Size (nm)
	Rutile	Anatase	TiO ₂ -II			
Powder	2.7	97.3	0	0.1	60	130 ± 40
HPT	11.8	27.5	60.7	0.8	9	32 ± 16
Anneal	23.2	20.0	56.8	0.5	11	43 ± 22

3.3.2. Microstructure of catalysts

Microstructural features examined using TEM BF and DF images and corresponding SAED patterns are shown in Fig. 3.4 for (a-c) initial powder, (d-f) sample processed by HPT and (g-i) samples processed by annealing, where DF images were taken with the diffracted beams indicated by arrows in the SAED patterns. The initial sample has an average particle size of 130 nm and most of the particles contain only one crystal. After HPT processing, the particle sizes become larger due to HPT-induced partial consolidation [18], but the grain sizes (white regions in the DF images) are refined to the range of nanometer due to the effect of severe strain [16,17]. The grain sizes remain at the range of nanometer even after annealing at 773 K and many crystals still exist in one particle. The ring shapes of the SAED patterns for the samples processed by HPT and annealing in comparison with the net shape of the SAED pattern for the initial powder also confirms smaller crystal sizes in these two samples. For the three samples, Table 3.2 gives the mean grain sizes measured by TEM, average crystallite sizes determined by XRD and the average lattice strain determined by XRD. The grain/crystallite size decreases and the lattice strain increases by HPT processing, while the grain/crystallite size slightly increases and the lattice strain increases by subsequent annealing.

To examine the nanostructural features, high-resolution TEM was performed as shown in Fig. 3.5 for (a, b) initial powder, (c, d) sample processed by HPT and (e, f) sample processed by annealing. The initial powder contains large crystals of anatase phase which was confirmed by FFT analyses. After HPT processing, many nanocrystals with the TiO₂-II structure are visible, while some small amount of the rutile phase is also detectable. The nanostructural changes after

annealing are not so significant and many nanocrystalline with the TiO₂-II phase are still visible within the microstructure. In addition to nanograins, many TiO₂-II/anatase phase boundaries are formed by HPT processing. These results are consistent with the phase analyses in Fig. 3.3, microstructural analysis in Fig. 3.4 and reported data about the effect of strain on the microstructural evolution of various oxides [17-21].

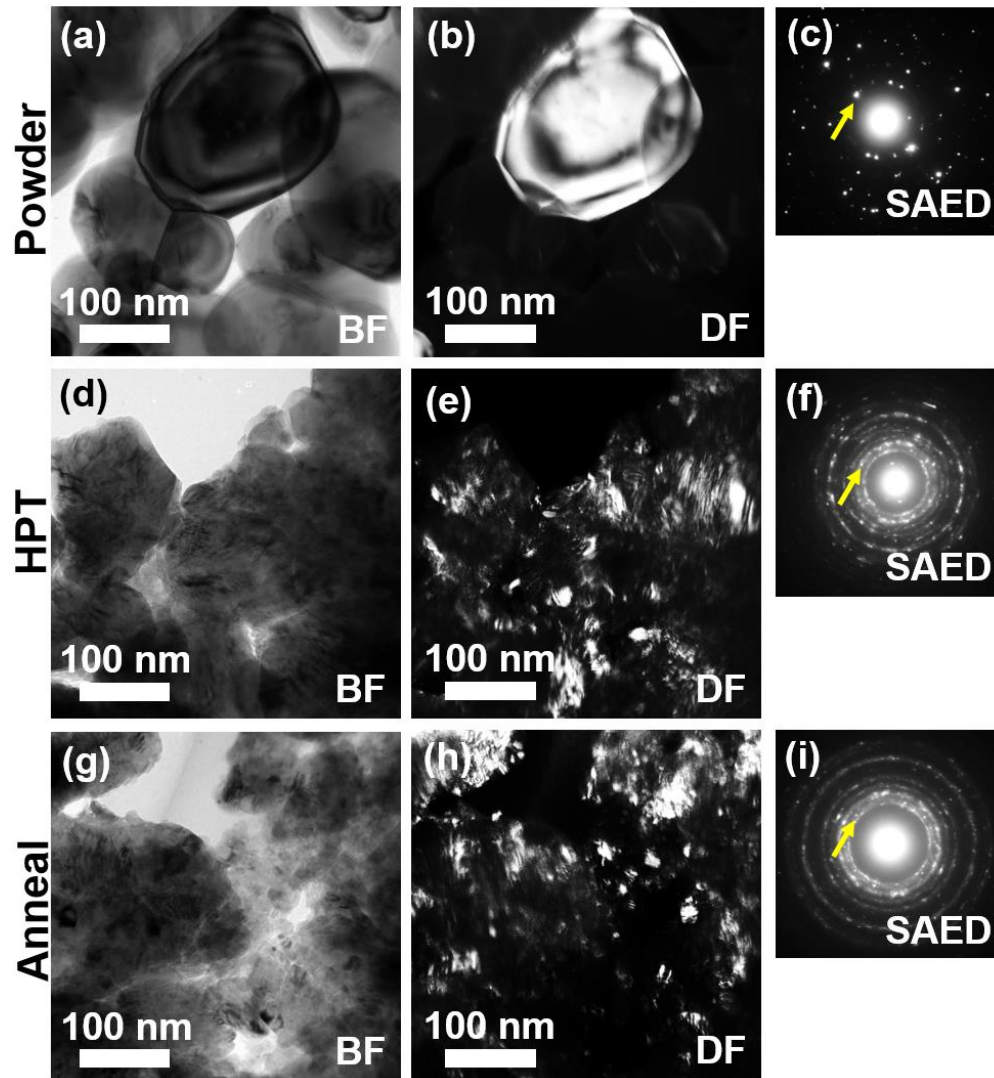


Fig. 3.4. Reduction of crystal size after HPT processing and stability of nanocrystals during annealing. TEM BF images, DF images and SAED patterns for (a-c) initial powder, (d-f) HPT-processed sample and (g-i) annealed sample.

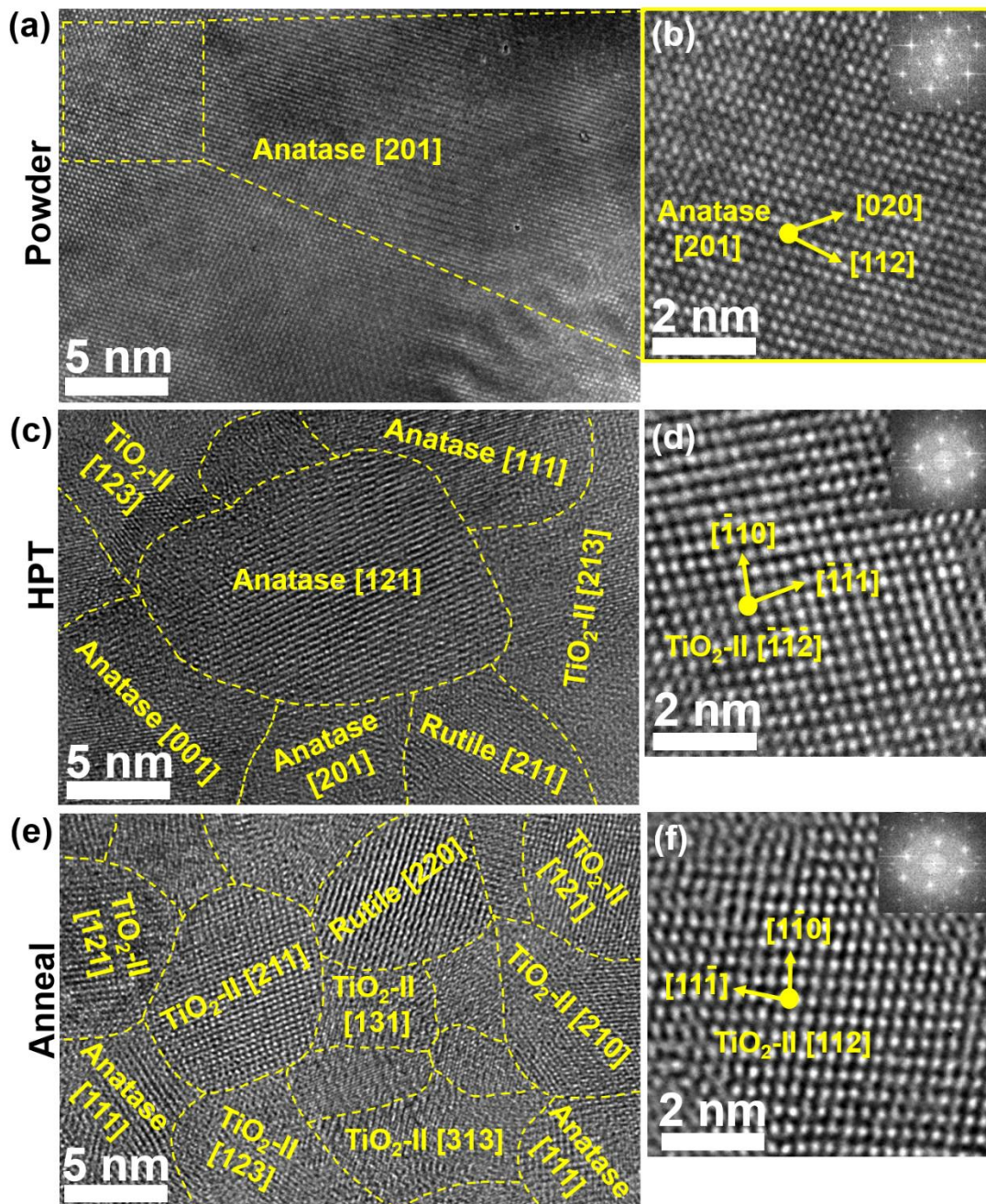


Fig. 3.5. Co-existence of nanocrystalline $\text{TiO}_2\text{-II}$ phase with anatase and rutile phases after HPT processing and annealing. TEM (a, c, e) high-resolution images and (b, d, f) lattice images and corresponding FFT diffractograms for (a, b) initial powder, (c, d) HPT-processed sample and (e, f) annealed sample.

3.3.3. Vacancy formation/annihilation

The appearance of three samples is shown in Fig. 3.1c. The initial anatase powder has a white color, but its color becomes dark green after HPT processing, indicating the formation of color centers such as oxygen vacancies in the sample [20,21]. The color of samples becomes almost white after annealing, indicating that many color centers are thermally annihilated. The formation of oxygen vacancies was examined by XPS, as shown in Fig. 3.6 for (a) O 1s and (b) Ti 2p.

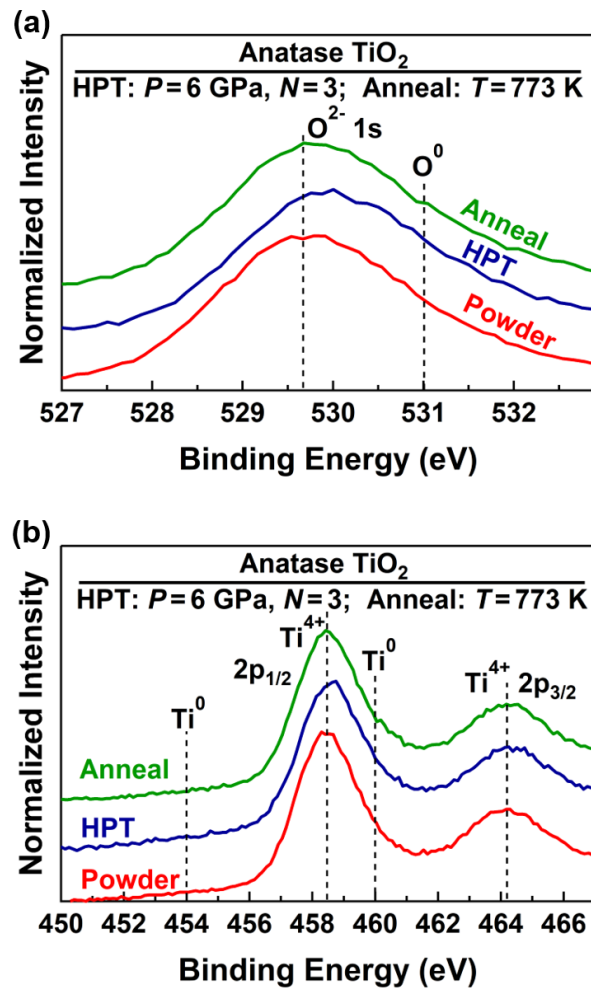


Fig. 3.6. Formation of oxygen vacancies by HPT processing and their annihilation after annealing. XPS profiles with (a) O 1s and (b) Ti 2p for initial powder and for samples processed by HPT and annealing.

Note that the XPS peak positions were corrected by considering the C 1s peak position at 284.8 eV. Figs. 3.6a and 5b demonstrate the presence of oxygen anion and titanium cations in all samples, but the main difference is that the O 1s peak shifts slightly to higher energies after HPT processing and return to its original position after annealing (Ti 2p peaks do not show clear changes). These shifts, which are consistent with the changes in the color of samples, are probably due to the formation and annihilation of oxygen vacancies after HPT processing and annealing, respectively. The formation of vacancies by HPT processing is a consequence of strain effect on microstructure which was reported in various metallic and ceramic materials [16-18].

The nature of oxygen vacancies was examined by EPR, as shown in Fig. 3.7. The initial powder contains various characteristic peaks of the anatase phase [27]. After HPT processing, characteristic peaks of anatase disappear and several pair peaks appear which corresponds to different defects: the one with a g value higher than 2 (i.e., inflection point at 2.004) is related to hole trapping sites such as oxygen-centered radicals or oxygen vacancies, and those with the g values smaller than 2 (i.e., inflection points at 1.991 and 1.984) are related to electron trapping sites such as Ti^{3+} radicals in the bulk [27-29]. Since the Ti^{3+} sites in the bulk are considered as recombination sites [30], their presence is not appropriate for photocatalytic activity [31,32]. After annealing, the peak pairs corresponding to the Ti^{3+} radicals disappear, while the oxygen vacancy peak is still visible. It should be noted that the presence of oxygen vacancies on the surface with an optimum concentration is positive for photocatalytic activity because they can act as active sites for reaction, but their presence in the bulk can enhance the electron-hole recombination and destroy the photocatalytic activity [30].

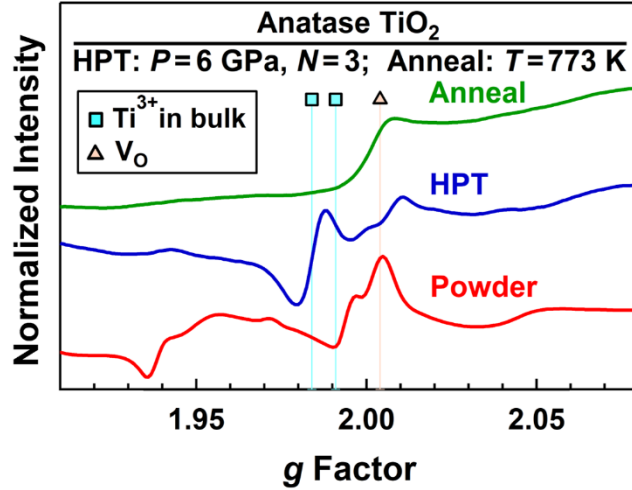


Fig. 3.7. Formation of Ti^{3+} radicals and oxygen vacancies (V_O) by HPT processing and disappearance of Ti^{3+} radicals in bulk by annealing. EPR spectra for initial powder and for samples processed by HPT and annealing.

3.3.4. Light absorbance and band structure

The light absorbance and band structure of samples were examined by UV-vis spectroscopy and XPS and compared with the energy requirements for the photocatalytic CO_2 conversion. UV-vis spectra are shown in **Fig. 3.8a**, indicating that the three samples mainly absorb the UV light. The light absorbance in the visible light region increases by HPT processing and decreases after annealing. **Fig. 3.8b** shows the bandgap calculation using the Kubelka-Munk theory. The bandgap of initial powder is 3.0 eV and it decreases to 2.5 eV after HPT processing and reaches 2.7 eV after annealing. The estimated bandgap for the HPT-processed sample is close to the reported theoretical bandgap of TiO_2 -II [33,34]. Therefore, the formation of TiO_2 -II phase and oxygen vacancies should have led to higher light absorbance and narrowing the optical bandgap. Increasing the bandgap after annealing should be due to the annihilation of oxygen vacancies and partial reduction of the TiO_2 -II phase fraction. **Fig. 3.8c** demonstrates the XPS profiles to calculate the top of valence band. Top of valence band for the initial powder, HPT-processed sample and annealed sample are 2.2, 2.1 and 2.0 eV, respectively. From the calculated bandgaps, the bottom of conduction band can be calculated by subtraction of bandgap from the top of valence band. **Fig. 3.8d** summarizes the band structure of three samples. All samples satisfy the energy requirement for CO_2 conversion [35,36]. Higher level of conduction band in the annealed sample compared

with the HPT-processed sample suggests that the annealed sample can thermodynamically be more appropriate for CO₂ reduction [37,38].

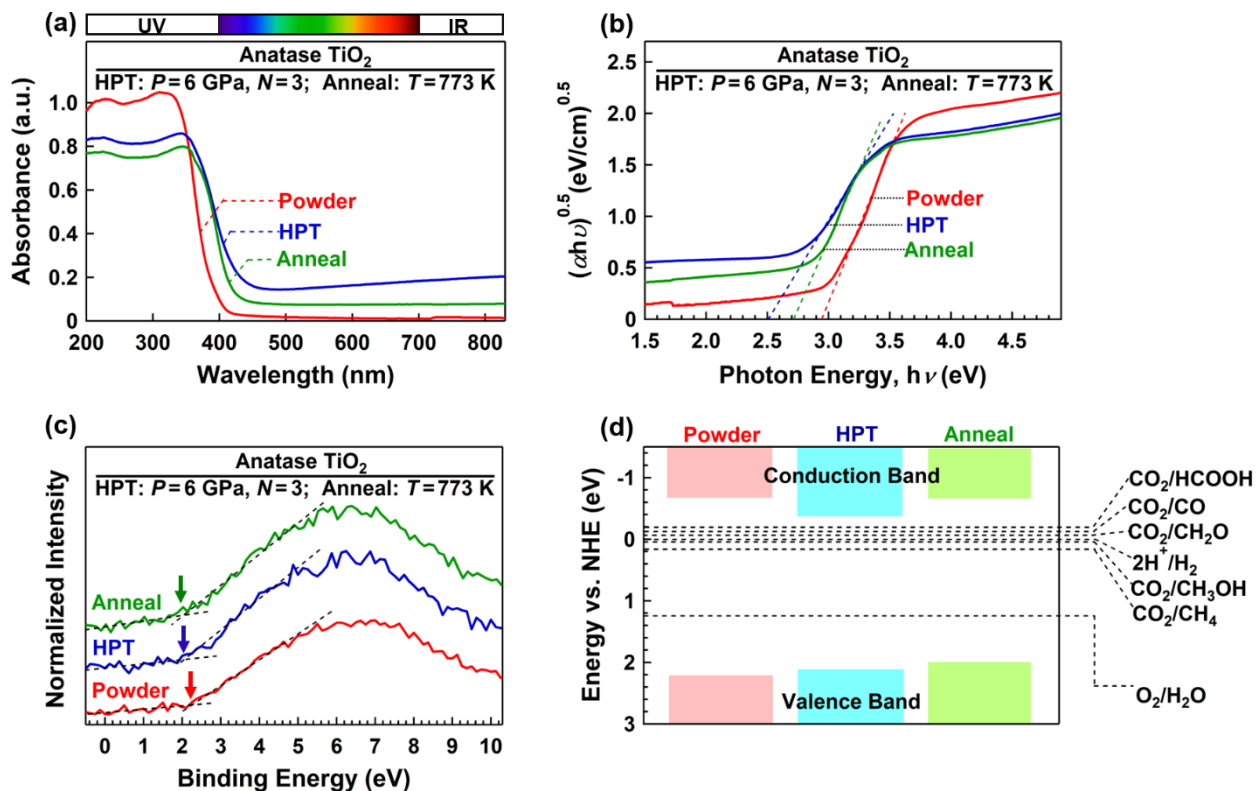


Fig. 3.8. Larger light absorbance and improvement of band structure for CO₂ conversion by introduction of TiO₂-II phase. (a) UV-vis spectra, (b) bandgap calculation using Kubelka-Munk analysis (α : light absorption, h : Planck's constant, ν : photon frequency), (c) XPS of top of valence band and (d) band structure compared to energy requirements for CO₂ conversion for initial powder and for samples processed by HPT and annealing.

3.3.5. Charge carrier decay and photocurrent generation

The steady-state PL emission spectroscopy (Fig. 3.9a), PL decay spectroscopy (Fig. 3.9b) and photocurrent measurements (Fig. 3.9c) were used to examine the mobility of excited charge carriers for movement to the activation/recombination sites. The steady-state PL spectra in Fig. 3.9a shows that the PL intensity for the HPT-processed and annealed samples is significantly lower than that for the anatase powder. This suggests that the recombination of electrons and holes is suppressed in these two samples. Three PL peaks are detectable in Fig. 3.9a at three wavelengths:

~420 nm corresponding to the band-to-band transition [39], ~510 nm corresponding to surface oxygen vacancies or shallow bulk traps in anatase [40], and ~830 nm (mainly in the annealed sample) corresponding to the deep traps in rutile [41].

The PL decay intensity versus time, $I(t)$, exhibits a two-exponential kinetic function in Fig. 3.9a [42].

$$I(t) = A_1 \exp\left(-\frac{t}{\tau_1}\right) + A_2 \exp\left(-\frac{t}{\tau_2}\right) \quad (3.3)$$

In this equation, A_1 and A_2 are the amplitude of each exponential function and τ_1 and τ_2 are the fast decay time and the slow decay time, respectively. As summarized in Table 3.3, the decay times decrease by HPT processing and decrease further by annealing. Here it should be noted that the decay time estimation of TiO₂ by PL decay spectroscopy can be influenced by both radiative and non-radiative recombination [43]. For the HPT-processed and annealed samples with surface defects as shallow traps or photocatalytic activation sites, the fast movement of charge carriers results in fast decay time [44-46]. The hetero-phase junctions such as anatase/rutile or anatase/TiO₂-II can also transfer the electron to the surface defects and reduce the decay time [47]. In contrast, the bulk defects which are formed by HPT processing and disappear by annealing can act as recombination centers and lead to a decrease in the lifetime and a decrease in the photocatalytic activity [48,49]. Decreasing the decay time by HPT processing and annealing also indicate that the change of the electronic structure between the ground state and the excited state decreases under light irradiation. This phenomenon can lead to decreasing the re-organization energy and improving the photocatalytic activity.

Photocurrent measurements, as shown in Fig. 3.9b, show that the current density for initial powder is smaller than the HPT-processed and annealed samples and the current density for the annealed sample is somehow higher than that for the HPT-processed sample. The reason for higher photocurrent density of annealed and HPT-processed samples can be due to lower bandgap, the presence of TiO₂-II/anatase heterojunctions which can ease the separation of electrons from the valence band, and the presence of vacancies as shallow traps [26]. Better photocurrent for the annealed sample should be due to the thermal annihilation of bulk vacancies as recombination sites [26]. The high photocurrent density shows the efficient charge separation which can contribute to increasing the CO₂ conversion efficiency on the HPT-processed and annealed samples. Here it should be noted that the authors' attempt to estimate the flat-band potential and carrier density

using the Mott-Schottky method [50] was not successful due to the technical limits in making uniform films with perfect FTO/TiO₂ and TiO₂/TiO₂ connections.

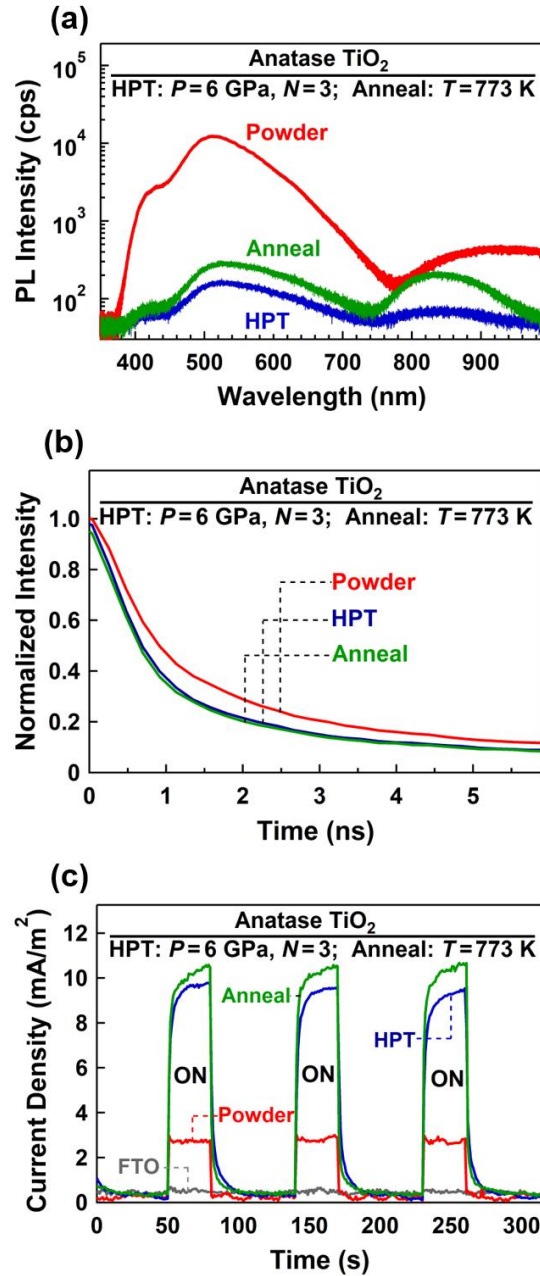


Fig. 3.9. Suppression of recombination and enhancement of charge carrier mobility by HPT processing and annealing. (a) Steady-state photoluminescence emission spectra, (b) time-resolved photoluminescence decay spectra and (c) photocurrent density against time for initial powder and for samples processed by HPT and annealing. ON in (c) indicates application of UV light irradiation.

Table 3.3. Decreasing decay time by HPT processing and subsequent annealing. Fitted parameters of PL decay spectra for initial powder and for samples processed by HPT and annealing.

Samples	τ_1 (ns)	τ_2 (ns)	A_1	A_2
Powder	1.24	11.46	41.98	58.02
HPT	0.95	11.40	40.23	59.77
Anneal	0.93	10.97	39.65	60.35

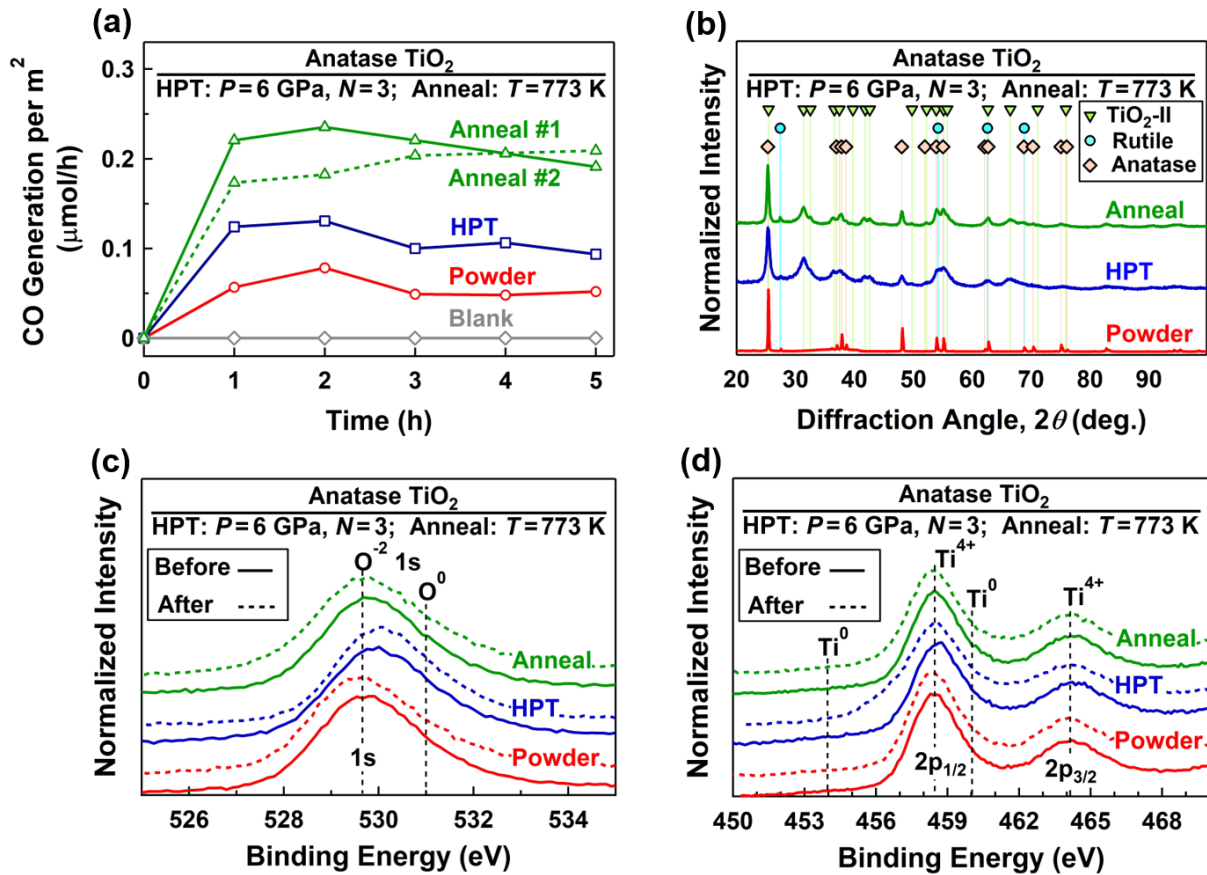


Fig. 3.10. Enhancement of photocatalytic CO₂ conversion by introduction of TiO₂-II phase via HPT processing and further activity enhancement by removing vacancies from bulk via annealing. (a) Amount of CO produced by photocatalytic CO₂ conversion under UV light, (b) XRD profiles after photocatalytic tests, and XPS profiles of (c) O 1s and (d) Ti 2p before and after photocatalytic test for initial powder and for samples processed by HPT and annealing.

3.3.6. Photocatalytic CO₂ conversion

The activity of samples under UV light was examined by photocatalytic CO₂ conversion, as attempted in earlier studies [2,35,36,51-54]. As shown in Fig. 3.10a, all samples convert CO₂ to CO without addition of a co-catalyst, and similar to photocurrent measurements, the efficiency of annealed sample is higher than the two other samples (the standard deviation of gas chromatography for three different tests was less than 10%). Independent synthesis of the annealed sample (as the most active catalyst) and repeating the photocatalytic test, as shown as Anneal #2 in Fig. 3.10a, also confirm a reasonable reproducibility of the synthesis method and CO generation. Examination of gas and liquid phases by gas chromatography and nuclear magnetic resonance suggests that CO is the only product of photoreaction within the detection limits of the measurement system. Blank test by illumination of solution and without the addition of catalysts does not lead to the production of CO, as shown in Fig. 3.10a. Moreover, none of the samples show CO generation after 1 h process in dark (i.e. the data points at the time of 0 in Fig. 3.10a). It is worth mentioning that the authors' earlier attempts for water splitting using these samples under UV light resulted in no H₂ or O₂ generation without addition of a co-catalyst and sacrificial agent.

Examination of the stability of catalysts by XRD after photocatalytic test, as shown in Fig. 3.10b, confirms that the bulk of all samples remains stable during the photocatalytic test. Moreover, XPS analysis before and after photocatalytic test, as shown in Fig. 10c for O 1s and in Fig. 3.10d for Ti 2p, confirms that the surface of samples also remains stable after photocatalytic test. Although the recyclability of these three samples were not examined in detail in this study, an annealed sample used for achieving the photocatalytic data of Annealed #1 in Fig. 3.10a, was reused for photocatalysis after 10-month storage in air atmosphere. Its CO production rate per surface area was determined as 0.15 μmol/h within 1 h which is 30% lower than the activity of the sample in the first cycle. Although this experiment suggests that the sample was deactivated after recycling, but its CO production rate is still better than the activity of anatase powder. Taken altogether, the current results confirm the significance of TiO₂-II phase for photocatalytic CO₂ conversion, although the quantum yield analysis should be conducted in future to quantify its yield efficiency compared to other photocatalysts [55].

To understand the binding mode of CO₂ with the surface of photocatalysts, the ATR-FTIR spectroscopy was conducted before and after immersion of the samples in 1 M NaHCO₃. Fig. 3.11a shows the difference between the intensity of ATR-FTIR spectra before and after immersion. For

the three samples, a peak appears at wavenumbers of 2300-2400 cm^{-1} and the main difference is the higher intensity of peak for the annealed sample compared to the two other samples. Such peaks are typical of CO_2 in air or physisorbed CO_2 on the surface of oxides [56,57]. The ATR-FTIR spectra in Fig. 3.11a provide no clear evidence for carboxylate, bicarbonate, monodentate, bidentate and tridentate modes [58]. Here, it should be noted that NaHCO_3 was used in the photocatalytic tests of this study mainly to enhance the solubility of CO_2 in the solution and to enhance the basicity [59].

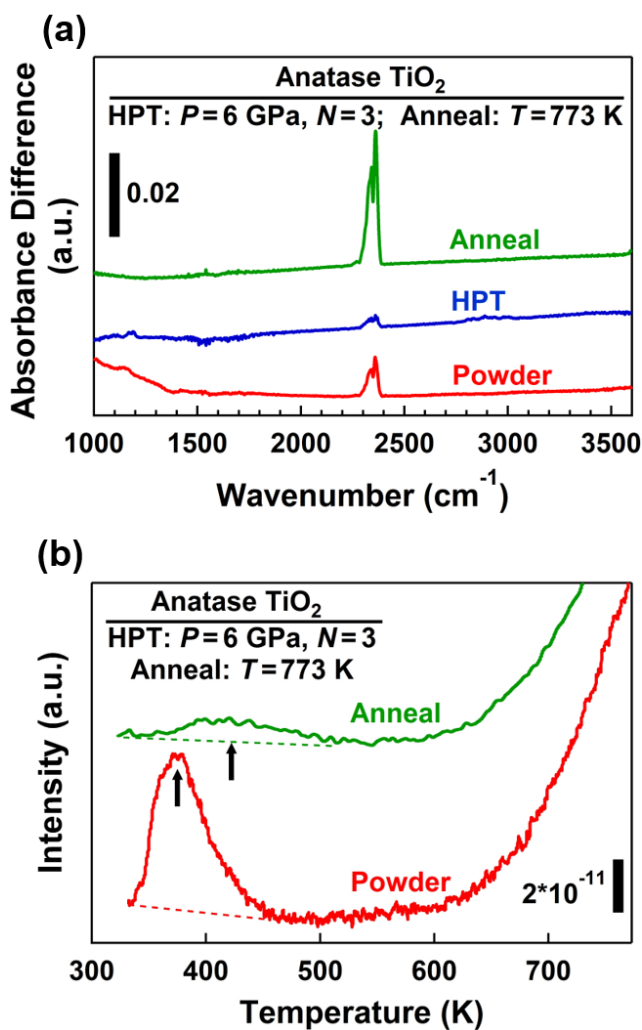


Fig. 3.11. Adsorption behavior of CO_2 on surface of photocatalysts. (a) Difference between ATR-FTIR spectra before and after immersion of photocatalysts in 1 M NaHCO_3 for initial powder and samples processed by HPT and annealing, and (b) TPD spectra after CO_2 adsorption at 773 K for initial powder and annealed sample.

To further investigate the adsorption of CO₂ to the surface of photocatalysts, TPD was conducted. None of the samples could adsorb CO₂ at 323 K perhaps because of weak basicity of the three samples. Here, it should be noted that a main difference between the TPD tests and photocatalytic reaction is that the photocatalytic test is conducted in an aqueous solution of 1 M NaHCO₃, but the TPD test is conducted under gaseous atmosphere. To have an insight into the CO₂ adsorption on the surface of samples, the adsorption temperature was increased to 773 K. At this temperature, the initial powder and the sample annealed at 773 K are stable, but the HPT-processed sample shows structural changes and cannot be tested by TPD. Both powder and annealed samples show small amounts of hydrogen adsorption, with the desorption behaviors shown in Fig. 3.11b. The powder sample has an intense peak at low temperature, but the annealed sample reveals a low-intensity broad peak but at higher temperatures (no CO or O₂ desorption was detected for the samples). These TPD results suggest that the bonding of CO₂ to the surface of annealed sample is somehow stronger than that for the initial anatase powder, although the amount of adsorbed CO₂ on the surface of initial powder is higher perhaps because of its larger surface area or its higher basicity in the gaseous atmosphere.

Three main pathways for CO₂ photoreduction including carbene, formaldehyde and glyoxal have been suggested by in situ examinations; however, there are still significant arguments regarding the photocatalytic CO₂ conversion pathways on TiO₂ due to the effects of impurities, structure and lattice defects [60,61]. The absence of HCOOH, CH₃OH and CH₄ in the liquid and gas phases within the detection limits of our experiments suggests that the formaldehyde may not be the main pathway for CO₂ photodegradation. Moreover, the absence of HCOOH and CH₄ in our study suggests that the glyoxal may not be the main pathway. The production of only CO suggests that the carbene can be the main pathway, although the nonappearance of CH₄ indicates that the carbene pathway probably stops at intermediate stages [58]. One reason for the absence of CH₄ can be explained by the presence of surface defects. Since the TiO₂ samples in this study have surface defects such as oxygen vacancies, CO₂ in combination with H₂O as a Lewis acid desires to adsorb on oxygen vacancy sites [58]. Such adsorption can lead to degradation of C=O bonding to produce [•]CO radicals and subsequently produce CO gas. Since CO is the main product and it is not willing to adsorb on the surface defects again [58], the carbene pathway at this stage does not move forward appreciably to produce detectable amount of CH₄. Some other studies also reported that although the formation of CH₄ is thermodynamically more favorable than the formation of

CO, the CH₄ formation is kinetically more difficult because it needs more electrons and protons [55].

3.3.7. Role of TiO₂-II and comparison with other photocatalysts

Taken altogether, the current results on the first application of TiO₂-II phase for photocatalytic CO₂ conversion confirm the high potential of the phase to enhance the activity compared to pure anatase phase. Unlike doped photocatalysts [3-11] which may experience impurity-induced recombination [29], the introduction of TiO₂-II modifies the band structure and enhances the activity for CO₂ conversion without addition of impurities. Examination of light absorbance, PL, PL decay and photocurrent suggests that the formation of TiO₂-II not only enhances the light absorbance but also suppress the recombination and improve the charge carrier mobility to surface to generate photocurrent and contribute to photocatalytic reactions. Moreover, the introduction of TiO₂-II induces heterojunctions with other phases such as anatase and rutile which can further enhance the charge separation for photocatalysis [47,62]. The presence of phase heterojunctions, formation of surface vacancies, narrowing the bandgap and easier separation and mobility of charge carriers should be responsible for the improvement of the photocatalytic activity of the samples after HPT processing and annealing. For the annealed sample, removing the vacancies from the bulk by annealing and increasing the energy position of conduction band should contribute to higher photocatalytic activity of this sample compared with the HPT-processed sample. Based on the ATR-FTIR and TPD spectra, CO₂ may make a stronger adsorption to the surface of annealed sample, which can be another reason for higher activity of this catalyst.

To have an insight into the photocatalytic activity of the annealed sample (i.e. the TiO₂-II-containing catalyst in this study), its CO production rate is compared with some reported data in the literature [63-73]. The reported data deviate in a wide range of 0.01-10.16 $\mu\text{mol g}^{-1}\text{h}^{-1}$. Such a wide range of variations is not only due to the differences in the activity of photocatalysts, but also due to the differences in the experimental setup such as light source, gas flow rate, catalyst concentration, specific surface area, etc. Despite differences in the experimental setup, Table 3.4 suggests that the first application of TiO₂-II phase for photocatalytic CO₂ conversion has led to a good activity ($1.39 \pm 0.13 \mu\text{mol.g}^{-1}\text{h}^{-1}$ CO production rate) compared to the reported data in the literature. We consider that this activity can be further improved in future by combining the effect of TiO₂-II with the other strategies developed so far to enhance the activity such as optimizing the

synthesis method, production of composites, and production of nanosheets, nanoparticles or mesoporous samples with large specific surface area.

Table 3.4. Photocatalytic CO production on annealed TiO₂ sample, compared with reported data for other photocatalysts.

Photocatalyst	Light Source	CO Production ($\mu\text{molg}^{-1}\text{h}^{-1}$)	Reference
TiO ₂ with TiO ₂ -II by HPT and Anneal	300 W Hg-Xe	1.39±0.13	This Study
TiO ₂ Nanosheet	150 W Xe	0.16	[63]
TiO ₂ Isotropic	150 W Xe	1.32	[63]
TiO ₂ Nanosheet - CN	150 W Xe	2.04	[63]
TiO ₂ Isotropic - CN	150 W Xe	1.55	[63]
TiO ₂ P25	150 W Xe	1.49	[63]
TiO ₂ Nanosheet - CN Nanosheet	150 W Xe	0.12	[63]
TiO ₂ Nanosheet - CN	150 W Xe	0.80	[63]
TiO ₂ P25	300 W Xe	1.84	[64]
TiO ₂ Mesoporous	300 W Xe	7.0	[64]
TiO ₂ - SiO ₂	300 W Xe	6.47	[64]
TiO ₂ - Graphitic Carbon	300 W Xe	10.16	[64]
TiO ₂ Nanosheets Exposed {001} Facet	2x18W Hg	0.12	[65]
TiO ₂ Pure	150W UV410	0.68	[66]
TiO ₂ - CoO _x	150W UV410	0.90	[66]
TiO ₂ - CoO _x Calcined in Air	150W UV410	0.57	[66]
TiO ₂ Calcined in N ₂ /H ₂	150W UV410	0.49	[66]
TiO ₂ Hydrogenated - CoO _x	150W UV410	0.81	[66]
TiO ₂ - CoO _x Hydrogenated	150W UV410	1.24	[66]
TiO ₂ P25	300 W Xe	1.8	[67]
TiO ₂ 3D Ordered Microporous	300 W Xe	2.9	[67]
TiO ₂ 3D Ordered Microporous - Au	300 W Xe	0.3	[67]
TiO ₂ 3D Ordered Microporous - Au ₃ Pd	300 W Xe	1.2	[67]
TiO ₂ 3D Ordered Microporous -Au ₂ Pd ₂	300 W Xe	1.8	[67]
TiO ₂ 3D Ordered Microporous - AuPd ₃	300 W Xe	2.6	[67]
TiO ₂ 3D Ordered Microporous - Pd	300 W Xe	3.9	[67]

C ₃ N ₄ Nanosheet	150 W Xe	0.19	[63]
C ₃ N ₄ Graphitic	150 W Xe	0.01	[63]
C ₃ N ₄ by Thermal Condensation	350 W Xe	4.83	[68]
C ₃ N ₄ by Solvothermal Process	350 W Xe	1.12	[68]
C ₃ N ₄ Solvothermal Melamine	350 W Xe	0.43	[68]
C ₃ N ₄ Solvothermal Melamine	350 W Xe	1.69	[68]
C ₃ N ₄ by Carbon Tetrachloride Process	350 W Xe	0.33	[68]
C ₃ N ₄ by Cyclohexane Process	350 W Xe	1.23	[68]
C ₃ N ₄ by Tetrahydrofuran Process	350 W Xe	2.13	[68]
C ₃ N ₄ by Solvothermal and Calcination	350 W Xe	3.52	[68]
Cd _{1-x} Zn _x S	100 W LED	2.9	[69]
BiOI	300 W Xe	4.1	[70]
xCu ₂ O - Zn _{2-2x} Cr	200-W Hg-Xe	2.5	[71]
CeO _{2-x}	300 W Xe	1.65	[72]
Cu ₂ O - RuO _x	150 W Xe	0.88	[73]

3.4. Conclusion

Effect of high pressure TiO₂-II phase on photocatalytic CO₂ conversion was examined for the first time. The introduction of nanostructured TiO₂-II phase and TiO₂-II/anatase heterojunctions modified the band structure and enhanced the photocurrent and photocatalytic activity, despite the formation of strain-induced vacancies in the bulk. Thermal annihilation of vacancies in the bulk led to further improvement of photocatalytic CO₂ conversion, confirming the significance of the TiO₂-II phase for photocatalytic reactions.

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Chapter 4. Defective high-entropy oxide photocatalyst with high activity for CO₂ conversion

4.1. Introduction

High-entropy ceramics are a new type of materials which show high stability and promising structural and functional properties due to the so-called cocktail effect, lattice strain/defects, heterogenous valence electron distribution and high configurational entropy [1]. As shown in Fig. 4.1a, high-entropy ceramics are defined as multi-component materials with at least five principal elements and a configurational entropy higher than $1.5R$ (R : the gas constant) [2]. These materials have a low Gibbs free energy due to their high entropy and this gives a high stability to these materials under different conditions [1] including catalytic reactions [3,4]. Moreover, the presence of at least five cations with different atomic sizes in these materials results in the formation of inherent lattice strain and defects [2]. Since lattice strain and defects are effective to enhance photocatalytic CO₂ conversion [5-8], the high-entropy ceramics are expected to show good activity for such a conversion. High-entropy oxides (HEOs) are the most popular high-entropy ceramics which have been investigated for various applications and properties such as thermal barrier coatings [9,10], magnetic components [11,12], dielectric components [13,14], Li-ion batteries [15,16], Li-S batteries [17], Zn-air batteries [18], catalysts [19,20], electrocatalysts [21], and photocatalytic hydrogen production [22,23]. Despite the inherent defective and strained structure of HEOs, there have been no attempts to employ these materials for photocatalytic CO₂ conversion.

In this chapter, a HEO photocatalyst, TiZrNbHfTaO₁₁, is synthesized and its activity for CO₂ conversion is examined. The transition elements titanium, zirconium, niobium, hafnium, and tantalum are selected simply because their binary oxides with the d⁰ electronic structure can act as photocatalysts. This first application of HEOs for photocatalytic CO₂ conversion confirms that the HEO photocatalyst shows higher activity compared to common binary or ternary photocatalysts such as TiO₂ and BiVO₄, suggesting HEOs as a new family of photocatalysts for CO₂ conversion.

4.2. Experimental

4.2.1. Sample preparation

Although various methods have been developed in recent years for the synthesis of HEO [1-4,9-23], a two-step high-pressure mechanical alloying and high-temperature oxidation which is available in the authors' laboratory was used to synthesize the HEO. In the first step, equiatomic amounts of Ti (99.9%), Zr (95.0%), Hf (99.5%), Nb (99.9%) and Ta (99.9%) powders were mixed in acetone, treated by ultrasonic and then dried. The dried powder mixture was processed by high-pressure torsion (HPT), shown in Fig. 4.1b, to fabricate TiZrNbHfTa high entropy alloy (HEA) with the body-centered cubic (BCC) structure (see the principles of HPT and its applications to oxides in [24,25]). To produce the TiZrNbHfTa alloy, a 10 mm diameter and 1 mm thick disc was prepared by compacting the powder mixture under a pressure of 400 MPa. The compacted disc was then compressed between two HPT anvils under a high pressure of 6 GPa at room temperature and simultaneously processed by rotating the lower HPT anvil with respect to the upper one for 100 turns with a rotation rate of one turn per minute. In the second step, the HPT-processed TiZrNbHfTa alloy was exposed to hot air at a temperature of 1373 K for 24 h to produce an oxide, with the appearance shown in Fig. 4.1c. Examination of the mass of sample before and after oxidation suggested a composition of TiZrNbHfTaO₁₁ for the produced oxide. The oxide was crushed after oxidation into the powder form using a mortar and examined by various characterization methods, as described below.

4.2.2. Characterization

To examine the crystal structure, X-ray diffraction (XRD) using the Cu K α radiation with a wavelength of $\lambda = 0.1542$ nm and micro-Raman spectroscopy using a laser source with a wavelength of $\lambda = 532$ nm were utilized.

Examination of microstructure was conducted by (i) scanning electron microscopy (SEM) with energy dispersive X-ray spectroscopy (EDS) analysis under 15 keV, (ii) transmission electron microscopy (TEM) with selected area electron diffraction (SAED), bright-field (BF) images, dark-field (DF) images, high-resolution images and fast Fourier transform (FFT) analysis under 200 keV, and (iii) scanning-transmission electron microscopy (STEM) with high-angle annular dark-field (HAADF) images and EDS analysis under 200 keV.

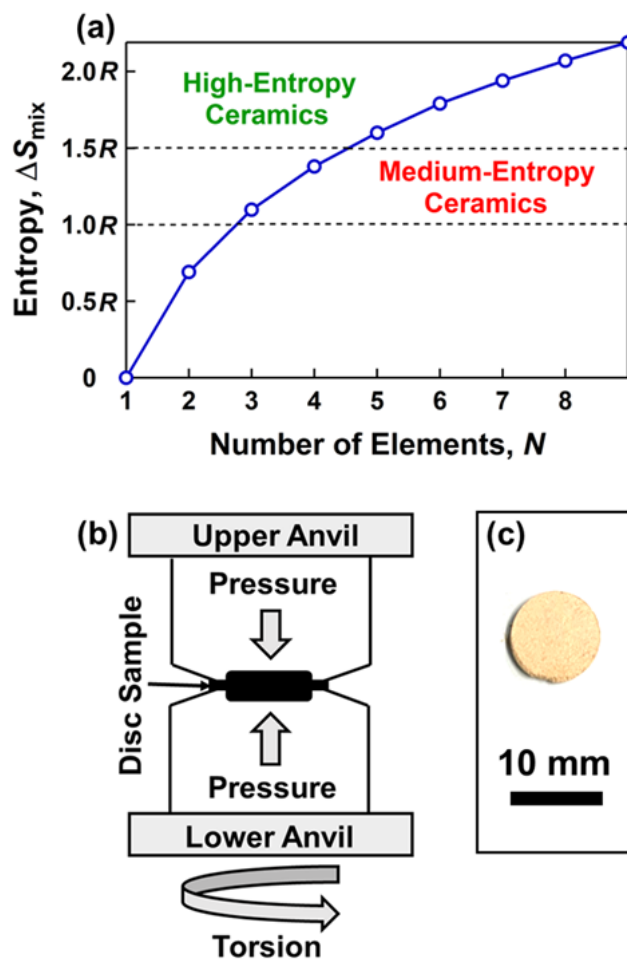


Fig. 4.1. (a) Relationship between number of elements and configurational entropy and definition of high-entropy ceramics with equiatomic fractions of elements, (b) schematic illustration of high-pressure torsion (HPT), and (c) appearance of high-entropy oxide synthesized in this study.

To investigate the presence of point defects such as oxygen vacancies, electron paramagnetic resonance (EPR) was performed at ambient temperature using a microwave source with a frequency of 9.4688 GHz.

To study the oxidation states of different elements and to estimate the valence band top position, X-ray photoelectron spectroscopy (XPS) using the Al K_{α} radiation with a wavelength of $\lambda = 0.989$ nm was used. The XPS energy position for each element was adjusted by considering the peak position of C 1s at 284.8 eV. After correction of the energy positions, the peaks for different elements were analyzed by peak deconvolution by considering the standard energy

relations and differences reported in the handbook [26]: $f_{7/2}:f_{5/2} = 4:3$, $d_{5/2}:d_{3/2} = 3:2$, $p_{3/2}:p_{1/2} = 2:1$, $Ti\ 2p_{1/2} - Ti\ 2p_{3/2} = 5.54\ eV$, $Zr\ 3d_{3/2} - Zr\ 3d_{5/2} = 2.43\ eV$, $Hf\ 4f_{5/2} - Hf\ 4f_{7/2} = 1.71\ eV$, $Nb\ 3d_{3/2} - Nb\ 3d_{5/2} = 2.72\ eV$, $Ta\ 4f_{5/2} - Ta\ 4f_{7/2} = 1.91\ eV$.

To investigate the light absorbance and bandgap (Kubelka-Munk analysis), UV-vis diffuse reflectance spectroscopy was conducted, and the band structure was calculated by considering both XPS and UV-vis spectra.

To study the lifetime of excited electrons, steady-state photoluminescence (PL) emission spectroscopy with a 325 nm laser source and time-resolved photoluminescence decay (PL decay) with a 285 nm laser source were conducted.

The specific surface area of powder was examined by nitrogen gas adsorption and using the Brunauer-Emmett-Teller (BET) method.

4.2.3. Photocurrent test

Photocurrent generation was examined using a thin film of sample in a 1 M Na_2SO_4 electrolyte under the full arc of Xe lamp (without using any filter), as described in detail earlier [27]. The thin film was prepared by deposition of HEO powder on FTO (fluorine-doped tin oxide) glass with 2.25 mm thickness and $15 \times 25\ mm^2$ surface area. About 5 mg of sample was crushed in 0.2 mL ethanol and carefully dispersed on the FTO glass using a drop and annealed at 473 K for 24 h. The average thickness of HEO on FTO glass was about 0.04 mm, which was estimated by measuring the thickness of glass before and after deposition of HEO using a micrometer with 0.01 mm accuracy. Photocurrent generation was examined by an electrochemical analyzer in the potentiostatic amperometry mode during time (30 s light ON and 60 s light OFF), while the counter electrode was Pt wire, the reference electrode was Ag/AgCl, and the external potential was 0.7 V vs. Ag/AgCl.

4.2.4. Photocatalytic test

Photocatalytic CO_2 conversion was conducted using the powder of HEO in a continuous flow quartz photoreactor. The photoreactor, as shown in Fig. 4.2a, had a cylindrical shape with a total inner volume of 858 mL. The reactor had an inner space to insert the light source. There were two holes on the top of photoreactor: one for the inlet of CO_2 flow, which was connected to a gas cylinder; and another one for the outlet of gas and sampling the reaction products for analysis,

which was connected to a vent and gas chromatograph. For the photocatalytic reaction, 120 mg of HEO was mixed with 500 mL of deionized water and NaHCO₃ with 1 M concentration and then bubbled with CO₂ with a flow rate of 3 mL/min. The temperature was controlled as 288 K using a water chiller and the suspension was continuously stirred using a magnetic stirrer. The process was first conducted for 2 h without light irradiation, and after confirmation that no reaction products appear, the photocatalytic test was conducted under irradiation with a high-pressure Hg light source (Sen Lights Corporation, HL400BH-8, 400 W, with the spectral composition shown in Fig. 4.2b). The light intensity irradiated on the photocatalysts was 0.5 W/cm² and no filter was used during the irradiation. The reaction products were analyzed by a gas chromatograph (Shimadzu GC-8A, Ar Carrier). A flame ionization detector equipped with a methanizer (Shimadzu MTN-1) was used to measure the CO and CH₄ production rate. A thermal conductivity detector also was utilized to evaluate the H₂ and O₂ production. To be sure about the absence of CO from other sources such as contamination, blank tests were conducted (i) under irradiation in the presence of CO₂, NaHCO₃ and H₂O and without the photocatalyst addition and (ii) under irradiation in the presence of Ar, NaHCO₃ and H₂O and with photocatalyst addition.

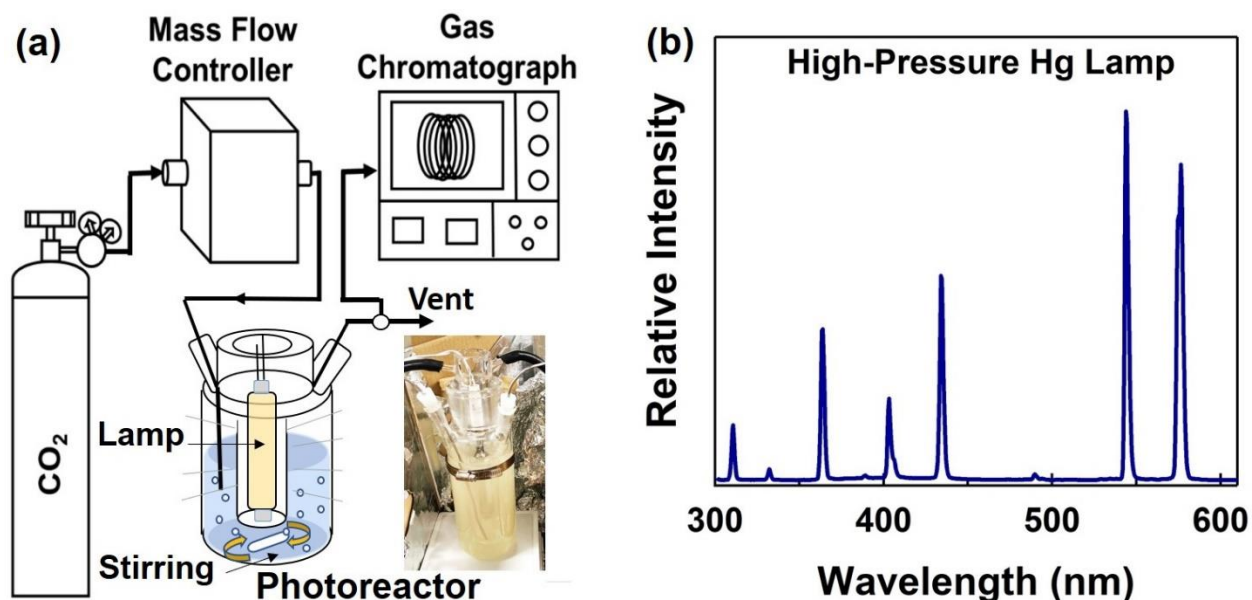


Fig. 4.2. (a) Description of experimental setting for photocatalytic CO₂ conversion including photograph of photoreactor, and (b) spectroscopy of light source used for photocatalytic test.

4.3. Results

4.3.1. Crystal structure and microstructure

Fig. 4.3 shows the SEM images of HEO in various scales. Particle size measured by SEM is 25 μm . The HEO contains particles with different sizes, as shown in Fig. 4.3, and its specific surface area, achieved by the BET method, is 0.66 m^2/g . Although big size of some particles can have negative effect on photocatalytic activity due to decreasing the active surface area, this issue can be addressed in the future by using other synthesis method or advanced crushing techniques. The presence of numerous nanograins in each particle is obvious in higher magnification images in Fig. 4.3b, c and d. The average grain size for this material is estimated to be 192 nm, while some pores are also visible within the particles. Here, it should be noted that low specific surface area and small grain size are characteristics of materials which are synthesized/processed by the HPT method [22-25].

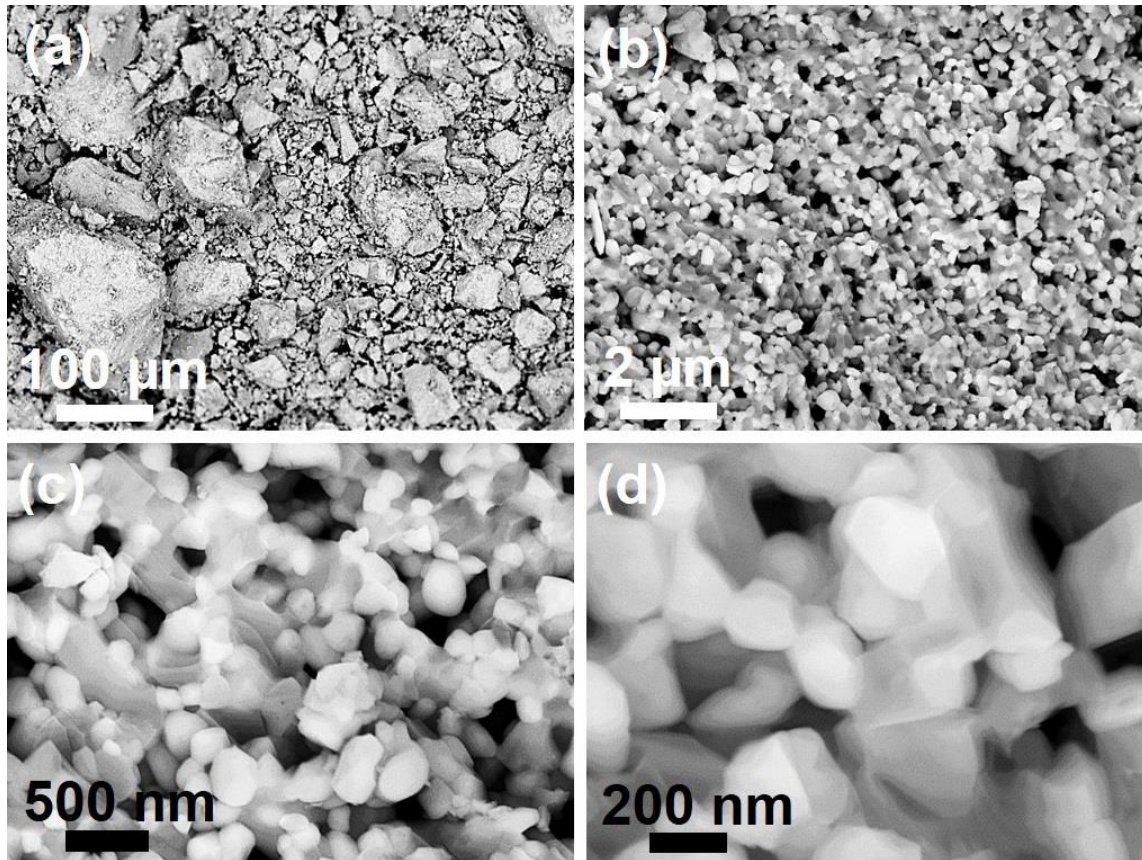


Fig. 4.3. Morphology of high-entropy oxide examined by SEM at different magnifications.

To confirm the successful oxidation of the material, electronic states of each element in the HEO are presented in Fig. 4.4 using the XPS analysis and corresponding peak deconvolution. Fig. 4.4 shows that the main cations in the sample are Ti^{4+} , Zr^{4+} , Nb^{5+} , Hf^{4+} and Ta^{5+} , suggesting that the material is successfully oxidized to a d^0 electronic configuration during the high-temperature oxidation [26]. However, it should be noted that the peaks for Ti, Zr, Nb, Hf and Ta have some shoulders to the lower energy sides, suggesting that some oxygen-deficient regions with lower oxidation states should exist within the material, as confirmed by the peak deconvolution analysis (i.e., some oxygen vacancies present). The presence of vacancies is not surprising as similar issue can be observed in other HPT-processed materials due to the strain effect [24,25] and in other HEOs due to the atomic size mismatch effect [1,2].

To confirm the distribution of elements in the material, Fig. 4.5a and Fig. 4.5b illustrate the elemental distribution mappings in the micrometer and nanometer scales, respectively. Fig. 4.5 shows that the elements distribute appropriately in both micrometer and nanometer scales. It is confirmed that the elements are successfully mixed by high-pressure mechanical alloying and their distribution remains reasonably homogeneous even after high-temperature oxidation. SEM-EDS analysis suggests that the material should have a general composition of TiZrNbHfTaO_{11} . Uniform distribution of elements is a general requirement of high-entropy materials [1-4,9-23].

Crystal structure of HEO was examined using XRD analysis, as shown in Fig. 4.6a. The material contains two phases with the monoclinic and orthorhombic structures. Based on the Rietveld analysis, the HEO consists of 40 wt% of monoclinic phase ($A2/m$ space group, $a = 1.193$ nm, $b = 0.381$ nm, $c = 2.044$ nm, $\alpha = 90^\circ$, $\beta = 120.16^\circ$, $\gamma = 90^\circ$) and 60 wt% of orthorhombic phase ($Ima2$ space group, $a = 4.092$ nm, $b = 0.493$ nm, $c = 0.527$ nm, $\alpha = \beta = \gamma = 90^\circ$). Raman spectra, shown in Fig. 4.6b from three different positions, illustrate similar patterns in different positions, suggesting the size of phases should be smaller than the spatial resolution of micro-Raman. Taken altogether, a combination of XPS, EDS and XRD confirms that a dual-phase HEO could be successfully produced in this study.

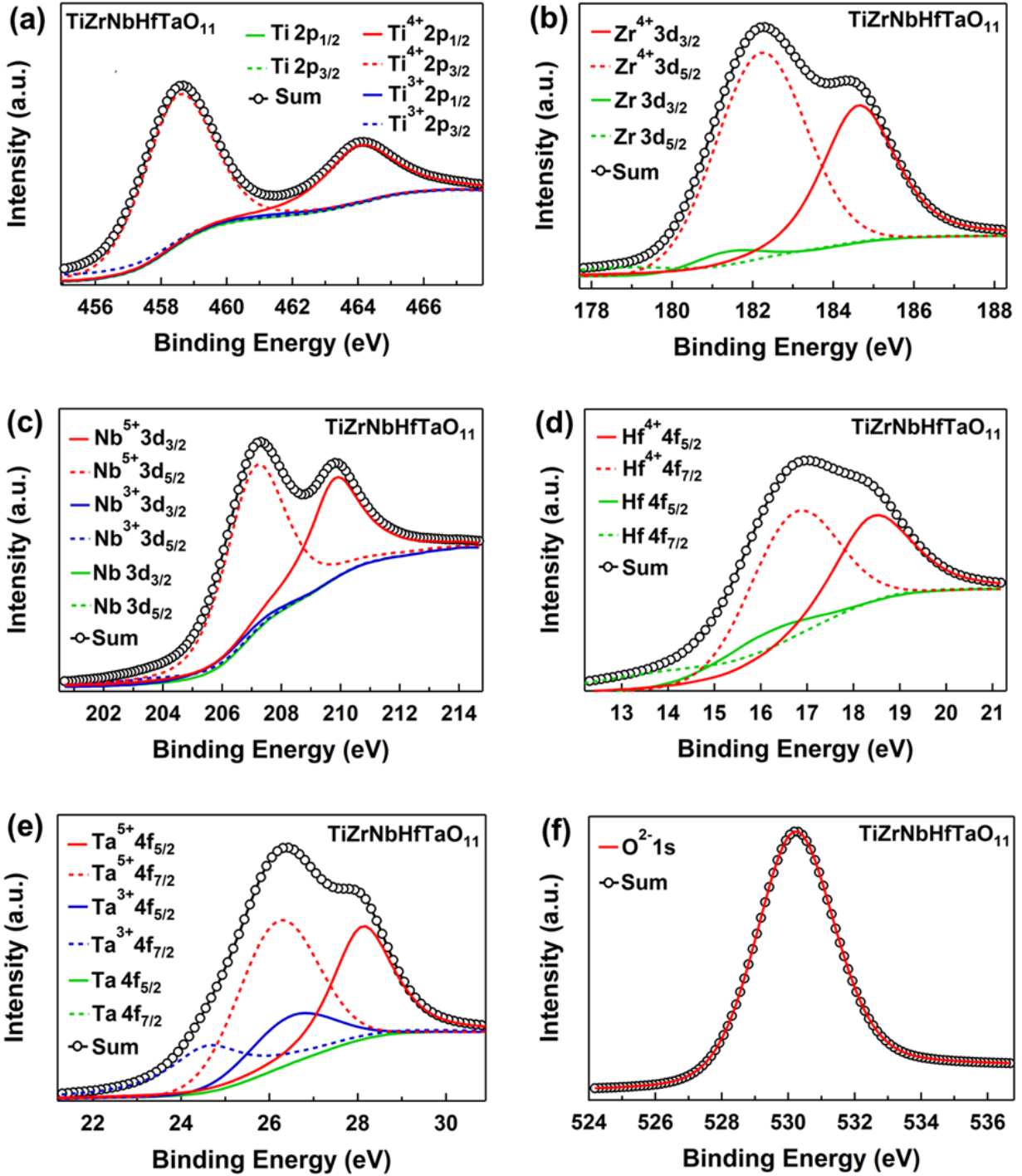


Fig. 4.4. Electronic states and relevant peak deconvolution of (a) Ti, (b) Zr, (c) Nb, (d) Hf, (e) Ta and (f) O in high-entropy oxide examined by XPS analysis.

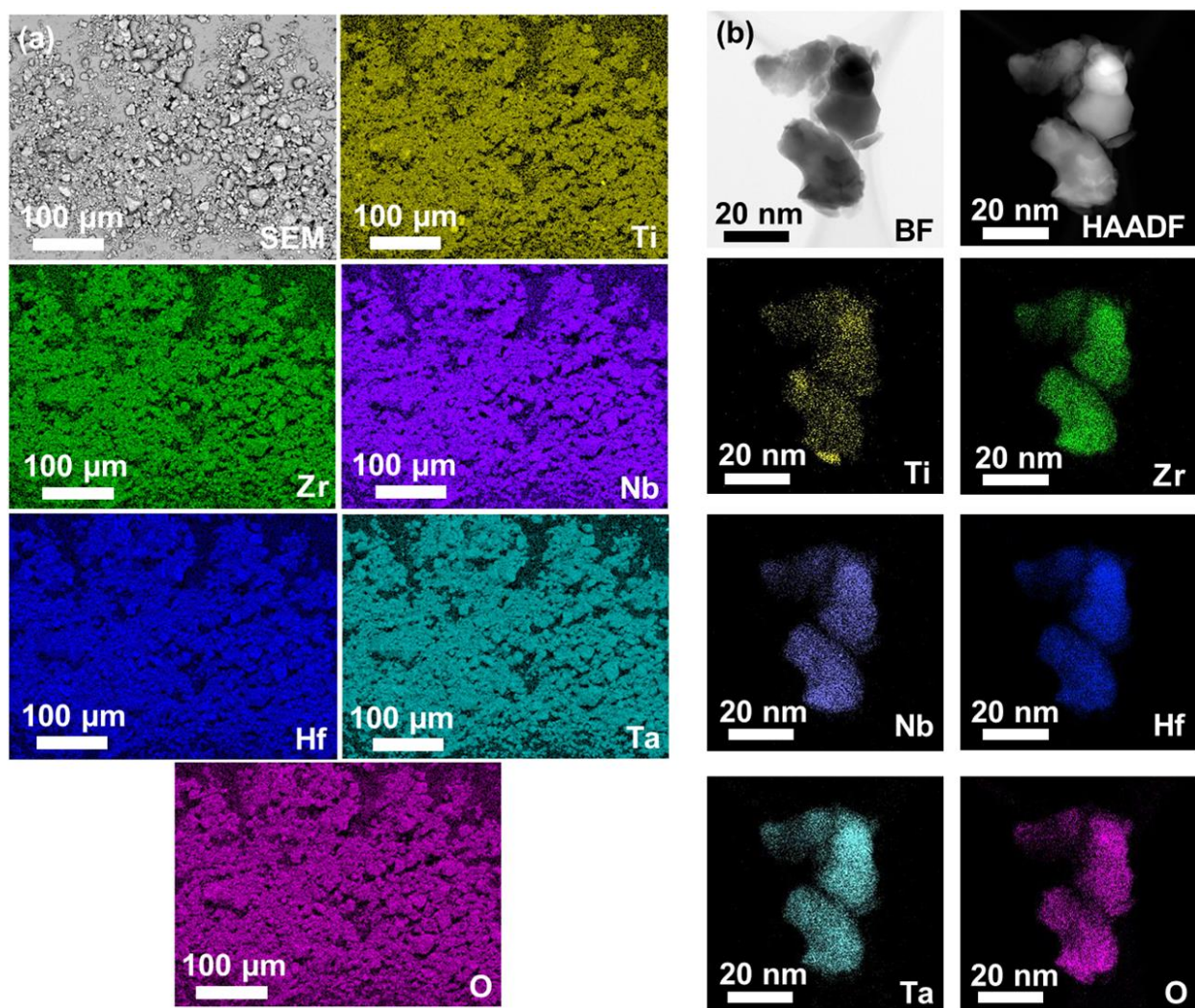


Fig. 4.5. Distribution of elements in high-entropy oxide examined at (a) micrometer scale using SEM-EDS and (b) nanometer scale using STEM-EDS.

Examination of microstructural/nanostructural features of this dual-phase HEO is shown in **Fig. 4.7**, where **a** is a BF image, **b** is a corresponding SAED pattern, **c** is a DF image, **d** and **e** are HR images, and **f** is a magnified lattice image of the selected squared region in **e**. **Fig. 4.7** reveals several important points. (i) A ring pattern of SAED image confirms the presence of many nanocrystals with random orientation in **Fig. 4.7a**. (ii) The BF and DF images confirm that the grain sizes are quite small and less than 100 nm. This indicates that there are still smaller crystals within the grain-like regions observed in the SEM images of **Fig. 4.3**. (iii) The HR images confirm the co-existence of two monoclinic and orthorhombic phases at the nanoscale and large fraction of

interphase boundaries. It was shown that the presence of interphases as charge heterojunctions can improve the photocatalytic activity through enhanced charge carrier separation and mobility [28]. (iv) The lattice images are quite distorted and close examination of the lattice confirms the presence of many dislocation defects within the grains. Since it was reported that the dislocations can enhance the light absorbance and photocatalytic activity at least in some semiconductors [29], the presence of dislocations in this HEO may positively act for enhancement of photocatalytic activity.

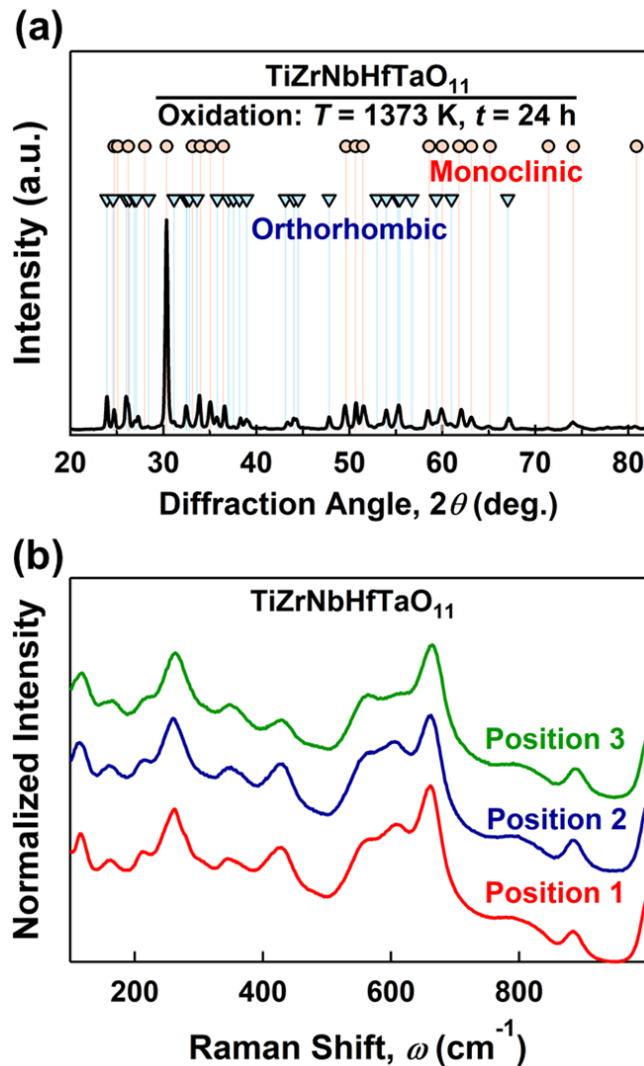


Fig. 4.6. Dual-phase structure of high-entropy oxide examined by (a) XRD profile and (b) micro-Raman spectra at three different positions.

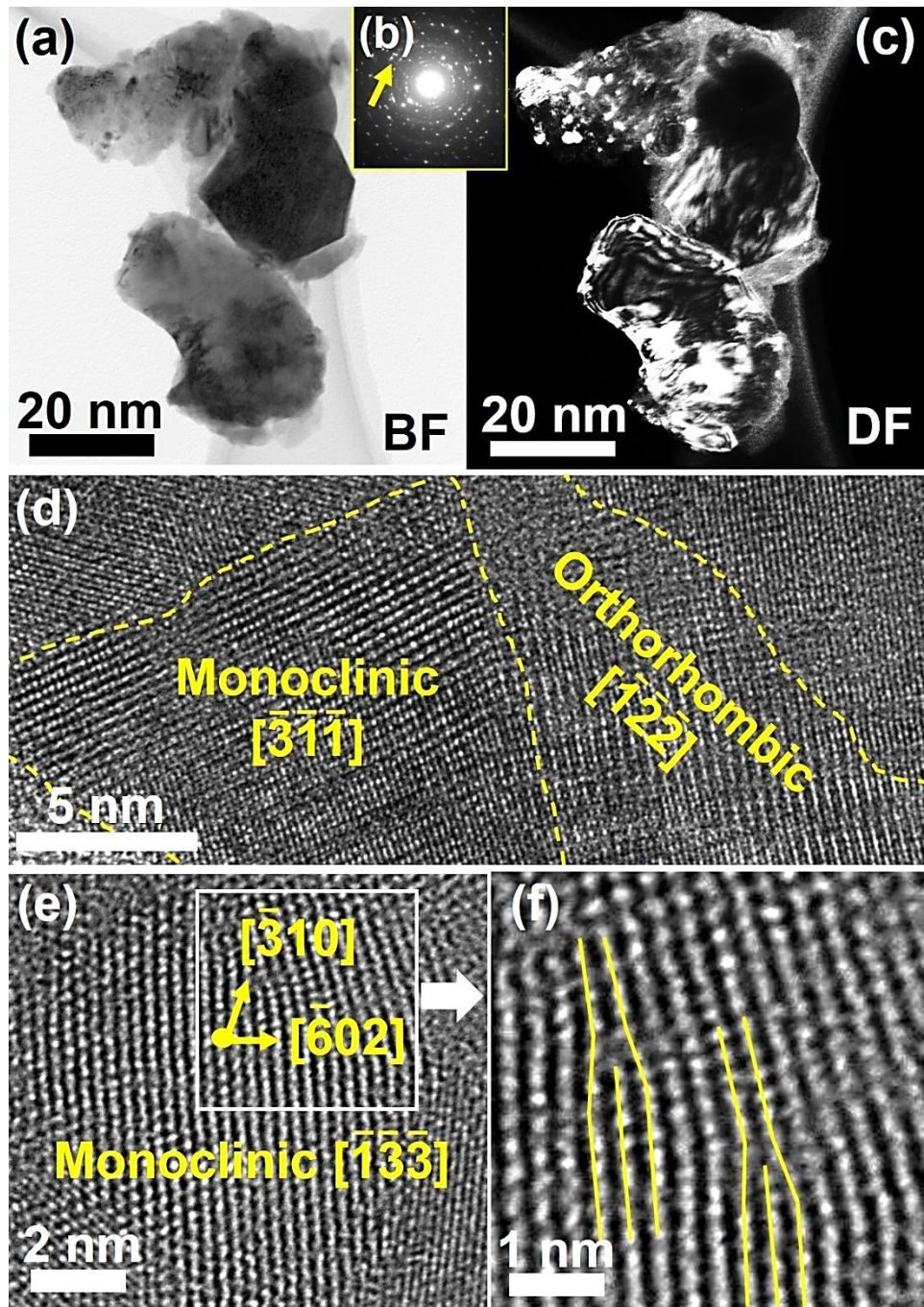


Fig. 4.7. Presence of nanoscaled dual phases with large fraction of interfaces and dislocations in high-entropy oxide examined by TEM (a) BF image, (b) SAED analysis, (c) DF image and (d-f) HR images, where (c) was taken with diffracted beams indicated by arrow in (a), and (f) is a magnified view of squared region in (e).

3.2. Electronic structure and defect states

Fig. 4.8 shows (a) UV-vis absorbance spectrum, (b) Kubelka-Munk plot, (c) XPS spectrum of top of valence band and (d) electronic structure determined by a combination of UV-vis and XPS analyses. Fig. 8a indicates that the HEO can absorb light in both ultraviolet and visible light regions, although the quantity of absorbed light in the ultraviolet region is higher than that in the visible light region. Such a visible light absorbance is not detected in binary oxides such as TiO_2 , ZrO_2 , HfO_2 , Nb_2O_5 and Ta_2O_5 [27,29].

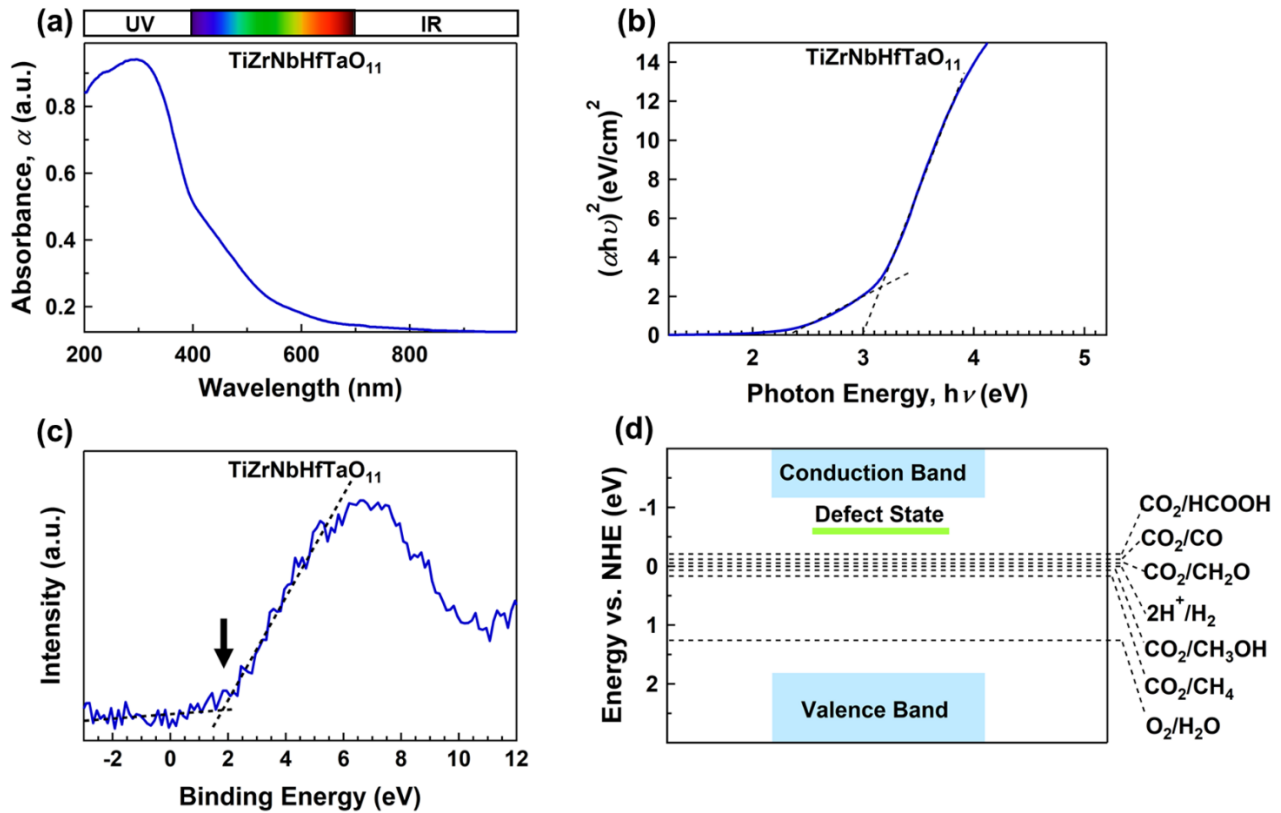


Fig. 4.8. Appropriate electronic structure of high-entropy oxide for photocatalytic CO_2 conversion. (a) UV-vis light absorbance spectrum, (b) Kubelka-Munk plot to calculate indirect bandgap (α : light absorption, h : Planck's constant, ν : light frequency), (c) XPS spectrum to estimate top of valence band, and (d) electronic band structure in comparison with potentials for photocatalytic CO_2 conversion.

Based on the Kubelka-Munk analysis, there are two apparent bandgaps of 3.0 and 2.3 eV for this HEO. The first energy gap should be related to the energy difference between the valence band and conduction band which is reasonably similar to the bandgap of TiO₂ and smaller than the bandgap of other binary oxides in the Ti-Zr-Hf-Nb-Ta-O system (3.1-5.7 eV) [30,31], and the second gap should be due to the defect level between the valence band and conduction band. The presence of defects (i.e., oxygen vacancies or color centers), which can be confirmed from the low energy shoulders in XPS spectra of cations, should be a main reason for the orange color of sample. The top of valence band calculated by XPS is 1.8 eV vs. NHE, which is shown by an arrow in Fig. 4.8c. The bottom of conduction band is calculated as -1.2 vs. NHE by considering an indirect bandgap of 3.0 eV and the defect state is estimated as -0.5 eV vs. NHE. As summarized in Fig. 4.8d, the potential of reactions for CO₂ conversion and water splitting (see Table 4.1) are between the energy levels for the top of valence band and the bottom of conduction band, and thus, this HEO can basically satisfy the requirements for photocatalytic reactions [28,32,33].

4.3.3. Charge-carrier dynamics

Charge-carrier dynamics were examined by (a) steady-state PL spectroscopy, (b) PL decay spectroscopy, (c) EPR spectroscopy and (d) photocurrent measurement, as shown in Fig. 4.9. The PL spectrum in Fig. 4.9a shows a peak at 580 nm which is equivalent to an energy level of 2.14 eV. Since this energy level is close to the energy gap of 2.3 eV, calculated using the Kubelka-Munk analysis for the defect states, it can be concluded that this peak corresponds to the recombination of excited electrons at the defect state. To have an insight into the significance of these electron-hole recombination, Table 4.2 compares the PL intensity and PL wavelength of the HEO with those measured by the current authors for anatase TiO₂ and BiVO₄ (as two popular photocatalysts for CO₂ conversion [34-36]). It is obvious that the PL intensity of HEO is lower than that of anatase TiO₂ and BiVO₂, despite its high light absorbance which is an indication of large electron-hole production. The lower PL intensity suggests that the recombination in this HEO is not higher than TiO₂ and BiVO₂, provided that the heat energy generation through the electron-hole recombination is considered identical for the three oxides.

Evaluation of PL decay intensity versus time, as shown in Fig. 4.9b, indicates that the PL decay of the HEO follows an exponential equation.

$$I(t) = A_1 \exp\left(-\frac{t}{\tau_1}\right) + A_2 \exp\left(-\frac{t}{\tau_2}\right) \quad (4.1)$$

where, $I(t)$, A_1 , A_2 , τ_1 and τ_2 are PL decay intensity at time t , amplitude of the first exponential function, amplitude of the second exponential function, fast decay time and slow decay time, respectively. Analysis of data in Fig. 4.9b suggests the values of 1.53 and 10.39 ns for τ_1 and τ_2 , respectively. Here, the following equation can be used to estimate the average lifetime, τ_{ave} [37].

$$\tau_{ave} = \frac{A_1\tau_1^2 + A_2\tau_2^2}{A_1\tau_1 + A_2\tau_2} \quad (4.2)$$

Table 4.2 compares the average lifetime for the HEO with those for anatase TiO₂ and BiVO₄. The average lifetime for the HEO is 10.5 ns which is close to the lifetime of TiO₂ anatase (10.7 ns). Low recombination intensity of the HEO, measured by steady-state PL spectroscopy, and an appropriate electron lifetime close to TiO₂ anatase, show that the excited electrons on the surface of this material can be active for appropriate time to take part in photocatalytic reaction before recombination with holes. One reason for the appropriate charge carrier lifetime and low-intensity recombination for this HEO can be the presence of oxygen vacancies on the surface [38,39].

Presence of oxygen vacancies, which was suggested by the orange color of sample in Fig. 4.1c, XPS spectroscopy in Fig. 4.4 and UV-vis spectroscopy in Fig. 4.8a, was examined further by EPR spectroscopy, as shown in Fig. 4.9c. Two symmetric hump peaks with a g factor of 2.15 appear which may be due to the oxygen vacancies, as reported in some oxides such as Nb₂O₅ [40]. It should be noted that the oxygen vacancies on the surface can act as active sites for electron-hole separation and photocatalytic reaction [36]. Moreover, it was shown that the surface oxygen vacancies have a direct effect on photocatalytic CO production rate: surface oxygen vacancies can absorb CO₂ and contribute to breaking the C=O bonds to produce CO [36].

Table 4.2. PL wavelength and intensity, fitted parameters of PL decay spectra and photocurrent density for high-entropy oxide in comparison with anatase TiO₂ and BiVO₄ photocatalysts.

PL	Wavelength (nm)		Intensity (cps)		
TiZrNbHfTaO ₁₁	580		190		
Anatase TiO ₂	510		12300		
BiVO ₄	640		300		
PL Decay	τ_1 (ns)	τ_2 (ns)	A_1	A_2	τ_{ave} (ns)
TiZrNbHfTaO ₁₁	1.53	10.39	42.34	57.66	10.5
Anatase TiO ₂	1.24	11.46	41.98	58.02	10.7
BiVO ₄	2.17	14.90	56.56	43.44	12.9
Photocurrent (mA/m ²)	Cycle 1	Cycle 2	Cycle 3	Cycle 4	
TiZrNbHfTaO ₁₁	9.6	8.9	8.4	8.2	
Anatase TiO ₂	43.5	32.2	27.9	25.2	
BiVO ₄	18.6	17.1	17.0	16.7	

Fig. 4.9d shows photocurrent measurement on HEO thin film. The material successfully generates photocurrent, although its photocurrent density decreases during the time due to the accumulation of holes with positive charge on the surface. **Table 4.2** compares the photocurrent density of the HEO with that of reference anatase TiO₂ and BiVO₄ for the first four cycles. It should be noted that the quantitative comparison of the photocurrent density of these three materials should be conducted by care due to the technical limits in making dense films with good FTO-oxide bonding by annealing at 473 K. The photocurrent density of HEO is apparently lower than that of the reference oxides. Despite the low photocurrent density of HEO, photocurrent generation on this material indicates that the excited electrons can have enough lifetime to separate from the surface of material and take part in the photocurrent generation. The generation of photocurrent is a positive sign for possible photocatalytic activity of this HEO, as discussed earlier for other photocatalysts [27].

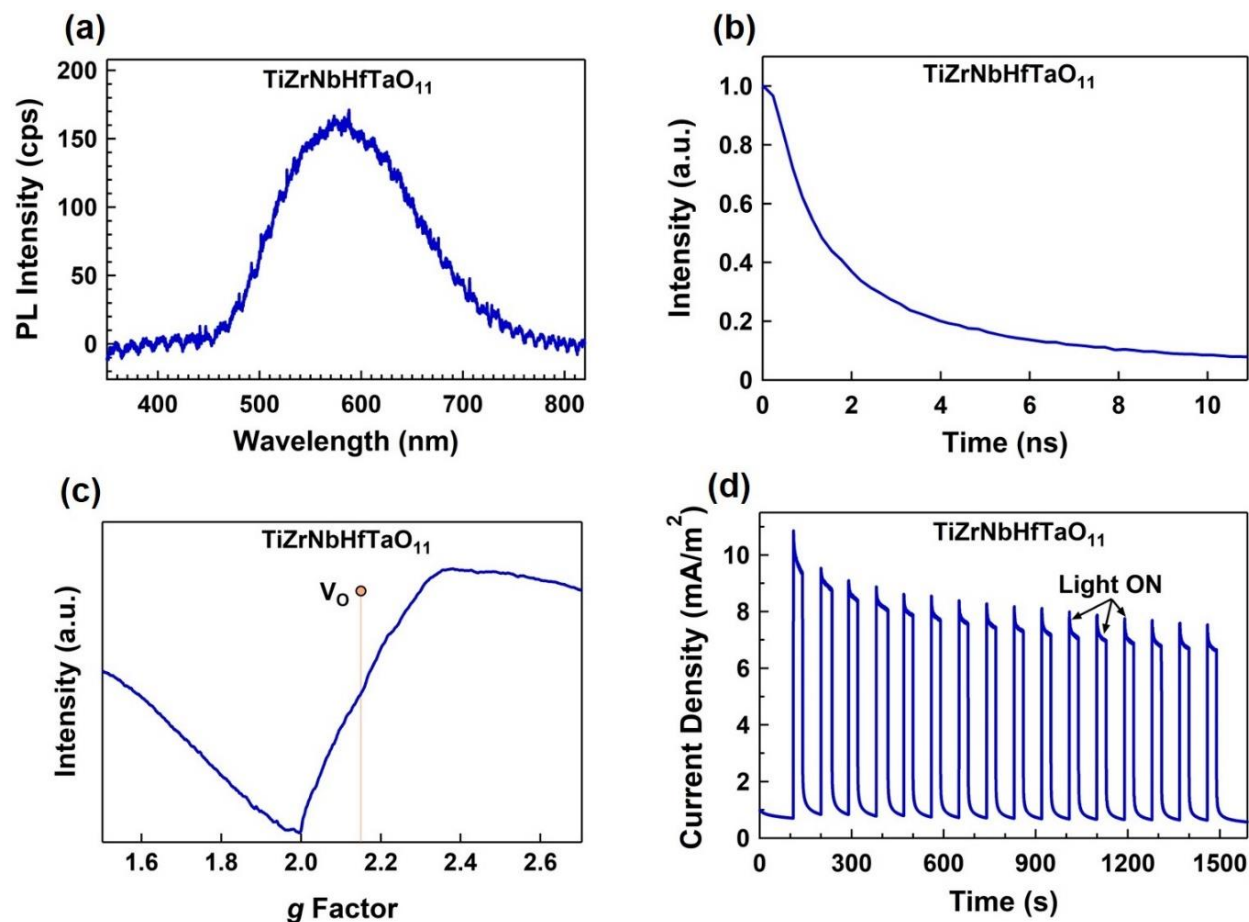


Fig. 4.9. (a) Steady-state PL emission, (b) time-resolved PL decay, (c) EPR spectra and (d) photocurrent generation for high-entropy oxide.

4.3.4. Photocatalytic activity

Photocatalytic activity of HEO for CO₂ conversion is summarized in Fig. 4.10. As shown in Fig. 4.10a and b, the HEO could successfully produce both CO and H₂ under the full arc emission of high-pressure Hg lamp without any co-catalyst addition, despite its low specific surface area as 0.66 m²/g (the error bar of gas amount measurement for three repeated tests was lower than 10%). Independent synthesis of the HEO material and repeating the photocatalytic test, as indicated as Sample #2 in Fig. 4.10a, also confirm the high activity of this material for photocatalytic CO₂ conversion with a reasonably constant CO and H₂ production rate within an extended irradiation time of 10 h. The amount of CO production is higher and the amount of H₂ production is lower for Sample #2 compared to Sample #1, suggesting that the activity of this HEO can be still improved by modification of the synthesis method. Two points should be noted here.

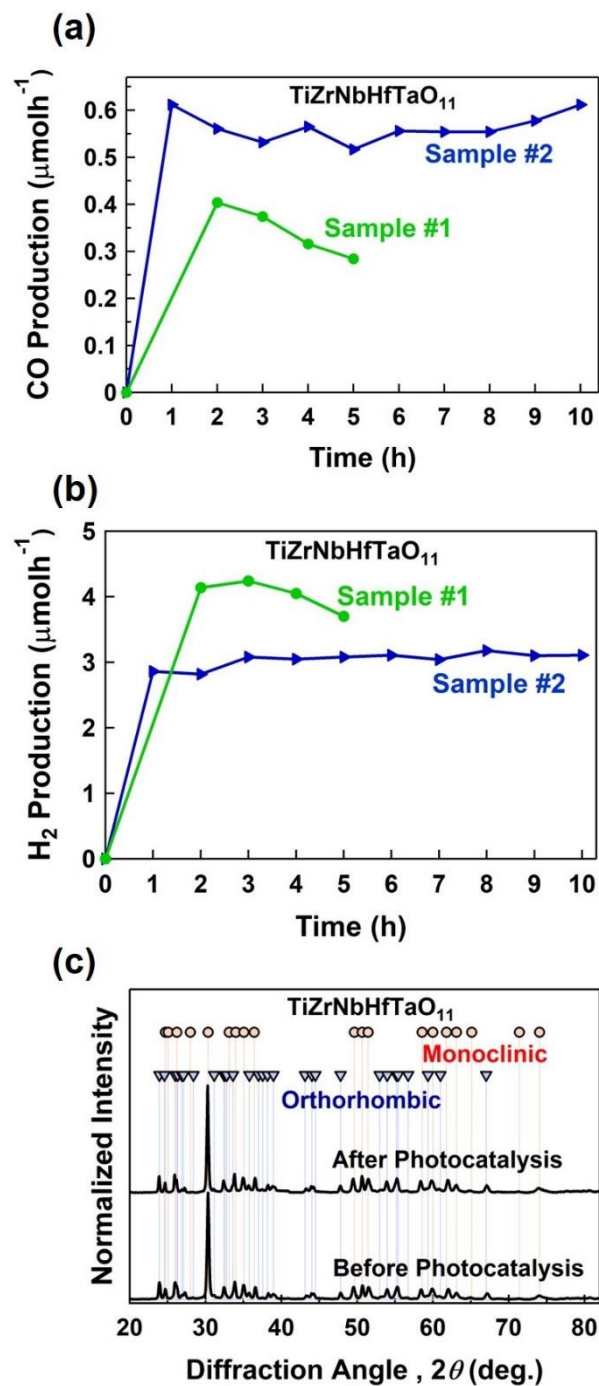


Fig. 4.10. Photocatalytic activity of high-entropy oxide for CO_2 conversion and H_2O decomposition. (a) CO production rate versus time, (b) H_2 production rate versus time, and (c) XRD pattern before and after photocatalytic test.

First, CO and H₂ were the only reaction products within the detection limits of analyses and no other products including methane could be detected. Second, blank tests confirmed that no CO and H₂ are produced by (i) CO₂ injection in the presence of HEO under the dark condition for 2 h, (ii) Ar injection in the presence of HEO under the light irradiation for 1 h, and (iii) CO₂ injection without the presence of HEO under the light irradiation for 5 h. The stability of HEO, examined by XRD analysis after the photocatalytic test, is shown in Fig. 4.10c, indicating that the crystal structure of the HEO is stable after photocatalytic test. The stability of HEOs, which was also reported for other applications such thermal barrier coatings [9,10], magnetic components [11,12], dielectric components [13,14], Li-ion batteries [15,16], Li-S batteries [17], Zn-air batteries [18], catalysts [19,20] and electrocatalysts [21], is usually due to their low Gibbs free energy resulting from their high entropy [1,2].

4.4. Discussion

Three issues need to be discussed in detail here: (i) comparison of photocatalytic activity of the HEO with available photocatalysts, (ii) factors influencing the photocatalytic activity of the HEO, and (iii) mechanism of CO₂ conversion on the HEO photocatalyst.

Although the current results confirm the potential of HEOs as a new family of photocatalysts for CO₂ conversion, their activity should be compared with other photocatalysts to have an insight into their significance. To understand this issue, photocatalytic CO₂ conversion activity of the HEO with a specific surface area of 0.66 m²/g was compared with anatase TiO₂ (99.8%), BiVO₄ (99.9%) and P25 TiO₂ (99.5%) with the surface areas of 10.2, 0.3 and 38.7 m²/g, respectively. Since various parameter such as catalyst concentration, temperature, reactor type, light source type and light intensity can influence the CO production rate, photocatalytic activity of these materials were compared in the same conditions. Fig. 4.11 shows the activity of these materials per 1 g of catalyst. The CO production rate for HEO is significantly higher than anatase TiO₂ and BiVO₄ which are some of the most popular photocatalysts for photocatalytic CO₂ conversion. Moreover, the CO and H₂ production rate on this HEO is comparable with P25 TiO₂ as a benchmark photocatalyst, although the surface area of current HEO is 60 times smaller than that of P25 TiO₂. It should be noted that the quantity of H₂ production on anatase TiO₂ and BiVO₄ was not within the detection limits of gas chromatograph. To get more insight on the significance of photocatalytic CO₂ conversion on this HEO, its activity was compared with some reported data

in the literature [5,41,42-50]. Although the experiments in the literature are not conducted under a consistent and standard condition, it is still useful to have a comparison. As given in **Table 4.3**, the amount of CO production rate varies in a wide range of 0.12-10.16 $\mu\text{molh}^{-1}\text{g}^{-1}$. The average CO production for HEO is $4.64\pm 0.30\ \mu\text{molh}^{-1}\text{g}^{-1}$ which is higher than many of the reported values in **Table 4.3**.

The reason for high CO production rate on current HEO can be attributed to various factors: the presence of lattice defects such as oxygen vacancies which can act as activation sites [5,6], the presence of five cations which can enhance the activity by straining effect [7,8], appropriate electronic structure which satisfy most of the reactions for CO₂ conversion and water splitting [32,33], and appropriate lifetime of charge carriers to participate in photocatalytic reaction due to the defective nature of HEOs [38,39]. Moreover, the presence of two phases can improve the charge carrier separation through interfaces and enhance the photocatalytic activity [27,28]. The presence of several cations in the HEO can also produce hybridized orbitals with higher activity for chemical reactions [1,2]. To further enhance the efficiency of current HEO for photocatalytic CO₂ conversion, future works are required to enhance its specific surface area by improving the synthesis or crushing techniques.

Table 4.3. Summary of some reported photocatalytic CO₂ conversion rates in literature in comparison with results of current study.

Photocatalyst	Light Source	CO Production ($\mu\text{molh}^{-1}\text{g}^{-1}$)	References
TiO ₂ Nanosheet -CN	150 W Xe lamp	2.04	[41]
TiO ₂ - Graphitic carbon	300 W Xe lamp	10.16	[43]
TiO ₂ nanosheets exposed {001} facet	2 *18W Hg lamps	0.12	[44]
TiO ₂ - Hydrogenated CoO _x	150W UV lamp	1.24	[45]
TiO ₂ 3D Ordered Microporous - Pd	300 W Xe lamp	3.9	[46]
C ₃ N ₄ by Thermal Condensation	350 W Xe lamp	4.83	[5]
Cd _{1-x} Zn _x S	100 W LED plate	2.9	[47]
BiOI	300 W Xe lamp	4.1	[48]
xCu ₂ O-Zn _{2-2x} Cr	200-W Hg-Xe lamp	2.5	[42]
CeO _{2-x}	300 W Xe lamp	1.65	[49]
Cu ₂ O/RuO _x	150 W Xe lamp	0.88	[50]
TiO ₂ Anatase	400 W Hg Lamp	0.58±0.12	This Work
BiVO ₄	400 W Hg Lamp	2.16±0.21	This Work
TiO ₂ P25	400 W Hg Lamp	4.63±0.33	This Work
TiZrNbHfTaO ₁₁	400 W Hg Lamp	4.64±0.30	This Work

Regarding the third issue, three main mechanisms for photocatalytic CO₂ reduction have been suggested, as summarized in Table 4.4: carbene pathway, formaldehyde pathway and glyoxal pathway [35,36]. The behavior of current HEO is similar to P25 TiO₂, suggesting that both materials probably follow the same pathway. Although even for TiO₂ with different impurities and lattice defects, there are still significant arguments regarding the CO₂ reduction pathways, it is still possible to discuss about the possible mechanisms for current HEO photocatalyst. The nonappearance of HCOOH, CH₃OH and CH₄ in the gas and liquid phases within the detection limits of analyses suggests that the formaldehyde pathway may not be the major mechanism [35,36]. The nonappearance of HCOOH and CH₄ also indicates that the glyoxal pathway may not be the major mechanism [35,36].

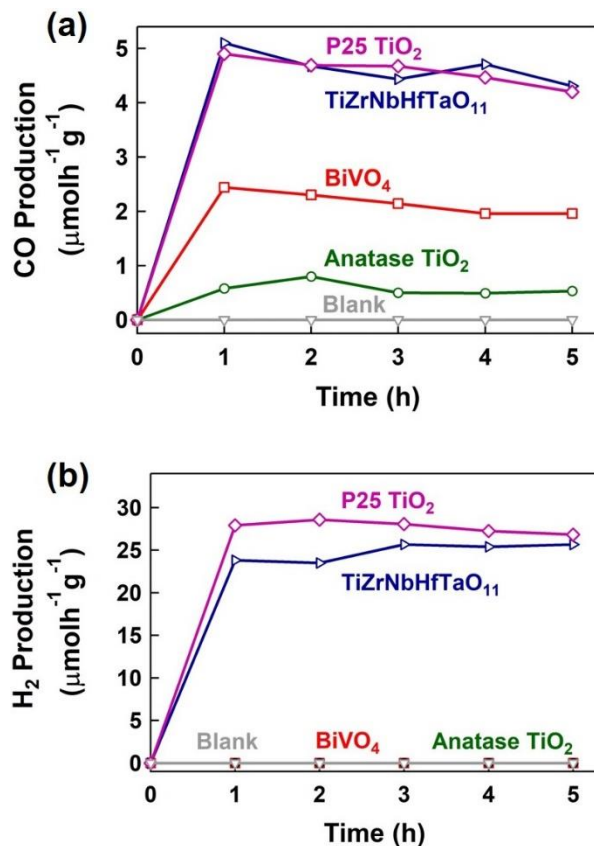


Fig. 4.11. High efficiency of high-entropy oxide compared with TiO₂ and BiVO₄ for photocatalytic CO₂ conversion. (a) CO production rate and (b) H₂ production rate versus time.

The production of CO suggests that the carbene pathway is probably the major mechanism. However, the absence of CH₄ and the presence of H₂, which is similar to the behavior of P25 TiO₂

in this study, indicates that the carbene pathway possibly stops at some intermediate stages due to the formation of H₂ gas [51]. The absence of CH₄ can also be explained by the defective structure of HEO. Since the HEO material has oxygen vacancies as surface defects, CO₂ in connection with H₂O as a Lewis acid tends to adsorb on oxygen vacancies [51]. This adsorption degrades C=O bonding and produce ^{*}CO radicals and consequently generates CO gas [36]. Compared with CO₂, the generated CO has lower tendency to be adsorbed on the surface defects [51], and thus, the carbene pathway does not continue to produce detectable quantity of CH₄. For TiO₂, it was also reported that although the CH₄ formation in the carbene pathway is thermodynamically more favorable than CO and H₂ formation, the formation of CH₄ is kinetically more difficult because it needs higher numbers of electrons and protons [52].

Table 4.4. Main mechanisms for CO₂ photocatalytic reduction pathway [35].

Carbene Pathway	Formaldehyde Pathway	Glyoxal Pathway
(1) CO ₂ + e ⁻ → CO ₂ ^{-*}	(1) CO ₂ + e ⁻ → CO ₂ ^{-*}	(1) CO ₂ + e ⁻ → CO ₂ ^{-*}
(2) CO ₂ ^{-*} + e ⁻ + H ⁺ → CO + OH ⁻	(2) CO ₂ ^{-*} + H ⁺ → [*] COOH	(2) CO ₂ ^{-*} + e ⁻ + H ⁺ → CHOO ⁻
(3) CO + e ⁻ → CO ^{-*}	(3) [*] COOH + e ⁻ + H ⁺ → HCOOH	(3) CHOO ⁻ + H ⁺ → HCOOH
(4) CO ^{-*} + e ⁻ + H ⁺ → C + OH ⁻	(4) HCOOH + e ⁻ + H ⁺ → H ₃ COO ^{-*}	(4) HCOOH + e ⁻ → HOC [*]
(5) C + e ⁻ + H ⁺ → CH [*]	(5) HCOOH ₂ [*] + e ⁻ + H ⁺ → HCOH + H ₂ O	(5) HOC [*] + OH ⁻ → C ₂ H ₂ O ₂
(6) CH [*] + e ⁻ + H ⁺ → CH ₂	(6) HCOH + e ⁻ → H ₂ C ^{-*} O	(6) C ₂ H ₂ O ₂ + e ⁻ + H ⁺ → H ₃ O ₂ C ₂ [*]
(7) CH ₂ + e ⁻ + H ⁺ → CH ₃ [*]	(7) H ₂ C ^{-*} O + H ⁺ → H ₂ OHC [*]	(7) H ₃ O ₂ C ₂ [*] + e ⁻ + H ⁺ → C ₂ H ₄ O ₂
(8) CH ₃ [*] + e ⁻ + H ⁺ → CH ₄	(8) H ₂ OHC [*] + e ⁻ + H ⁺ → CH ₃ OH	(8) C ₂ H ₄ O ₂ + e ⁻ + H ⁺ → H ₃ OC ₂ [*] + H ₂ O
(9) CH ₃ [*] + OH ⁻ → CH ₃ OH	(9) CH ₃ OH + e ⁻ + H ⁺ → [*] CH ₃ + H ₂ O	(9) H ₃ OC ₂ [*] + e ⁻ + H ⁺ → C ₂ H ₄ O
	(10) [*] CH ₃ + e ⁻ + H ⁺ → CH ₄	(10) C ₂ H ₄ O + h ⁺ → H ₃ OC ₂ [*] + H ⁺
		(11) H ₃ OC ₂ [*] → CH ₃ [*] + CO
		(12) CH ₃ [*] + e ⁻ + H ⁺ → CH ₄

Taken altogether, this study introduces HEOs as active photocatalysts for CO₂ conversion, and this opens a path to explore numerous photocatalysts by considering the state-of-art on engineering of catalysts for CO₂ photoreduction [53]. Despite high activity of current HEO, future studies are required to clarify the exact CO₂ conversion mechanism on this new family of materials. It should be noted that although the material in this study was synthesized by a two-step high-pressure mechanical alloying and high-temperature oxidation, other methods developed earlier for the synthesis of high-entropy ceramics [54] can be used in the future to synthesize powders with high specific surface area and low economical cost. Moreover, since earlier studies showed that

the semiconductor photocatalysts with CO₂ conversion capability can have good activity for degradation of organic pollutants as well [55,56], it is expected that the photocatalytic activity of HEOs is not limited to CO and H₂ production.

4.5. Conclusion

A high-entropy oxide with a general composition of TiZrNbHfTaO₁₁ was synthesized and used for photocatalytic CO₂ conversion. Due to appropriate electronic band structure, good charge carrier lifetime and a defective and strained dual-phase structure, the material acted as a photocatalyst for CO₂ to CO conversion and H₂O to H₂ production without addition of any co-catalyst. The photocatalytic activity of this oxide was better than well-known anatase TiO₂ and BiVO₄ photocatalysts and comparable with P25 TiO₂ as a benchmark photocatalyst, suggesting high-entropy oxides as a new family of photocatalysts for CO₂ conversion.

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Chapter 5. Significant CO₂ photoreduction on a high-entropy oxynitride

5.1. Introduction

Oxide photocatalysts such as TiO₂ [1,2] are the most common photocatalysts with high stability for CO₂ photoreduction application, but they have large bandgaps such as 3.1 eV for TiO₂ [1,2]. In contrast, there are some reports on photocatalytic activity of nitrides such as TaN [3] and C₃N₄ [4] for CO₂ conversion which have lower bandgaps compared to oxides, but nitrides are not chemically so stable [3,4]. To solve the problem of oxide and nitride photocatalysts in terms of large bandgap and low stability, respectively, oxynitrides were recommended as low bandgap and highly stable catalysts for photocatalysis [5]. Oxynitrides have been widely used for photocatalytic water splitting; however, only limited oxynitrides such as α -Fe₂O₃/LaTiO₂N [6] and TaON [7,8] were used for photocatalytic CO₂ conversion. Significant electron-hole recombination, sluggish kinetics, the low tendency for CO₂ adsorption and relatively low stability in the co-presence of CO₂ and water are some reasons for limited application of metal oxynitrides for photocatalytic CO₂ conversion [6,9]. Therefore, introducing a strategy to solve all or some of these problems is a key issue in using the benefits of oxynitrides for CO₂ conversion. Simultaneous addition of several principal elements and production of high-entropy oxynitride ceramics can be an effective strategy, although there have been few attempts in this regard. A combination of the perception of metal oxynitrides as low bandgap photocatalysts and high-entropy ceramics as highly stable materials can be a new strategy to expand the application of oxynitrides for photocatalytic CO₂ conversion. Although high-entropy oxynitrides (HEONs) were successfully synthesized in a few studies [10,11], there are no reports on the photocatalytic performance of HEONs for CO₂ photoreduction.

In this chapter, a two-phase TiZrNbHfTaO₆N₃ was synthesized as the first HEON for photocatalytic CO₂ conversion by high-pressure torsion mechanical alloying [12] and subsequent oxidation and nitriding. The HEON showed better light absorbance, lower charge carrier recombination rate, higher CO₂ adsorbance and larger photocatalytic CO₂ conversion compared to relevant HEO (TiZrNbHfTaO₁₁) and benchmark photocatalyst P25 TiO₂. Moreover, the activity of the HEON was higher than almost all photocatalysts developed in the literature for CO₂

photoreduction. These findings open a path to develop new high-entropy photocatalysts with significant efficiency for CO₂ photoreduction.

5.2. Experimental procedures

Despite various synthesis methods reported in the literature to produce HEOs [13-27] and HENs [13,14,28-30], the HEON was fabricated using a three-step synthesis method for this study [33]: (i) severe plastic deformation through the high-pressure torsion (HPT) method for alloying pure elemental powders [12,31], (ii) oxidation at elevated temperature and (iii) nitriding at elevated temperature. First, titanium (99.9%), zirconium (95.0%), niobium (99.9%), hafnium (99.5%) and tantalum (99.9%) powders with the same molar fraction of 0.2 were dispersed in acetone, mixed using ultrasonic and dried in air. About 700 mg of powder mixture was compacted into a 10 mm diameter disc under a pressure of 0.4 GPa and further proceeded by HPT under 6 GPa at room temperature using a rotation rate of one turn per minute for 100 turns to achieve a single-phase (body-centered cubic, BCC) alloy. Second, the HPT-processed high-entropy alloy was crushed in a mortar and inserted into a furnace for 24 h under a hot (1373 K) air atmosphere to generate the HEO, TiZrNbHfTaO₁₁ with dual monoclinic (40 wt%) and orthorhombic (60 wt%) structures. Third, the HEO was processed by nitriding in ammonia at 1373 K for 7 h using a heating rate of 20 Kmin⁻¹ with an NH₃ flow of 150 mLmin⁻¹ to generate a two-phase (40 wt% monoclinic + 60 wt% face-centered cubic, FCC) HEON, TiZrNbHfTaO₆N₃. The fabricated HEON was characterized by different methods, as follows.

The crystallographic features were analyzed by X-ray diffraction (XRD) using a Cu K α source having 0.1542 nm wavelength. Phase fractions and lattice parameters were measured by the Rietveld analysis in the PDXL software.

The composition was examined by dispersing the sample on a carbon tape and conducting energy-dispersive X-ray spectroscopy (EDS) in a scanning electron microscope (SEM) under 15 keV.

The microstructure was examined by dispersing the crushed sample on carbon grids and employing a transmission electron microscope (TEM) under 200 keV by taking high-resolution (HR) images and analyzing them by fast-Fourier transform (FFT). Moreover, the distribution of elements was examined by a scanning-transmission electron microscope (STEM) under 200 kV by taking high-angle annular dark-field (HAADF) micrographs and conducting EDS analysis.

X-ray photoelectron spectroscopy (XPS) was performed to determine the top of the valence band and the electronic state of each element using a Mg K α source.

The absorbance of light, and band structure including the level of bandgap were evaluated by UV-vis diffuse reflectance spectroscopy (followed by Kubelka-Munk calculation) and X-ray/UV photoelectron spectroscopy (XPS and UPS). The valence band top was determined by the UPS and XPS analyses and the conduction band bottom was determined by subtracting the bandgap value from the valence band top.

The electron-hole recombination was evaluated by photoluminescence (PL) spectroscopy using a UV laser (325 nm wavelength).

Photocurrent measurement on thin films of the samples was performed using the full arc of a Xe lamp in a 1 M Na₂SO₄ electrolyte. The experiments were conducted in the potentiostatic amperometry mode during the time (180 s light ON and 180 s light OFF) using an electrochemical analyzer. The counter and reference electrodes were Pt wire and Ag/AgCl, respectively, and the external potential was 0.7 V vs. Ag/AgCl. To prepare the thin films, 5 mg of each sample was crushed in 0.2 mL ethanol, spread on the FTO glass (fluorine-doped tin oxide with 2.25 mm thickness and 15 × 25 mm² surface area), and annealed at 473K for 2 h.

Diffuse reflectance infrared Fourier transform (DRIFT) spectrometry was performed to understand the adsorbance mode of CO₂ on the surface of each photocatalyst. First, 50 mg of each sample was treated at 773 K for 1 h in an argon atmosphere. Then, argon was replaced by 100% CO₂ gas at 773 K and the samples were kept under this condition for 30 min. The samples were then cooled down to room temperature and the CO₂ gas was replaced with argon. After keeping the samples in argon for 30 min, the DRIFT spectroscopy was conducted.

CO₂ photoreduction was examined in a cylindrical-shaped quartz photoreactor with 858 mL inner volume and specifications described in detail earlier [27]. Light source was placed in a space inside the photoreactor and CO₂ flow entered the reactor from a gas cylinder by a hole on the top of the reactor. Outlet gas from the reactor partly entered a gas chromatograph for the gas analysis and mainly flew to a vent. For the photoreduction experiments, 100 mg of HEON photocatalyst were dispersed in a 500 mL solution of 1 M NaHCO₃ and pure water. CO₂ gas was injected into the mixture (3 mLmin⁻¹) and the mixture was continuously stirred by a magnetic stirrer. It should be noted that the temperature was kept constant at 288 K utilizing a water chiller. To be sure about the nonappearance of reaction products without irradiating the mixture, the experiments were first

performed for 2 h in dark conditions. Then the photocatalytic experiment was performed under irradiation of a 400 W high-pressure mercury lamp (HL400BH-8 of Sen Lights Corporation) with 0.5 Wcm^{-2} light intensity without any filtration. The gas of the photoreactor was analyzed using gas chromatography (GC-8A of Shimadzu). The generation of CH_4 and CO was analyzed using a methanizer and flame-ionization detector. The generation of oxygen and hydrogen was analyzed using a thermal conductivity detector. A blank test was conducted without catalyst addition under light irradiation and CO_2 injection to confirm that no CO was produced from other sources in the experimental system. Another blank test was conducted with catalyst addition under light irradiation and argon injection to confirm that CO was not produced without CO_2 injection.

5.3. Results

Fig. 5.1a illustrates the XRD profile of HEON. The HEON has two cubic (Fm3m space group, $a=b=c=0.459 \text{ nm}$; $\alpha=\beta=\gamma=90^\circ$) and monoclinic (P21/c space group, $a=0.512$, $b=0.517$, $c=0.530 \text{ nm}$; $\alpha=\gamma=90^\circ$, $\beta=99.2^\circ$) phases with 60 and 40 wt% fractions, respectively. **Fig. 5.1b** illustrates the EDS profile of the HEON. The EDS analysis suggests a general composition of $\text{TiZrHfNbTaO}_6\text{N}_3$ for the synthesized HEON. The presence of two phases should be due to the thermodynamics of the Ti-Zr-Hf-Nb-Ta-O-N system at the synthesis temperature. The existence of two phases can be beneficial for charge separation in photocatalysis because the phase boundaries can act as heterojunctions for charge carrier migrations [32,33]. Although first-principle electronic structure calculations are required to clarify the migration direction of charge carriers in this HEON, it is expected that photoexcited electrons in the conduction band of one phase with a higher energy level move to the conduction band of the other phase and exited holes transfer from the valence band of one phase with the lower energy level to the valence band of another phase [32,33].

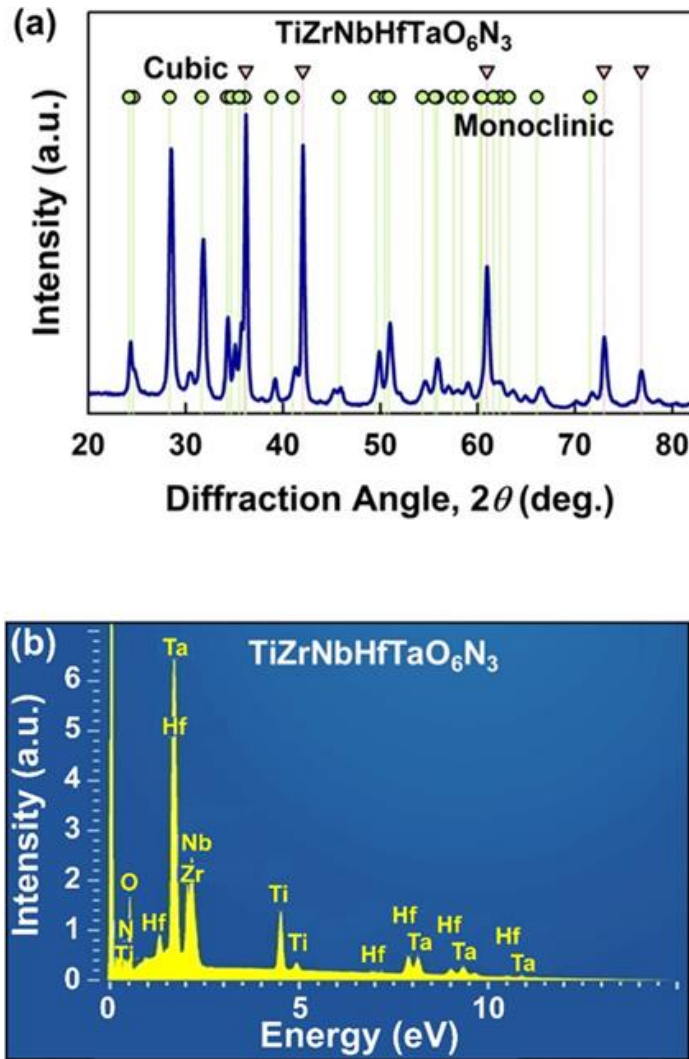


Fig. 5.1. Formation of high-entropy oxynitride with cubic and monoclinic phases and chemical composition of $\text{TiZrNbHfTaO}_6\text{N}_3$. a) XRD profile and b) EDS spectrum of high-entropy oxynitride.

The microstructure of the HEON is shown in **Fig. 5.2** using different methods. **Fig. 5.2a** illustrates a micrograph taken by SEM, which indicates that the HEON contains large powders with an average size of $20\ \mu\text{m}$. **Fig. 5.2b** shows a HR image taken by TEM confirming the existence of nanocrystals of cubic and monoclinic phases which agrees with the XRD analysis. It also indicates the existence of a large fraction of interphases that can act as heterojunctions [34]. **Fig. 5.2c** illustrates a HAADF micrograph with relevant EDS mappings taken by STEM, showing a reasonably homogenous distribution of elements at the nanometer scale. Slight differences in the

distribution of metallic elements, oxygen and nitrogen should be mainly due to the presence of two phases. Here, it should be noted that XPS analyses, shown in [Supporting Information Fig. S1](#), confirm that the main states of elements are Ti^{4+} , Zr^{4+} , Hf^{4+} , Nb^{5+} , Ta^{5+} , O^{2-} and N^{3-} .

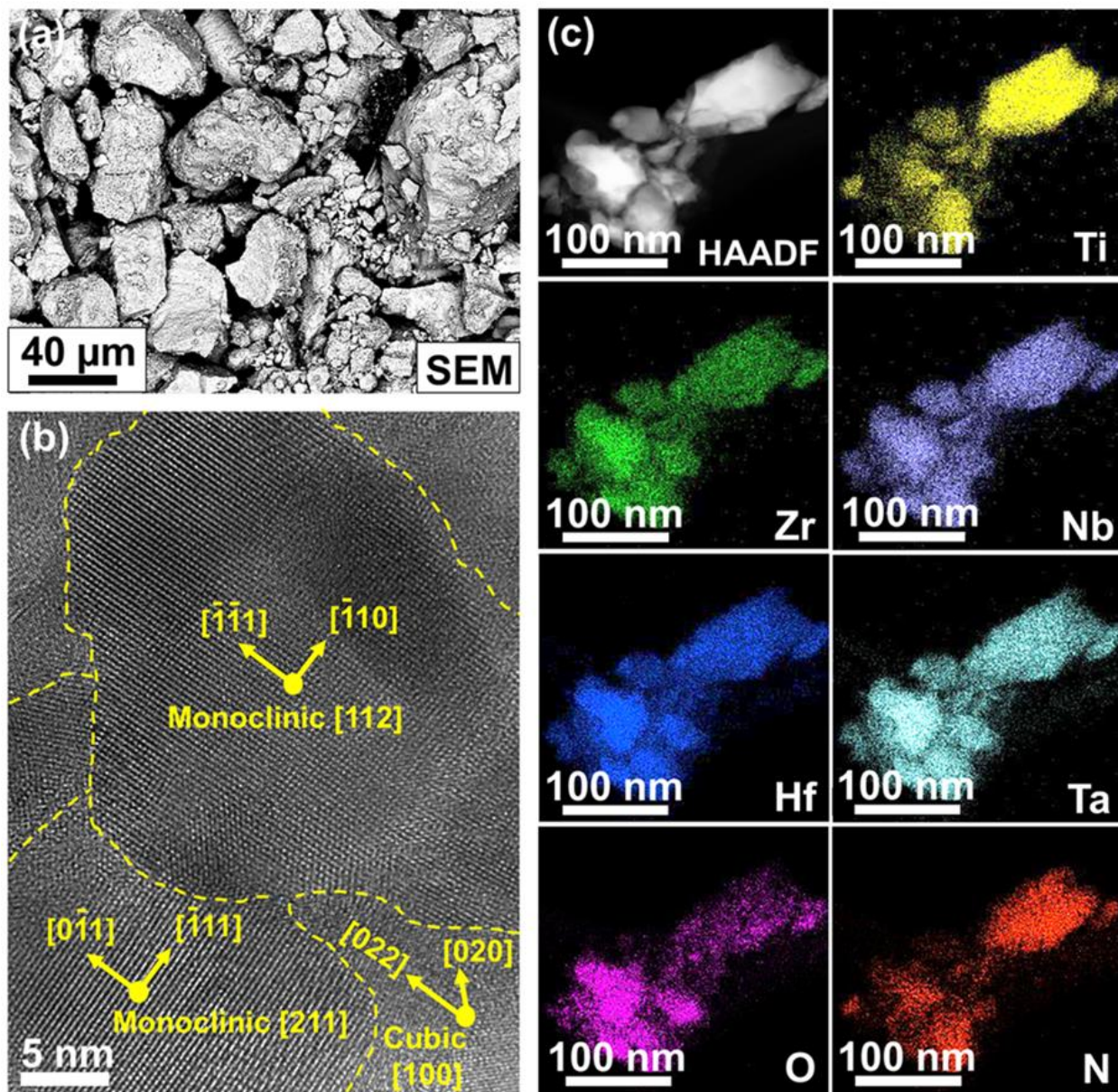


Fig. 5.2. Formation of nanocrystalline monoclinic and cubic phases with uniform elemental distribution in high-entropy oxynitride powder. a) SEM micrograph, b) HR micrograph and c) STEM-HAADF micrograph and relevant EDS mappings for high-entropy oxynitride.

Fig. 5.3a shows the light absorbance of the HEON in comparison with the relevant HEO as well as the P25 TiO₂ photocatalyst. The HEON exhibits significant light absorbance compared with the HEO and P25 TiO₂. According to the Kubelka-Munk calculation, the bandgap for the HEON is 1.6 eV which is extremely narrower compared with the bandgap of the HEO (3.0 eV) and P25 (3.1 eV) [27]. **Fig. 5.3b** shows the electronic band structure of the three mentioned materials including the appearance of three samples. A color change from white and orange for P25 TiO₂ and the HEO occurs to dark brown for the HEON, confirming the high light absorbance of the HEON in good agreement with the UV-vis absorbance data [35]. The electronic band structures were determined by considering the bandgaps calculated using the Kubelka-Munk theory, the top of the valence band was measured by XPS spectroscopy, and the bottom of the conduction band was calculated by subtracting the bandgap from the top of the conduction band. The bandgap for the P25 TiO₂, the HEO and the HEON are 3.0, 3.0 and 1.6 eV, respectively; the values for the top of the valence band are 2.2, 1.8 and 1.3 eV vs. NHE for P25 TiO₂, the HEO and the HEON, respectively; and the values for the bottom of the conduction band are -0.8, -1.2 and -0.3 eV vs. NHE for P25 TiO₂, the HEO and the HEON, respectively. As shown in **Fig. 5.3b**, the band structure of the HEON indicates its low bandgap with appropriate positions of the valence band top and the conduction band bottom for various CO₂ conversion reactions [36]. The low bandgap of this HEON can lead to easy separation of electrons and holes during photocatalysis.

Fig. 5.3c shows the photoluminescence spectra of the three materials to examine the electron-hole recombination. P25 TiO₂ and the HEO have almost the same photoluminescence intensity while the HEON shows the lowest photoluminescence. The absence of an intensive photoluminescence peak for the HEON confirms the significant suppression of electron-hole recombination which is a principal requirement for the enhancement of photocatalytic reactions [37,38]. These results show that the main problem of metal oxynitrides in terms of high electron-hole recombination [7,9] can be solved by the strategy used in this study through the concept of high-entropy ceramics.

Fig. 5.3d shows the photocurrent measurement for the three samples. Due to the different particle sizes of these three samples and their dissimilarities in making binding to the FTO glass, their current density cannot be compared quantitatively; however, the shape of their photocurrent curves can clarify their different behaviors. For P25 TiO₂ there is a spike peak at the beginning of irradiation, but the current density decreases rapidly to a steady state, suggesting that electron-hole

separation is followed by fast recombination. For the HEO the photocurrent curve under irradiation is almost a straight horizontal line which shows a better electron-hole separation of the HEO compared to P25 TiO₂. For the HEON, the current density increases by irradiation and reaches a steady state, suggesting that the ratio of electron-hole recombination to separation is the lowest for the HEON. After stopping the irradiation, the HEON still shows some reduced photocurrent due to the remained excited charge carriers, while P25 TiO₂ exhibits almost no photocurrent under the dark condition. The successful photocurrent generation on this HEON with an appropriate ratio of electron-hole separation to recombination suggests the potential of this material to act as a catalyst for CO₂ photoreduction [27,35].

Fig. 5.4a compares the CO production rate of the HEON with the relevant HEO and P25 TiO₂ benchmark photocatalyst. Photoreduction of CO₂ to CO on the HEON is considerably better than the HEO and P25. The CO production rate for the HEON reaches 14.3 μmolh⁻¹g⁻¹ after 1 h; and then decreases to a constant level after 3 h. The deviations in the reaction rate in the first two hours should be due to the time needed to reach an equilibrium condition in the measurement system. The average CO production rate for this HEON is 11.6 ± 1.5 μmolh⁻¹g⁻¹ after 5 h, while the HEO and P25 have lower photocatalytic CO production rates of 4.6 ± 0.3 μmolh⁻¹g⁻¹. Fig. 5.4b shows the photocatalytic activity of the HEON for H₂ production and compares it with the HEO and P25 TiO₂. This figure indicates that the efficiency of the HEON is much better than the HEO and P25 for photocatalytic H₂ production. The average H₂ production rate for the HEON is 5.1 ± 0.5 μmolh⁻¹g⁻¹, while the rate for the HEO and P25 TiO₂ is 1.3 ± 0.1 and 1.5 ± 0.1 μmolh⁻¹g⁻¹, respectively. To confirm the high activity and stability of this HEON for CO₂ photoreduction, a long-term photocatalytic experiment for 20 h was conducted on the sample after storage in air for 7 months. As shown with dashed-line curves in Fig. 5.4a and 5.4b, the material still shows high activity with a constant CO and H₂ production rate, although the reaction rates are slightly lower than in the first experiment.

Here, three issues regarding the photocatalytic tests should be mentioned. First, no CO was detected in three blank tests: (i) with catalyst addition and CO₂ injection under dark conditions, (ii) without catalyst addition under light irradiation and CO₂ injection, and (iii) with catalyst addition under light irradiation and argon injection. Second, despite the high light absorbance of HEON in the visible light and near-infrared region, the material did not show any photocatalytic activity in these regions within the detection limits of gas chromatographs. Third, despite the

higher feasibility of CH_4 production compared to CO generation in terms of thermodynamics, no CH_4 was detected for these materials which can be explained by the kinetics of reactions. Due to the requirement of CH_4 production to more electrons and protons, its production is not kinetically more feasible than CO production [39,40]. On the other hand, once CO is produced, it does not tend to be adsorbed on active sites and thus the reaction terminates with the CO production [41].

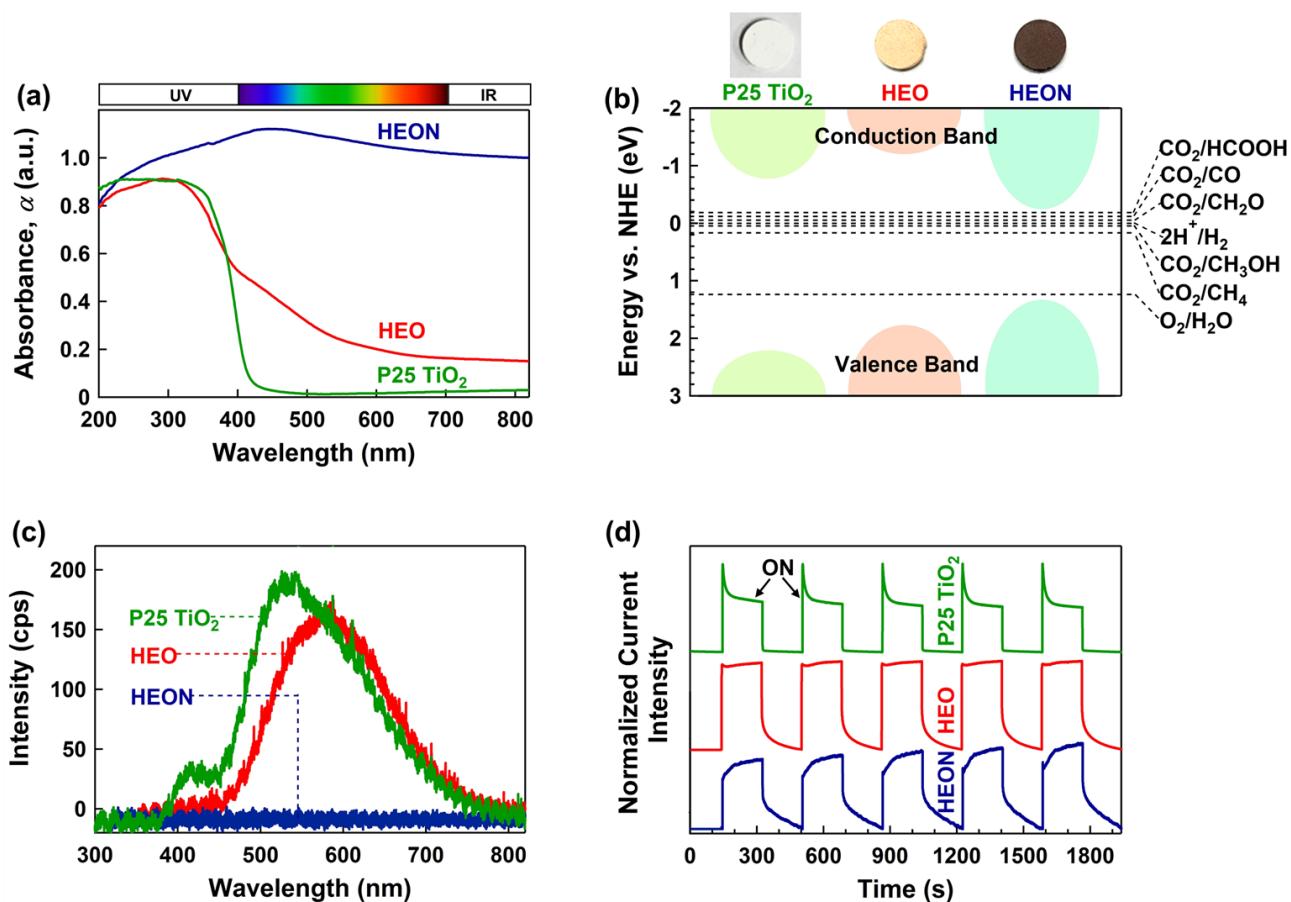


Fig. 5.3. High light absorbance, appropriate electronic band structure and low electron-hole recombination in high-entropy oxynitride. a) UV-vis spectra, b) electronic band structures including chemical potential for CO_2 conversion reactions and sample color, c) steady-state photoluminescence spectra and d) photocurrent generation of high-entropy oxynitride (HEON) in comparison with corresponding high-entropy oxide (HEO) and P25 TiO_2 .

Fig. 5.4c illustrates the XRD profiles of HEON before and after photocatalysis. The profiles indicate that the crystal structures do not change after photoreduction, suggesting that the

HEON remains stable after photocatalytic CO₂ conversion. The high stability of TiZrHfNbTaO₆N₃ is partly because of the entropy-stabilization concept which leads to low Gibbs free energy in the presence of a large number of elements [13,14]. This high stability is an important issue that has led to the utilization of high-entropy ceramics for various applications with superior performance [10-30].

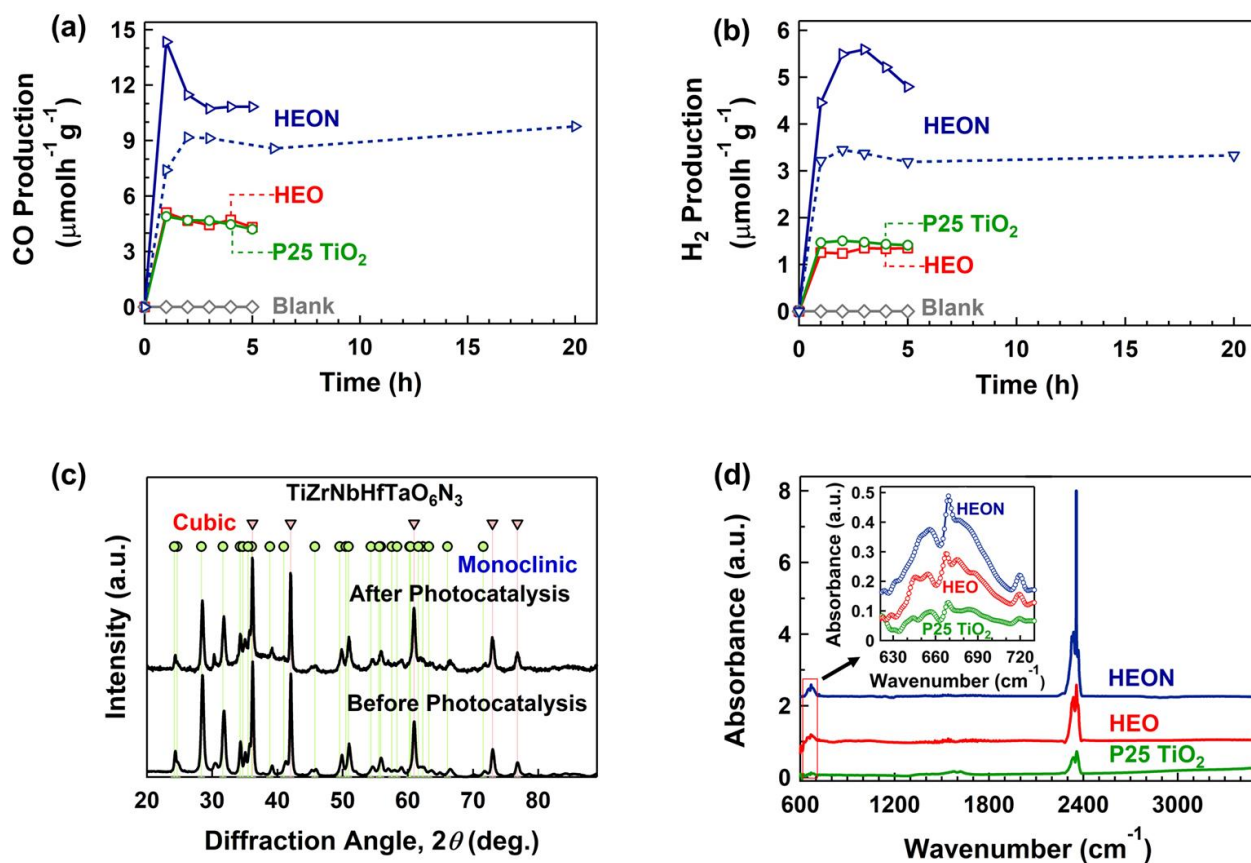


Fig. 5.4. High efficiency of high-entropy oxynitride for photocatalytic CO and hydrogen production. Rate of (a) CO₂ to CO photoreduction and (b) hydrogen generation versus UV irradiation time for high-entropy oxynitride (HEON) compared to corresponding high-entropy oxide (HEO) and P25 TiO₂. (c) XRD profiles before and after photocatalysis for high-entropy oxynitride. (d) DRIFT spectra of three samples.

Fig. 5.4d shows the DRIFT spectra for the three samples to investigate the adsorbance mode of CO₂ on the surface of each photocatalyst. There is a peak at 665 cm⁻¹ which corresponds to CO₃²⁻ [423] and another one at 2340-2360 cm⁻¹ which is relevant to CO₂ gas in the beamline of

spectrometer or to physically adsorbed CO₂ on photocatalysts [43]. The intensity of both peaks is the maximum for the HEON, but it is hard to discuss about physically adsorbed CO₂ using the peak at 2340-2360 cm⁻¹ due to the possible differences in the CO₂ gas concentration in the beamline. The peak at 665 cm⁻¹ for P25 TiO₂ is so weak, but its intensity is the highest for the HEON, suggesting CO₂ can bond to the surface as carbonate. Since CO₂ is a Lewis acid, the basic active sites have a significant role in the adsorption and activation of this molecule [44]. P25 TiO₂ is considered a weak acid and chemisorption of CO₂ in the form of carbonate is weak on the surface of this material. The high intensity of carbonate peak on the HEON suggests that the concentration of basic active sites is higher in this material. These DRIFT experiments indicate the higher capability of the HEON for physisorption and chemisorption of CO₂ compared to the HEO and P25 TiO₂.

5.4. Discussion

Three points need to be discussed in detail here: (i) the mechanism of photocatalytic CO production, (ii) the reasons for the high activity of the HEON, and (iii) the comparison of the activity of current HEON with other photocatalysts reported so far in the literature.

Regarding the first issue, it should be noted that the first step in photocatalytic CO₂ reduction process is the formation of CO₂^{•-} intermediate which is produced by sharing the electrons between CO₂ and photocatalyst surface [45]. Chemisorption of CO₂ molecules on photocatalyst surface to produce CO₂^{•-} occurs in three modes. (1) Nucleophilic bonding between oxygen atoms and catalyst surface (oxygen coordination), (2) electrophilic bonding between carbon atoms and catalyst surface (carbon coordination) and (3) mixed coordination between both oxygen and carbon atoms in CO₂ molecules with catalyst surface [45]. The chemistry of photocatalyst influences the bonding of CO₂^{•-} with catalyst surface and determine the reaction pathway [45]. If a photocatalyst contains Sn, Pb, Hg, In, and Cd metals, then it has a tendency to oxygen coordination to produce [•]OCHO as an intermediate and formic acid (HCOOH) as the final product. Photocatalysts containing noble and transition metals have a tendency to carbon coordination which leads to producing [•]CO and [•]OCHO as intermediates. Since the bonding between [•]CO and catalyst surface is weak, CO is usually the main product in this coordination. If a photocatalyst consists of Cu atoms, then [•]CO and [•]OCHO are produced as intermediates, but because of the strong bonding between [•]CO and Cu, other hydrocarbons such as methane (CH₄) and ethanol

(C₂H₅OH) are usually formed as final products [45]. In this study, since the HEON, the HEO and P25 TiO₂ consist of transition metals, the carbon coordination pathway occurs for these photocatalysts which leads to CO production. This pathway has the following reactions [45].



Regarding the second issue, it should be considered that combining the concepts of metal oxynitrides and high-entropy ceramics was the spark starter of this study. Metal oxynitrides have been introduced as promising low-bandgap photocatalysts particularly for water splitting [5], while their application for CO₂ conversion is still in the initial steps [6-8]. The reason for the low bandgap of oxynitrides compared to oxide photocatalysts is that the valence band top of these materials is generated using hybridized 2p oxygen and nitrogen orbitals, but it is generated using only 2p oxygen orbitals in oxides. Since the energy level for nitrogen 2p orbitals is higher than oxygen 2p orbitals, the bandgap of oxynitrides is smaller than oxides [5]. However, these metal oxynitrides suffer from significant recombination of charge carriers and modest stability [6,9].

High-entropy ceramics are promising new materials with interesting properties because of the presence of multiple elements which leads to superior stability, large lattice defects/strain and heterogenous valence electron distribution [13,14]. The presence of various elements in the lattice of high-entropy ceramics leads not only to high configurational entropy and resultant high chemical stability for catalysis but also to lattice distortion and formation of inherent point defects such as vacancies which can act as active sites for catalysis [46]. Although future theoretical studies are required to determine the active sites in high-entropy photocatalysts, it was shown in conventional photocatalysts that vacancies on the surface adsorb CO₂ and activate it by decreasing the bonding energy between carbon and oxygen [41,45]. These vacancies trap electrons and act as active sites and lead to the improved photocatalytic activity for CO₂ reduction, but their fraction should be optimized to achieve the highest activity and best reaction selectivity [47,48].

The combination of the two concepts of oxynitrides and high-entropy ceramics led to the introduction of TiZrNbHfTaO₆N₃ as a highly stable and low-bandgap photocatalyst for CO₂ conversion with much better photocatalytic performance compared with P25 TiO₂. Such a high activity is particularly interesting because the surface area of the HEON is much smaller than P25

TiO₂: 2.3 m²g⁻¹ for the HEON and 38.7 m²g⁻¹ for P25 TiO₂ measured by the Brunauer-Emmett-Teller (BET) technique in the nitrogen atmosphere. The high activity of the HEON for CO₂ photoreduction can be attributed to high light absorbance (i.e., easy electron-hole separation), appropriate band positions compared to chemical potentials for reactions, and low electron-hole recombination, and high surface CO₂ adsorption [34,37]. The presence of interphase boundaries in this HEON can also partly contribute to the easy separation of charge carriers and improvement of photocatalytic activity [34].

Regarding the third issue, although the comparison between the current HEON and P25 TiO₂ using similar experimental procedures confirms the high photocatalytic activity of the HEON, it is worth comparing the activity of this HEON with the given data in the literature. Photocatalysis in various studies is performed in different conditions in terms of photoreactor type, temperature, catalyst concentration, CO₂ flow rate, type of light source and concentration of reactants, and thus, a comparison between different studies should be evaluated with care. The CO production rate per catalyst mass and catalyst surface area are given in Table 5.1 in comparison with reported photocatalysts in the literature [48-79]. Since photocatalysis occurs on the surface, normalizing the CO production rate per surface area should be more reasonable for comparison purposes. According to Table 5.1, the rate of CO production in the literature fluctuates in the 0.00095-1.33 μmolh⁻¹m⁻¹ range, while the CO generation rate of 4.66 ± 0.3 μmolh⁻¹m⁻¹ on the HEON is higher than all these reported data. Moreover, the HEON shows much better activity per both surface area and mass unit compared to other oxynitrides reported in the literature. Although these findings introduce HEON as the most effective photocatalyst for CO₂ photoreduction, future studies should focus on decreasing the particle size of these materials to increase their specific surface area.

Table 5.1. Photocatalytic CO₂ to CO conversion rate on high-entropy oxynitride compared to reported photocatalysts. For some catalysts which surface area was not reported in literature, CO production rate in μmolh⁻¹m⁻¹ was not given.

Photocatalyst	Catalyst Concentration	Light Source	CO Production Rate (μmolh ⁻¹ g ⁻¹)	CO Production Rate (μmolh ⁻¹ m ⁻¹)	Ref.
TiO ₂ / Carbon Nitride Nanosheet	25 mg (Gas System)	150 W Xenon	2.04	----	[49]
TiO ₂ / Graphitic Carbon	100 mg (Gas System)	300 W Xenon	10.16	0.04	[50]
TiO ₂ Nanosheets Exposed {001} Facet	1 gL ⁻¹ (Liquid system)	Two 18 W Low-Pressure Mercury	0.12	0.00095	[51]
TiO ₂ / CoOx Hydrogenated	50 mg (Gas System)	150 W UV	1.24	0.0045	[52]
TiO ₂ 3D Ordered Microporous / Pd	100 mg (gas system)	300 W Xenon	3.9	0.066	[53]
Pt ²⁺ -Pt ⁰ / TiO ₂	100 mg (Gas System)	300 W Xenon	~12.14	0.7	[54]

Anatase TiO ₂ Hierarchical Microspheres	200 mg (Gas System)	40 W Mercury UV	18.5	0.37	[55]
TiO ₂ and Zn(II) Porphyrin Mixed Phases	60 mg (Gas System)	300 W Xenon	8	0.062	[56]
Anatase TiO ₂ Hollow Sphere	100 mg (Gas System)	40 W Mercury UV	14	0.16	[57]
Anatase TiO ₂ Nanofibers	50 gL ⁻¹ (Liquid System)	500 W Mercury Flash	40	----	[58]
C ₃ N ₄ by Thermal Condensation	100 mg (Gas System)	350 W Mercury	4.83	-----	[59]
Cd _{1-x} Zn _x S	45 mg (Gas System)	UV-LED irradiation	2.9	0.015	[60]
BiOI	150 mg (Gas System)	300 W High-Pressure xenon	4.1	0.03	[61]
xCu ₂ O / Zn _{2-2x} Cr	4 gL ⁻¹ (Liquid System)	200 W Mercury-Xenon	2.5	0.018	[62]
CeO _{2-x}	50 mg (Gas System)	300 W Xenon	1.65	0.08	[63]
Cu ₂ O / RuO _x	500 mg (Gas System)	150 W Xenon	0.88	---	[64]
Bi ₂ Sn ₂ O ₇	0.4 gL ⁻¹ (Liquid System)	300 W xenon	14.88	0.24	[65]
Ag / Bi / BiVO ₄	10 mg (Gas System)	300W Xenon Illuminator	5.19	0.42	[66]
g-C ₃ N ₄ / BiOCl	20 mg (Gas System)	300 W High-Pressure Xenon	4.73	---	[67]
Fe / g-C ₃ N ₄	1 gL ⁻¹ (Liquid System)	300 W Xenon	-22.5	0.06	[68]
Bi ₂ MoO ₆	0.7 gL ⁻¹ (Liquid System)	300 W Xenon	41.5	1.26	[69]
Bi ₂₄ O ₃₁ C ₁₁₀	50 mg (Gas System)	300 W High-Pressure Xenon	0.9	---	[70]
g-C ₃ N ₄ / Zinc Carbodiimide / Zeolitic Imidazolate Framework	100 mg (Gas System)	300 W Xenon	-0.45	0.014	[71]
BiVO ₄ / C / Cu ₂ O	---	300 W Xenon	3.01	----	[72]
g-C ₃ N ₄ / α-Fe ₂ O ₃	200 mg (Gas System)	300 W Xenon	5.7	-----	[73]
Bicrystalline Anatase/Brookite TiO ₂ Microspheres	30 mg (Gas System)	150 W Solar Simulator	145	0.95	[74]
10 wt % In-Doped Anatase TiO ₂	250 mg (Gas System)	500 W Mercury Flash	81	1.33	[75]
10 wt % Montmorillonite-Loaded TiO ₂	50 mg (Gas System)	500 W Mercury	103	1.25	[76]
Bi ₄ O ₅ Br ₂	20 mg (Gas System)	300 W High-Pressure Xenon	63.13	0.58	[77]
ZnGaON	---	1600 W Xenon	1.05	---	[78]
WO ₃ / LaTiO ₂ N	10 mg (Gas System)	300 W Xenon	2.21	0.4	[79]
α-Fe ₂ O ₃ / LaTiO ₂ N	20 mg (Gas System)	300 W Xenon	9.7	0.65	[6]
RuRu / Ag / TaON	1 gL ⁻¹ (Liquid System)	High-Pressure Mercury	5	----	[7]
RuRu / TaON	1 gL ⁻¹ (Liquid system)	High-Pressure Mercury	3.33	----	[7]
Ag / TaON / RuBLRu'	2 gL ⁻¹ (Liquid System)	500 W High-Pressure Mercury	0.056	----	[8]
TiZrHfNbTaO ₆ N ₃	0.2 gL ⁻¹ (Liquid System)	400 W High-Pressure Mercury	10.7 ± 1.8	4.7 ± 0.3	This study

5.4. Conclusion

This chapter introduced a high-entropy oxynitride with low bandgap, low electron-hole recombination, high CO₂ adsorbance and high chemical stability for photocatalytic CO₂ conversion. The material, which was synthesized using high-pressure torsion and subsequent oxidation and nitriding, had two phases of face-centered cubic and monoclinic with a chemical composition of TiZrNbHfTaO₆N₃. The material had better photocatalytic CO₂ conversion performance compared with corresponding high-entropy oxide, benchmark photocatalyst P25 TiO₂ and all reported photocatalysts in the literature. These findings open a new path to developing highly efficient photocatalysts for CO₂ conversion.

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Chapter 6. Concluding remarks and outlook

Global warming has become a significant concern in recent years which seriously threatens the life of humans and creatures. Conversion of CO₂ molecules to other components such as CO is a way to stand this event. In this regard, photocatalytic CO₂ conversion, which uses solar irradiation as a clean energy source, has been introduced as a new and promising strategy in recent years. Despite the introduction of various materials, which are modified by various strategies, the efficiency of CO₂ photoreduction is still low compared to conventional methods for CO₂ conversion, as discussed in chapter 1 of this thesis. The high-pressure torsion (HPT) as a severe plastic deformation (SPD) method has been used in this research work to produce some of the most active photocatalysts for CO₂ conversion. The HPT method could increase the CO₂ photoreduction efficiency by (i) oxygen vacancy and strain engineering, (ii) stabilization of high-pressure phases, (iii) formation of defective high-entropy oxides, and (iv) synthesis of low-bandgap oxynitrides.

In chapter 2, a photocatalyst with low bandgap, improved band structure and low recombination rate of electrons and holes was produced by simultaneous introduction of oxygen vacancies and lattice strain in BiVO₄ via the HPT process. The material showed high photocatalytic activity for CO₂ to CO conversion with an activity comparable to P25 TiO₂ as a benchmark photocatalyst. This chapter suggests that simultaneous engineering of lattice strain and defects is an effective strategy to produce active photocatalysts for CO₂ conversion.

In chapter 3, effect of high pressure TiO₂-II phase on photocatalytic CO₂ conversion was examined for the first time. The introduction of nanostructured TiO₂-II phase and TiO₂-II/anatase heterojunctions modified the band structure and enhanced the photocurrent and photocatalytic activity, despite the formation of strain-induced vacancies in the bulk. Thermal annihilation of vacancies in the bulk led to further improvement of photocatalytic CO₂ conversion, confirming the significance of the TiO₂-II phase for photocatalytic reactions.

In chapter 4, a high-entropy oxide with a general composition of TiZrNbHfTaO₁₁ was synthesized and used for photocatalytic CO₂ conversion. Due to appropriate electronic band structure, good charge carrier lifetime and a defective and strained dual-phase structure, the material acted as a photocatalyst for CO₂ to CO conversion and H₂O to H₂ production without

addition of any co-catalyst. The photocatalytic activity of this oxide was better than well-known anatase TiO_2 and BiVO_4 photocatalysts and comparable with P25 TiO_2 as a benchmark photocatalyst, suggesting high-entropy oxides as a new family of photocatalysts for CO_2 conversion.

In chapter 5, a high-entropy oxynitride with low bandgap, low electron-hole recombination, high CO_2 adsorbance and high chemical stability was introduced for photocatalytic CO_2 conversion. The material, which was synthesized using HPT and subsequent oxidation and nitriding, had two phases of face-centered cubic and monoclinic with a chemical composition of $\text{TiZrNbHfTaO}_6\text{N}_3$. The material had better photocatalytic CO_2 conversion performance compared with corresponding high-entropy oxide, benchmark photocatalyst P25 TiO_2 and all reported photocatalysts in the literature. This chapter opens a new path to developing highly efficient photocatalysts for CO_2 conversion.

All taken together, the main reason for the successful application of HPT to develop highly active photocatalysts for CO_2 conversion is that the method can simultaneously introduce some structural/microstructural features as effective strategies to enhance the activity. It was shown that the HPT method improves photocatalytic activity by increasing light absorbance, decreasing the optical bandgap, optimizing the electronic band structure, accelerating the electron-hole separation and migration and reducing the recombination rate of electrons and holes. The HPT method also contributed to the introduction of high-entropy photocatalysts which show high activity for CO_2 conversion.

Despite the high activity of HPT-processed photocatalysts, the mechanism for their high efficiency needs to be investigated further by employing theoretical calculations. Moreover, the catalysts produced by HPT have usually low surface area and new methods are expected to be developed to enhance the surface area of these catalysts. Although the HPT method has opened a new path to employing high-pressure and high-entropy phases as new photocatalysts, other synthesis methods should be used in the future for the large-scale production of these new catalysts. It is expected that HPT will continue its contribution to the field of photocatalysis, particularly due to the significance of developing a carbon-neutral society in recent years.

Conference Presentations

- [1] S. Akrami, M. Watanabe, T.H. Ling, T. Ishihara, M. Arita, M. Fuji, K. Edalati, High-pressure TiO₂-II polymorph as an active photocatalyst for CO₂ to CO conversion, *NanoSPD Conference*, Bangalore, India, February 26-March 3, 2023.
- [2] S. Akrami, M. Watanabe, T.H. Ling, T. Ishihara, M. Arita, M. Fuji, K. Edalati, High-pressure TiO₂-II polymorph as an active photocatalyst for CO₂ to CO conversion, *The 59th Symposium on Powder Science and Technology*, Kyoto, Japan, December 21-23, 2022.
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- [1] S. Akrami, M. Watanabe, T.H. Ling, T. Ishihara, M. Arita, M. Fuji, K. Edalati, High-pressure TiO₂-II polymorph as an active photocatalyst for CO₂ to CO conversion, *14th Pacific Rim Conference on Ceramic and Glass Technology (PACRIM 14)*, Vancouver, Canada, December 14, 2021.

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- [1] S. Akrami, P. Edalati, Y. Shundo, M. Watanabe, T. Ishihara, M. Fuji, K. Edalati, Significant CO₂ photoreduction on high-entropy oxynitride, *Chemical Engineering Journal*, Vol. 449, p. 137800.
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Acknowledgments

I would like to express my deepest gratitude and sincere thanks to my supervisor, Prof. Masayoshi Fuji, for his full support and thoughtful guidance. I also appreciate Dr. Kaveh Edalati for his valuable contributions to perform this research work. I am also very grateful to Prof. Nobuyasu Adachi, Prof. Masaaki Haneda and Dr. Chika Takai for serving on my thesis committee and taking their time and energy to appraise my thesis. In addition, I am grateful for the support and encouragement that I have received from my parents and my wife during this difficult period of my life.